

# 活動銀河核の降着流からの ニュートリノ放射

東北大学 (学振PD)

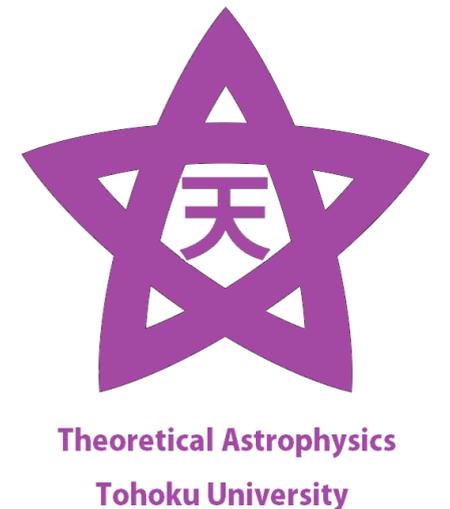
木村成生

## References

- 1) SSK, Murase, Meszaros, 2019, PRD, 100, 083014
  - 2) SSK, Murase, Meszaros in preparation
  - 3) Murase, SSK, Meszaros, arXiv:1904.04226
- see also: SSK, Murase, Toma, 2015, ApJ, 806, 159

## Collaborators

Peter Meszaros (Penn State)  
Kohta Murase (Penn State; YITP)  
Kenji Toma (Tohoku Univ.)



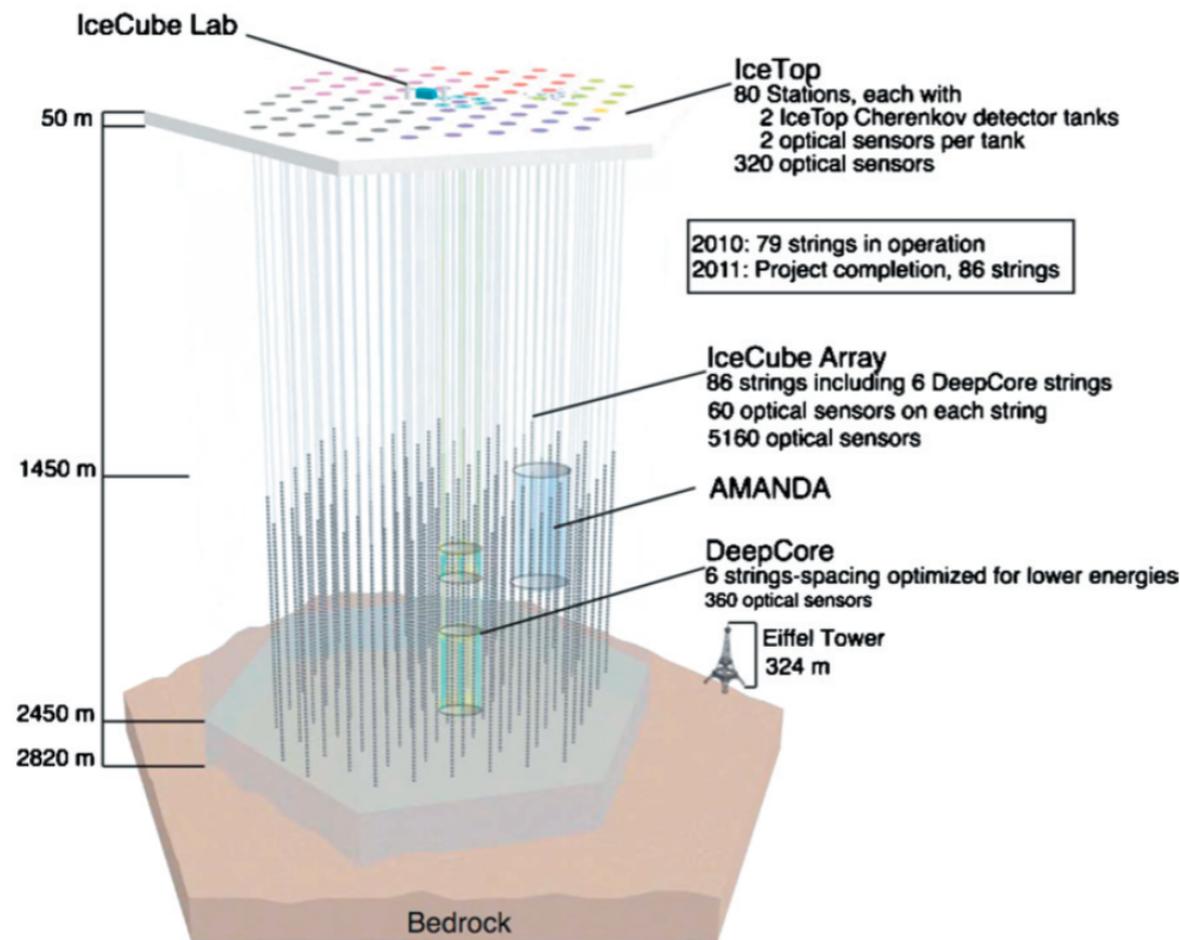
# Index

- IceCube Neutrinos
- Accretion Flow in AGN
- AGN Corona model Murase, SSK, Meszaros, arXiv:1904.04226
- LLAGN RIAF model SSK, Murase, Meszaros, 2019, PRD, 100, 083014  
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- Summary

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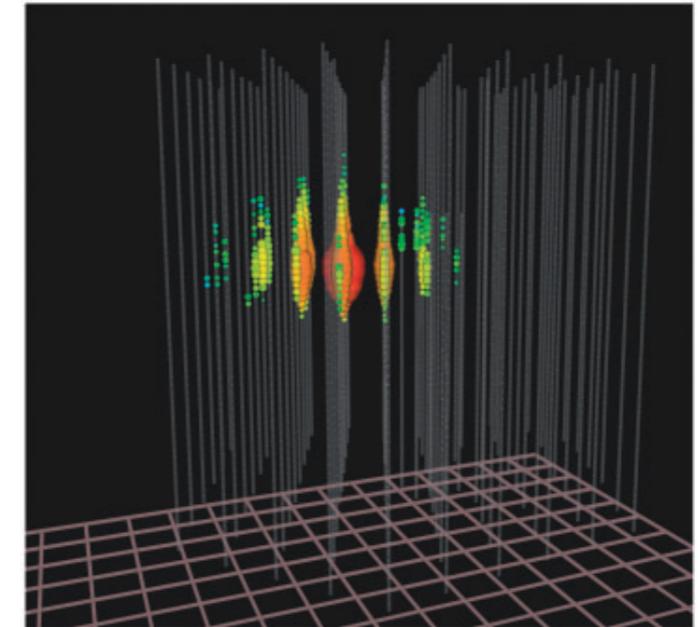
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# Detection of Astrophysical Neutrinos

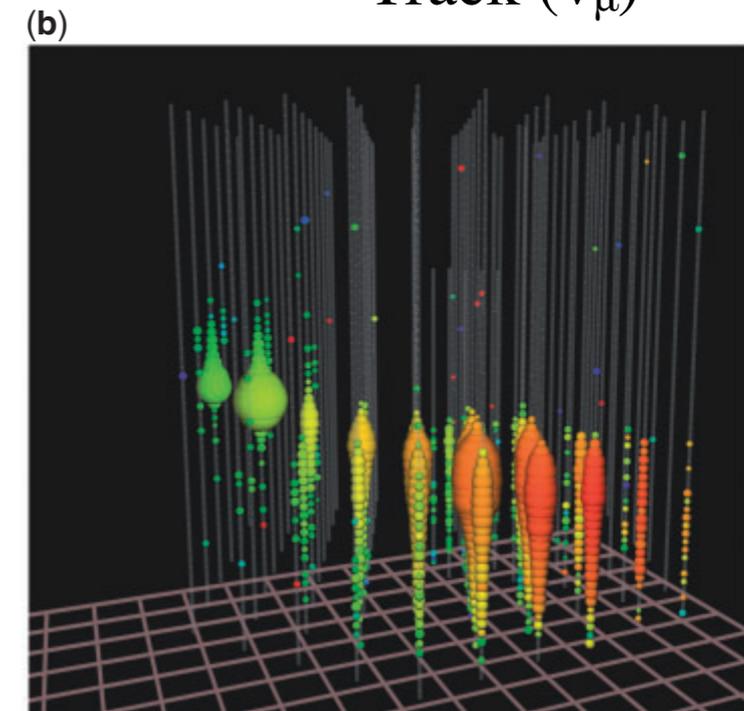


- IceCube experiment reported detection of astro- $\nu$  ( $E \sim \text{PeV}$ ) in 2013
- Shower: good for spectrum
- Track: good for source search

(a) Shower or Cascade ( $\nu_e$ )

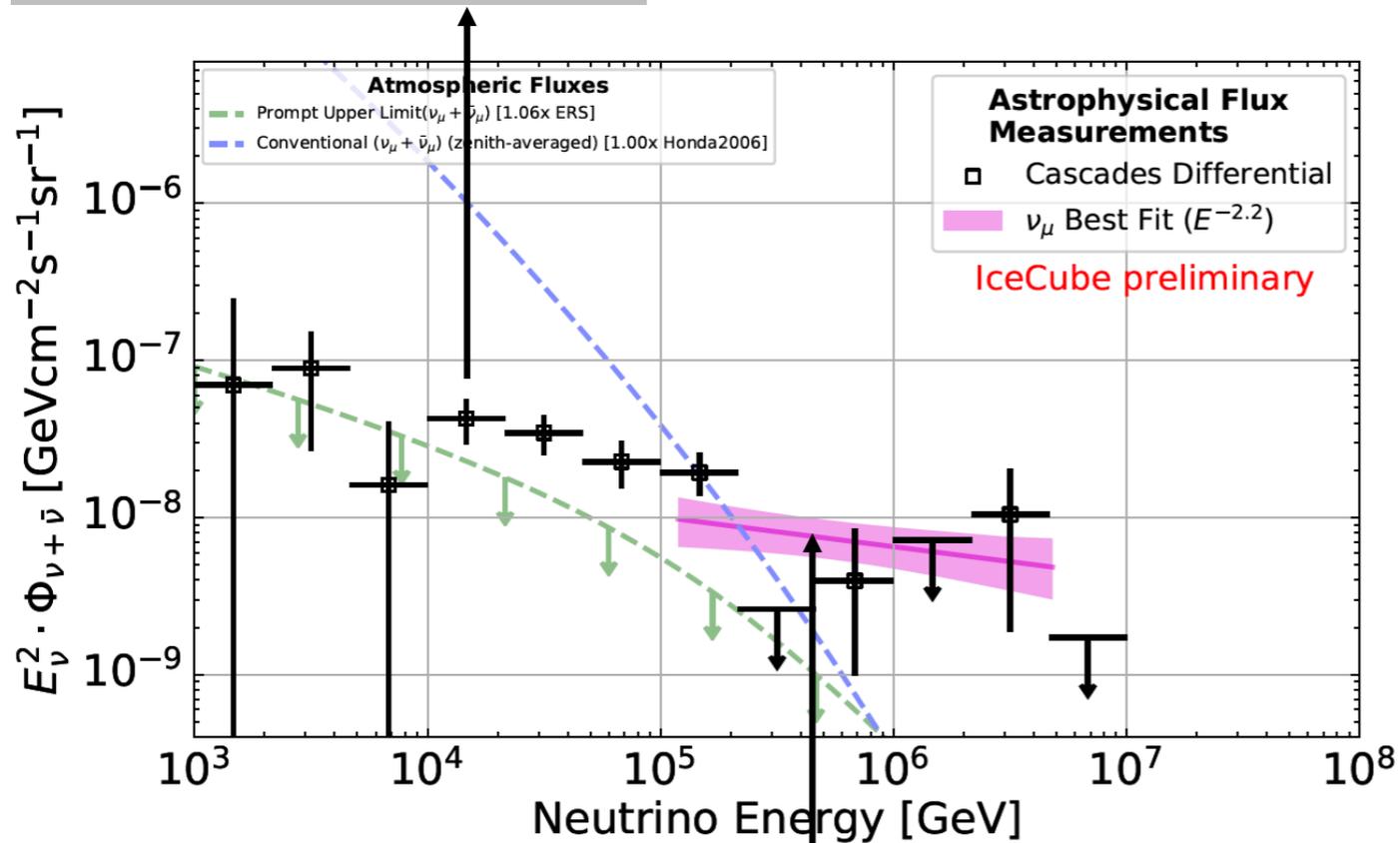


(b) Track ( $\nu_\mu$ )

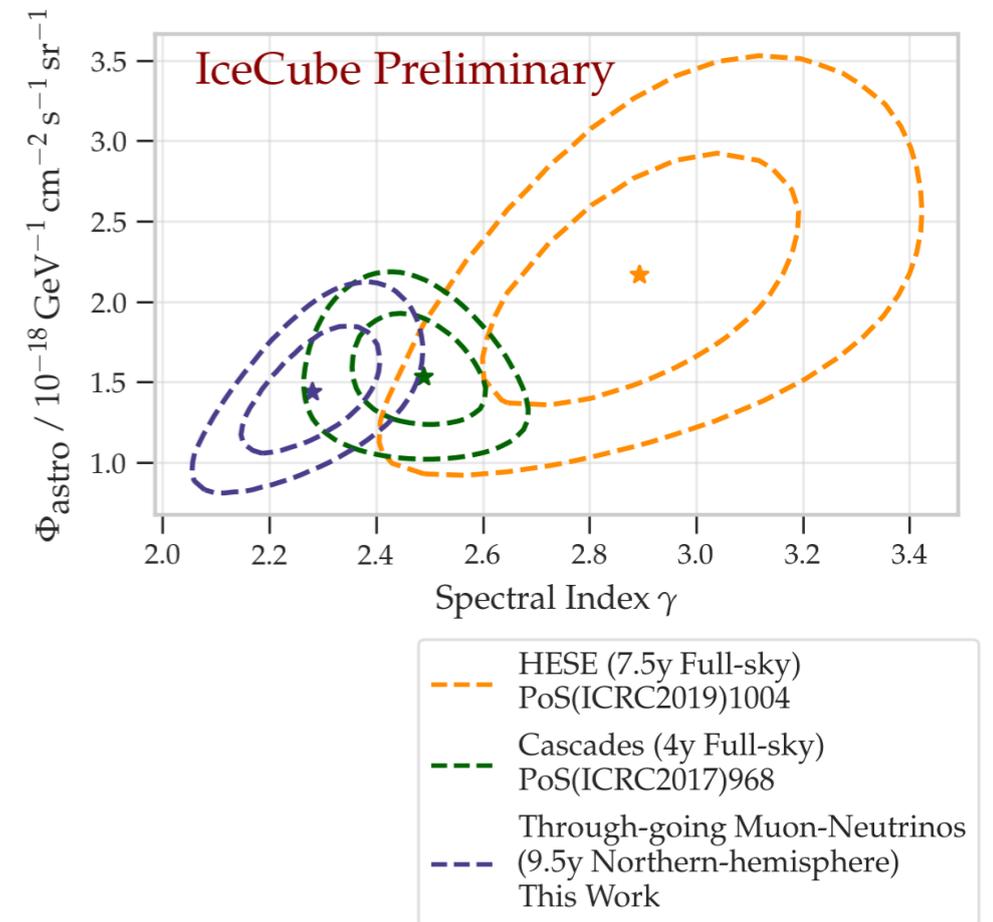


# Neutrino Spectrum

electron and tau neutrinos



muon neutrinos

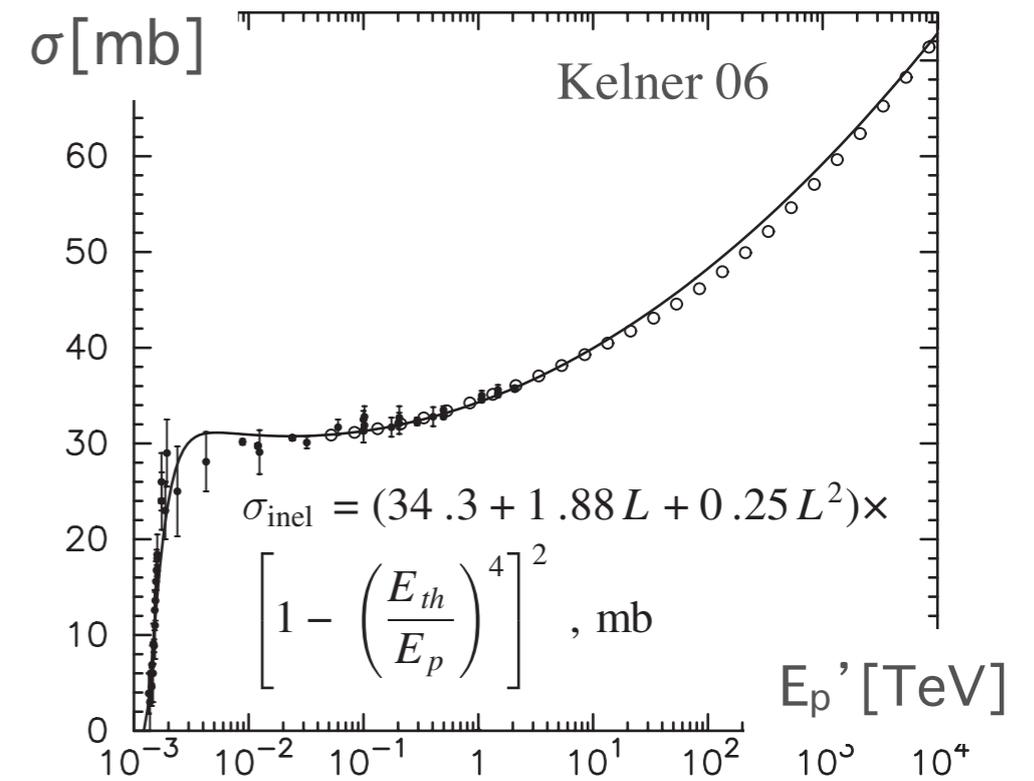
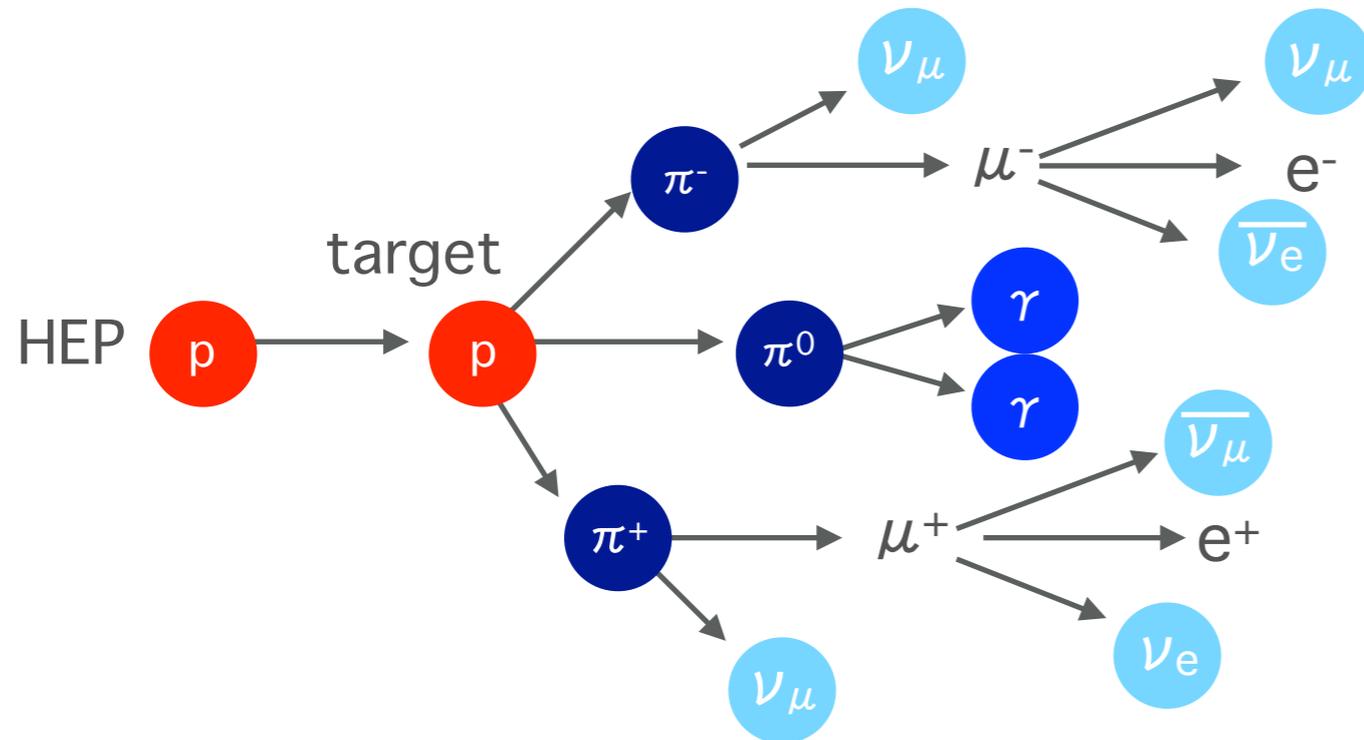


- Track analysis: flat spectrum ( $E > 200 \text{ TeV}$ )
- Cascade analysis: soft spectrum ( $E > 1 \text{ TeV}$ )
- Hint of 2 component??      Uncertainty of analyses??

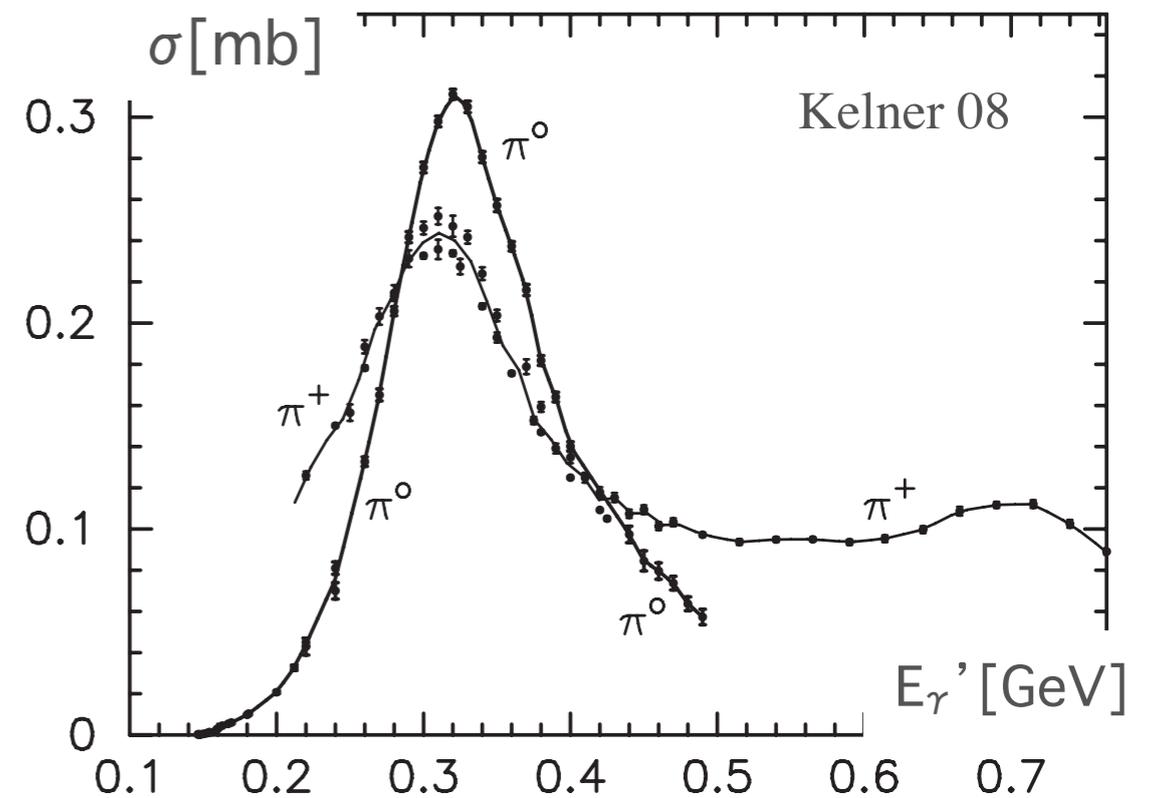
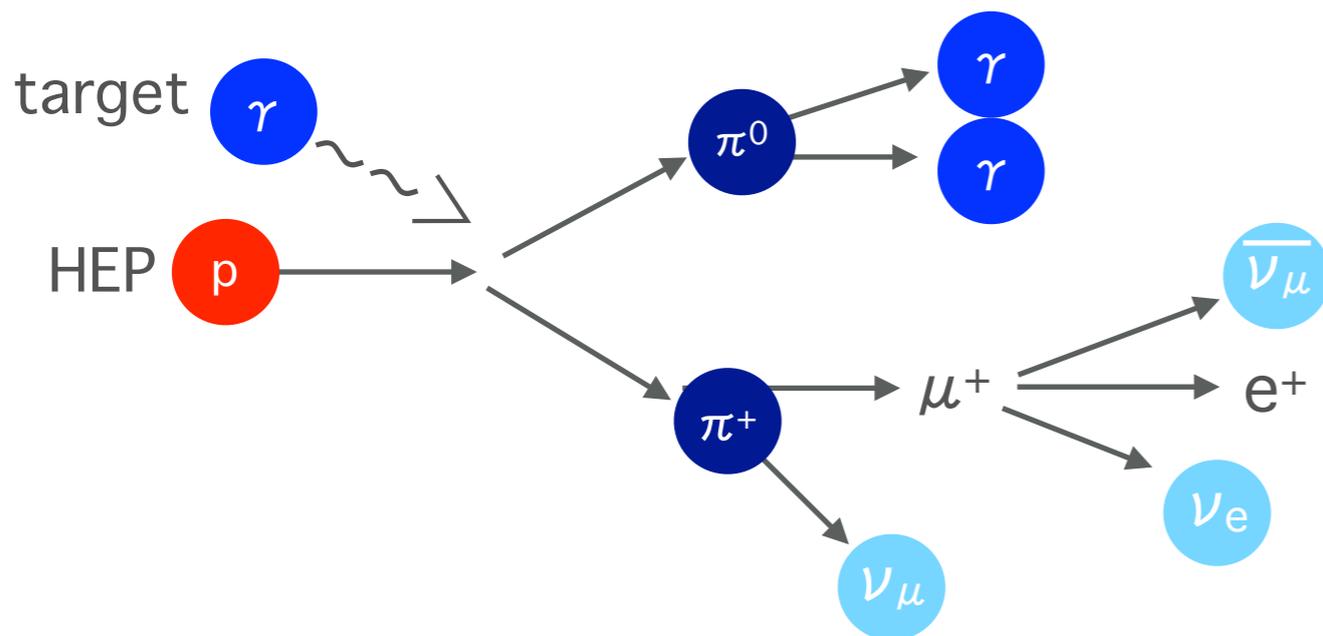


# Neutrino Production Process

- pp inelastic collision

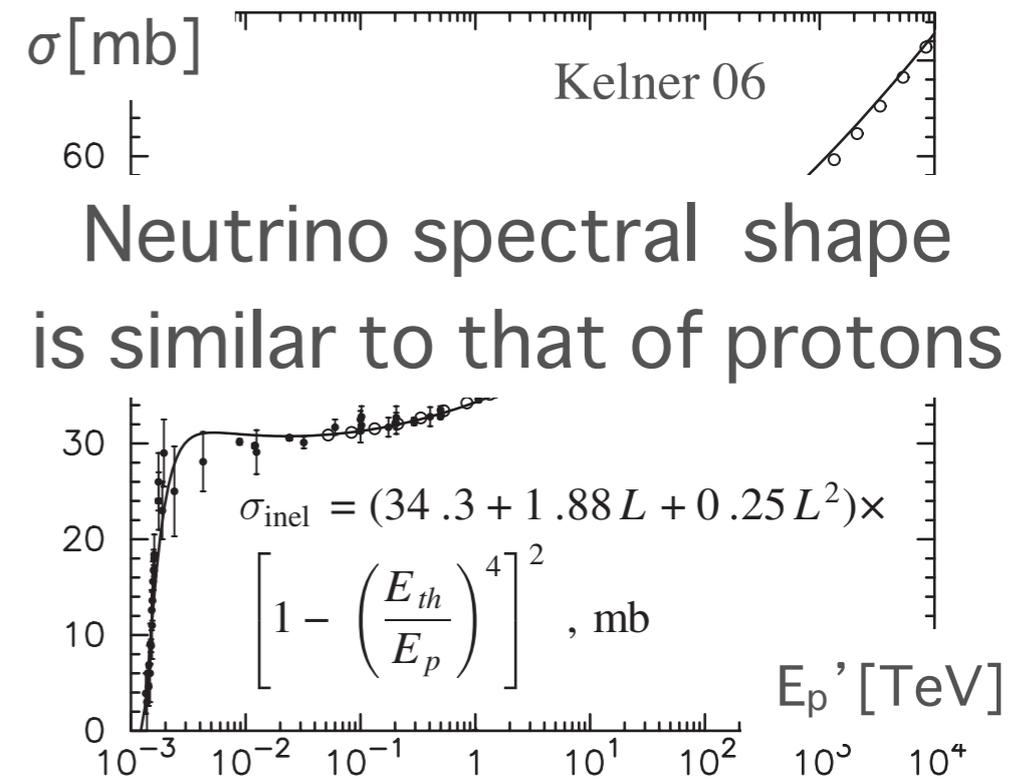
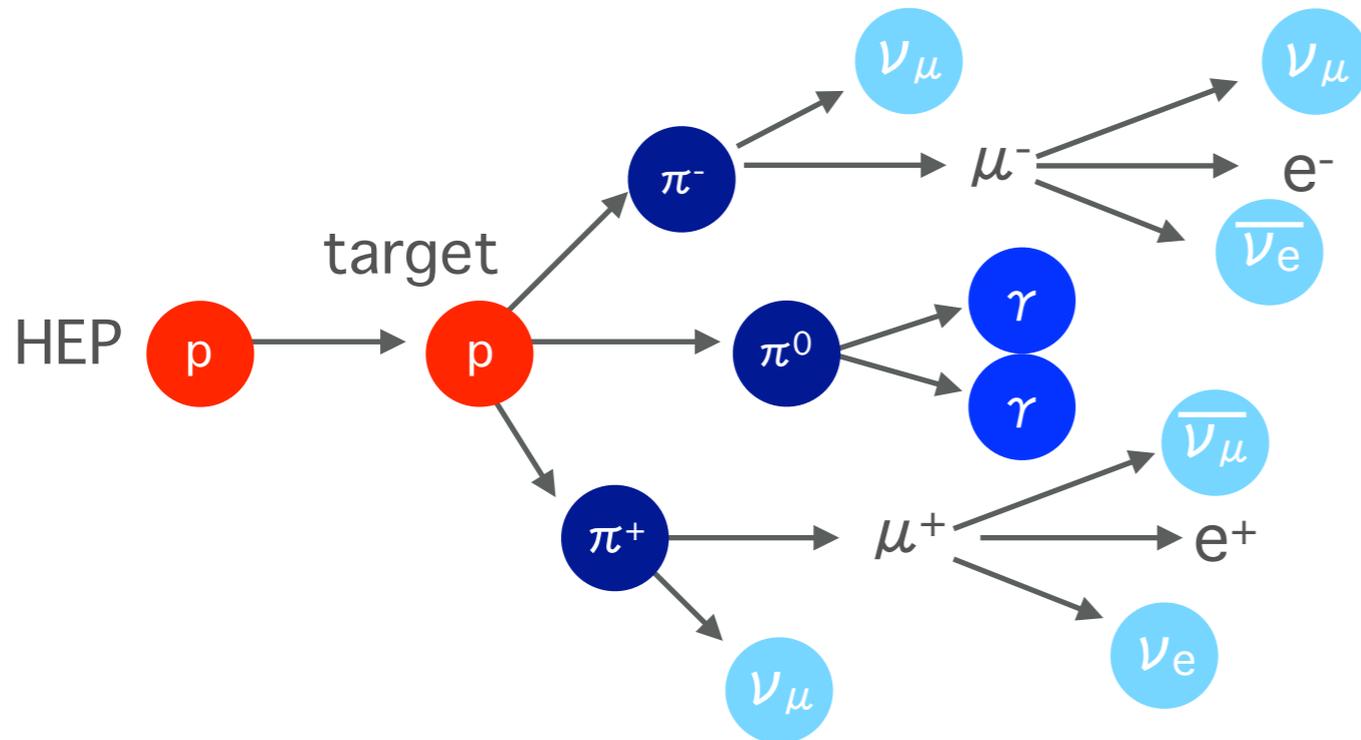


- Photomeson production ( $p\gamma$ )

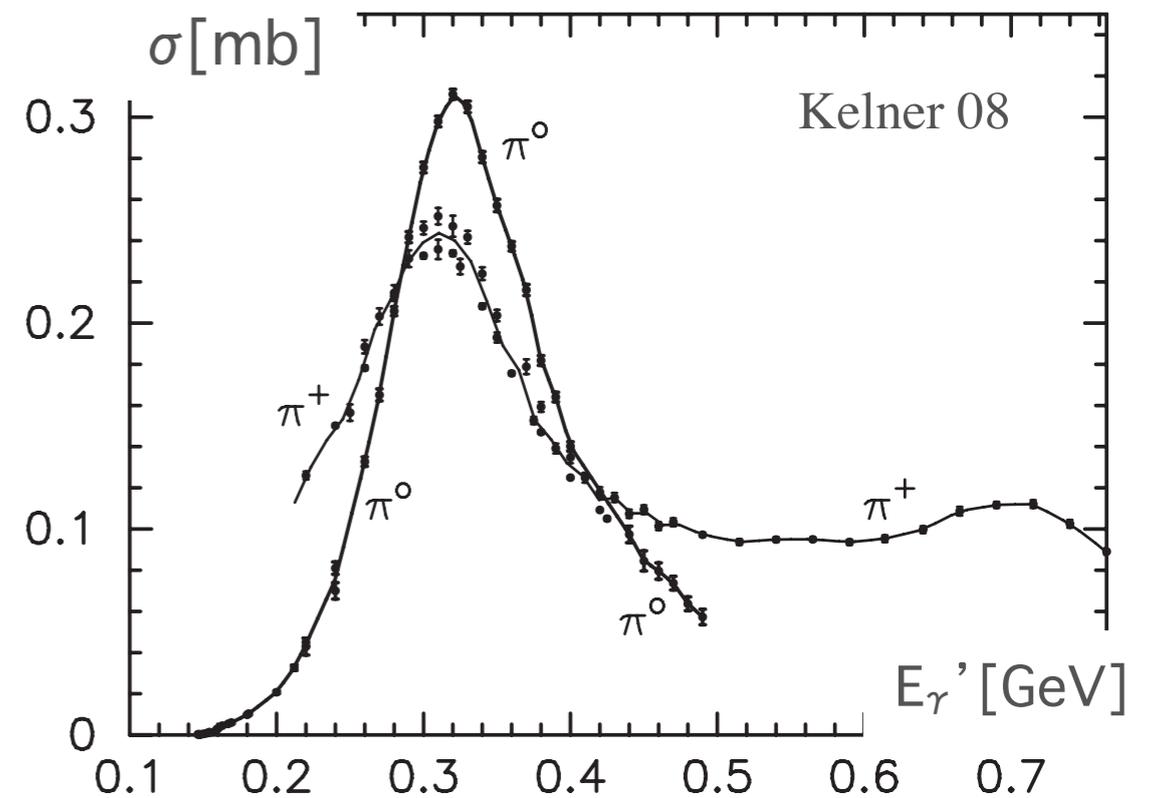
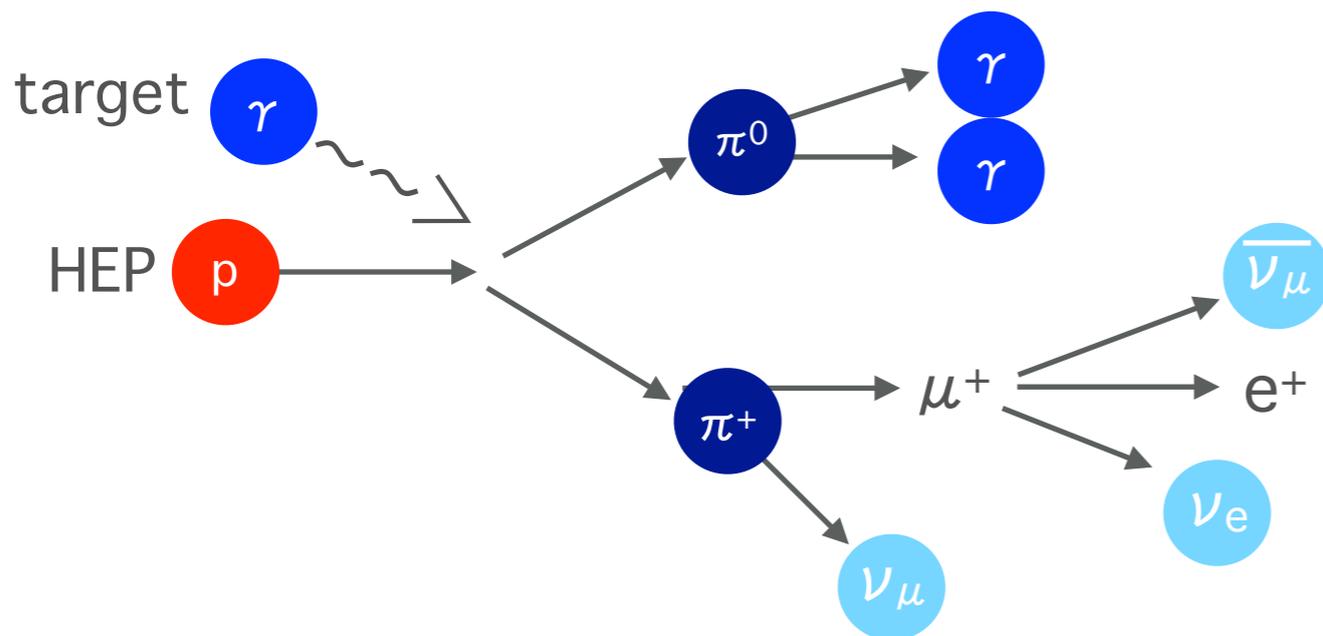


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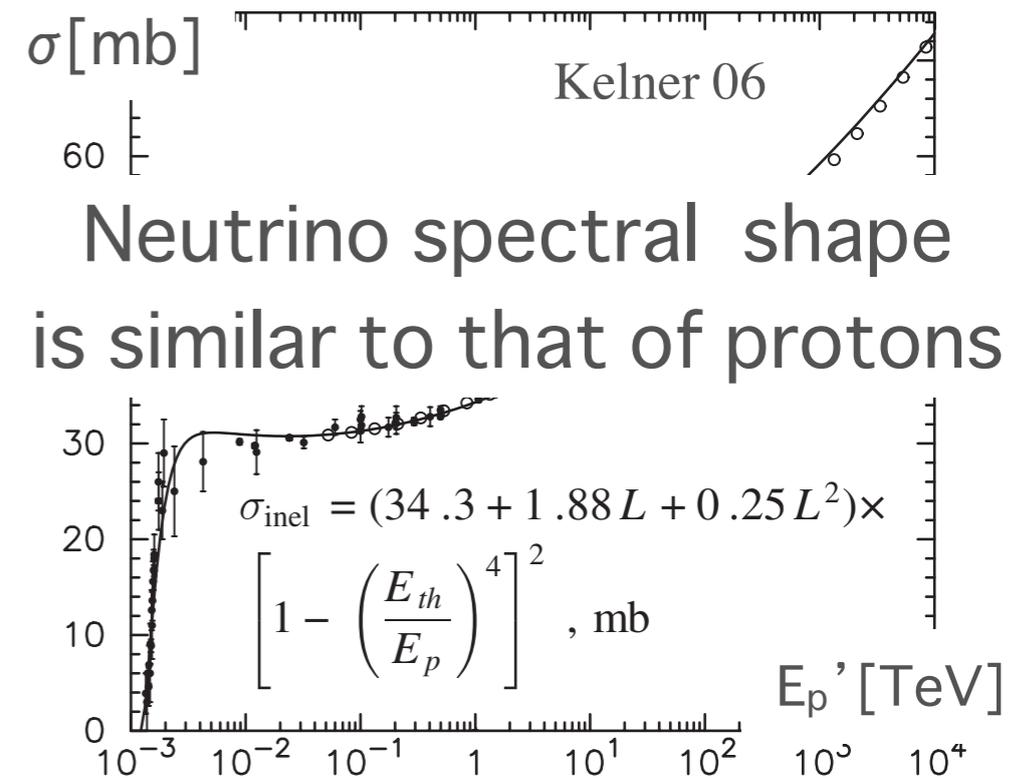
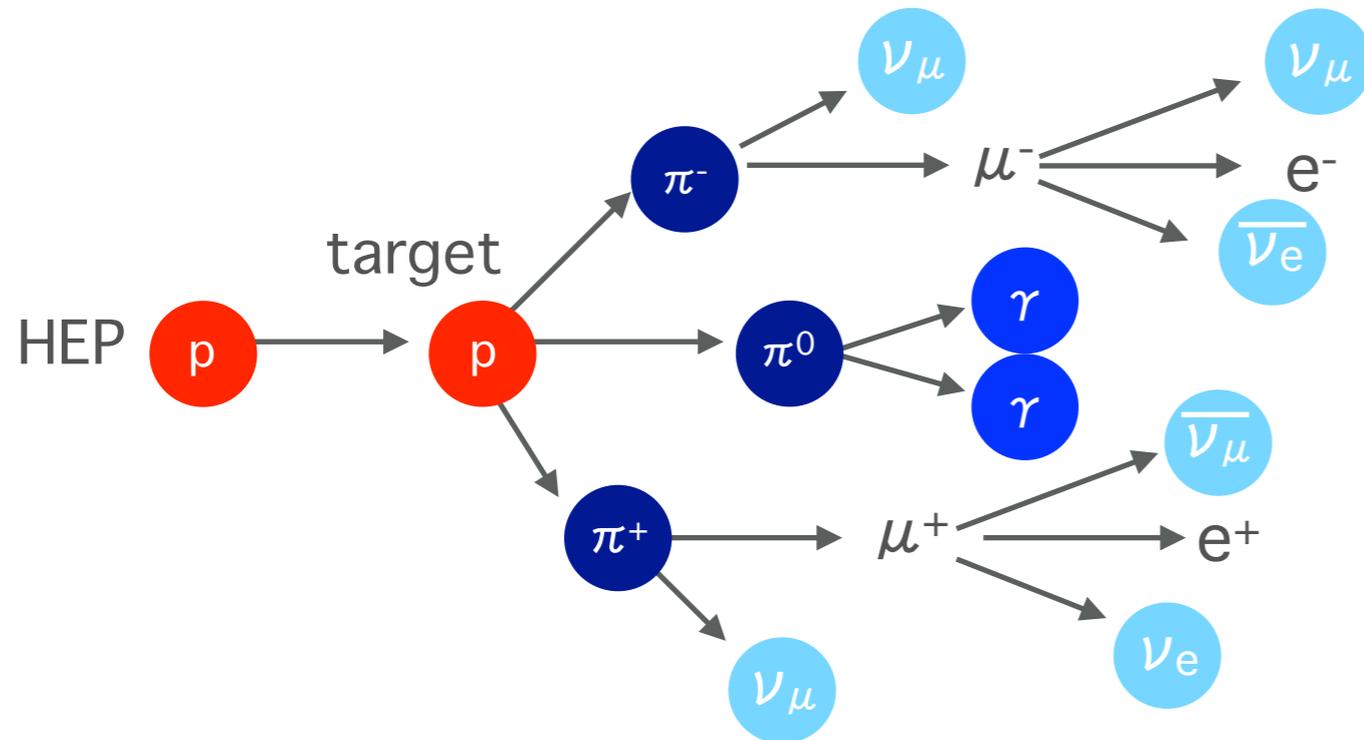


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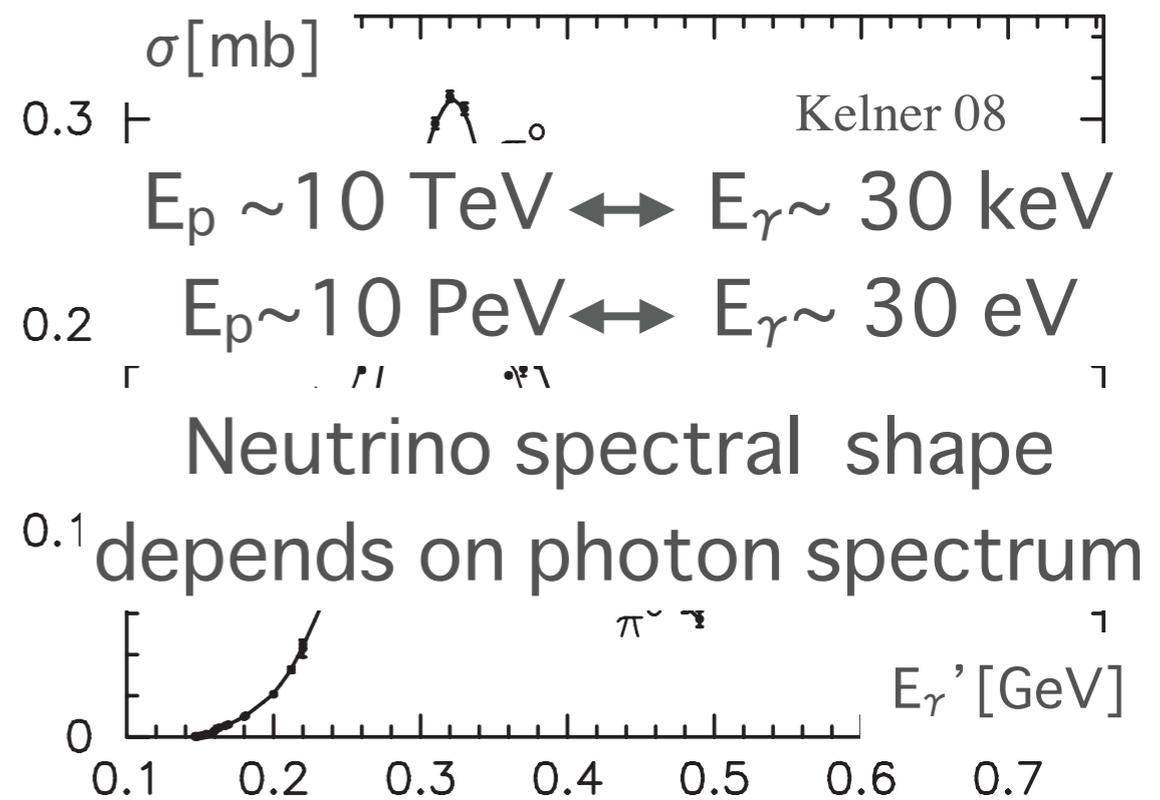
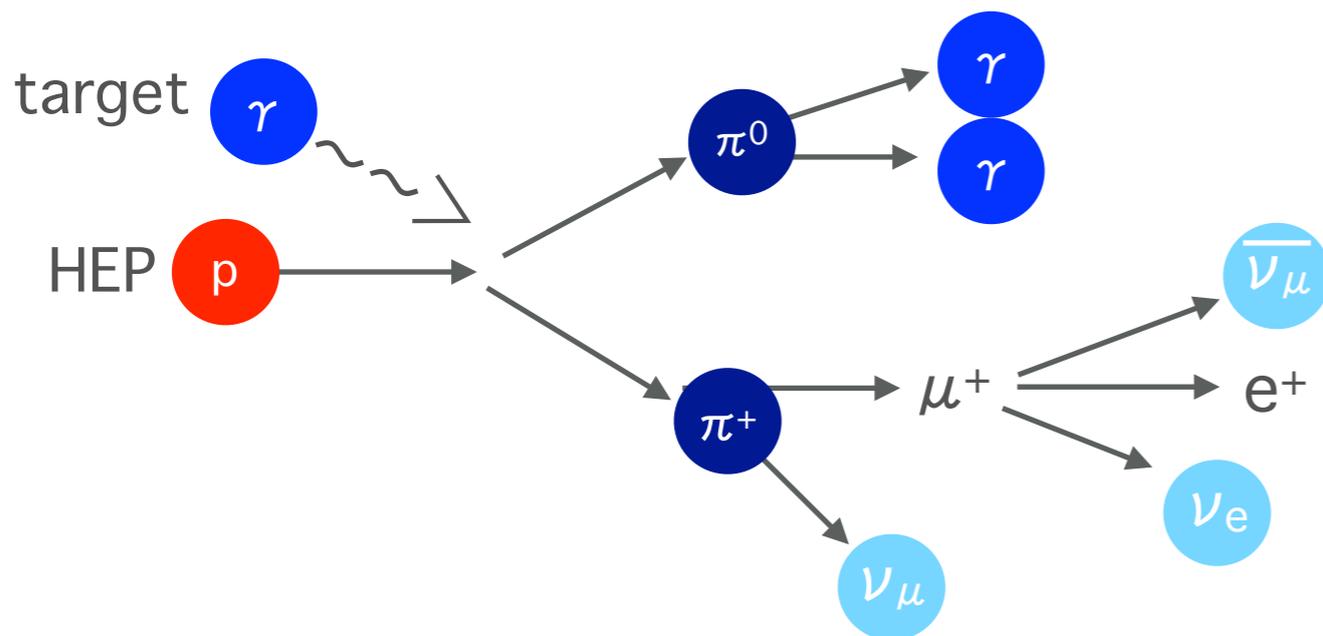
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- pp inelastic collision



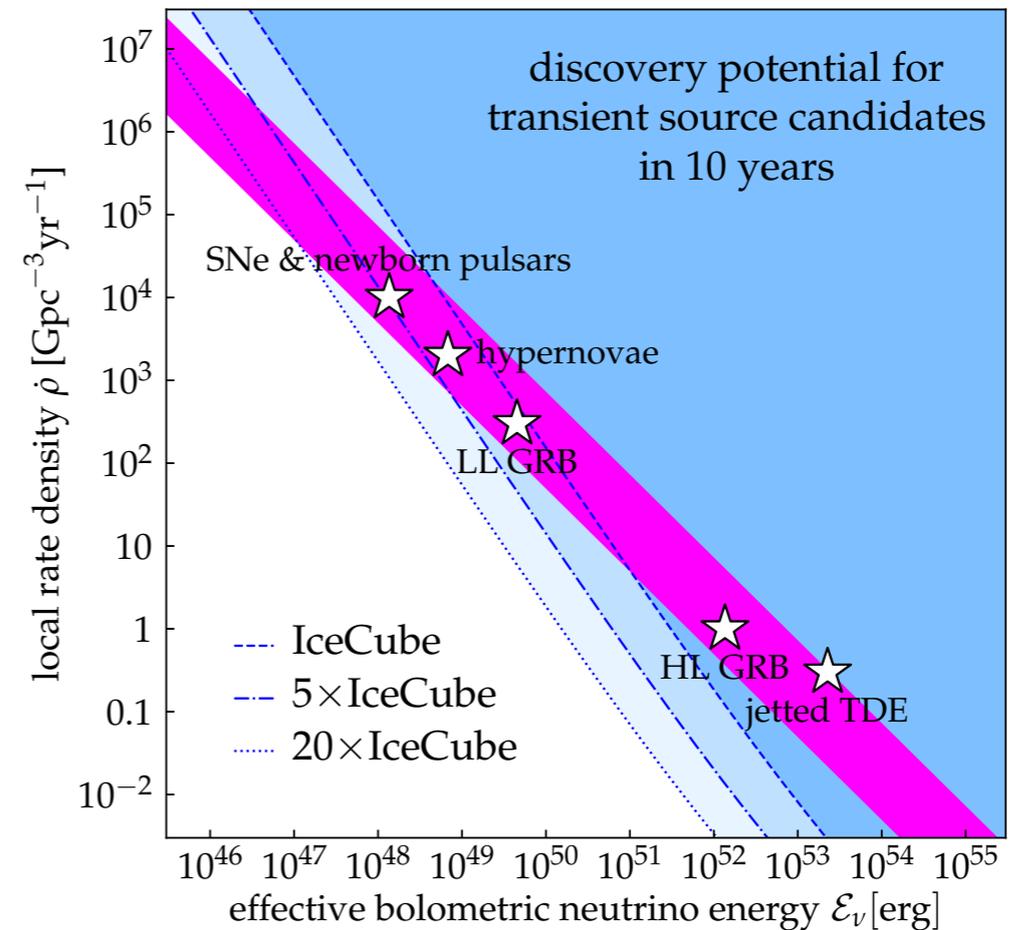
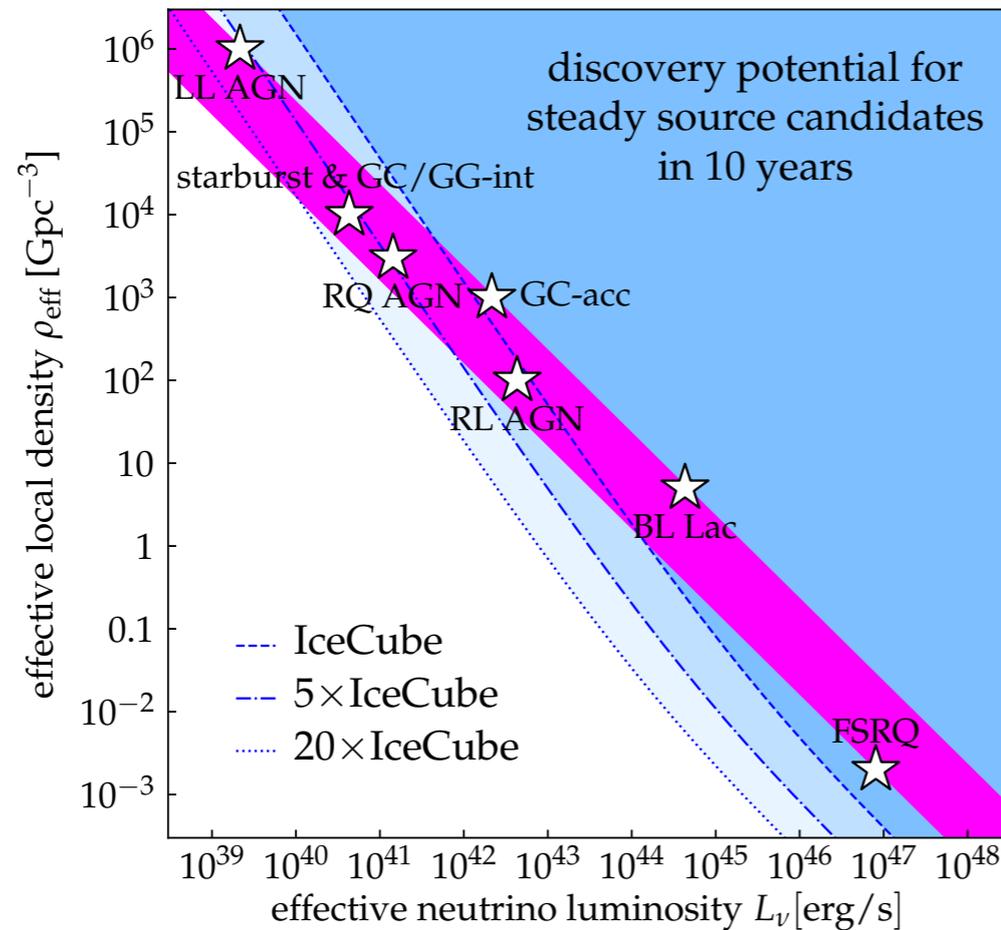
Neutrino spectral shape is similar to that of protons

- Photomeson production ( $p\gamma$ )



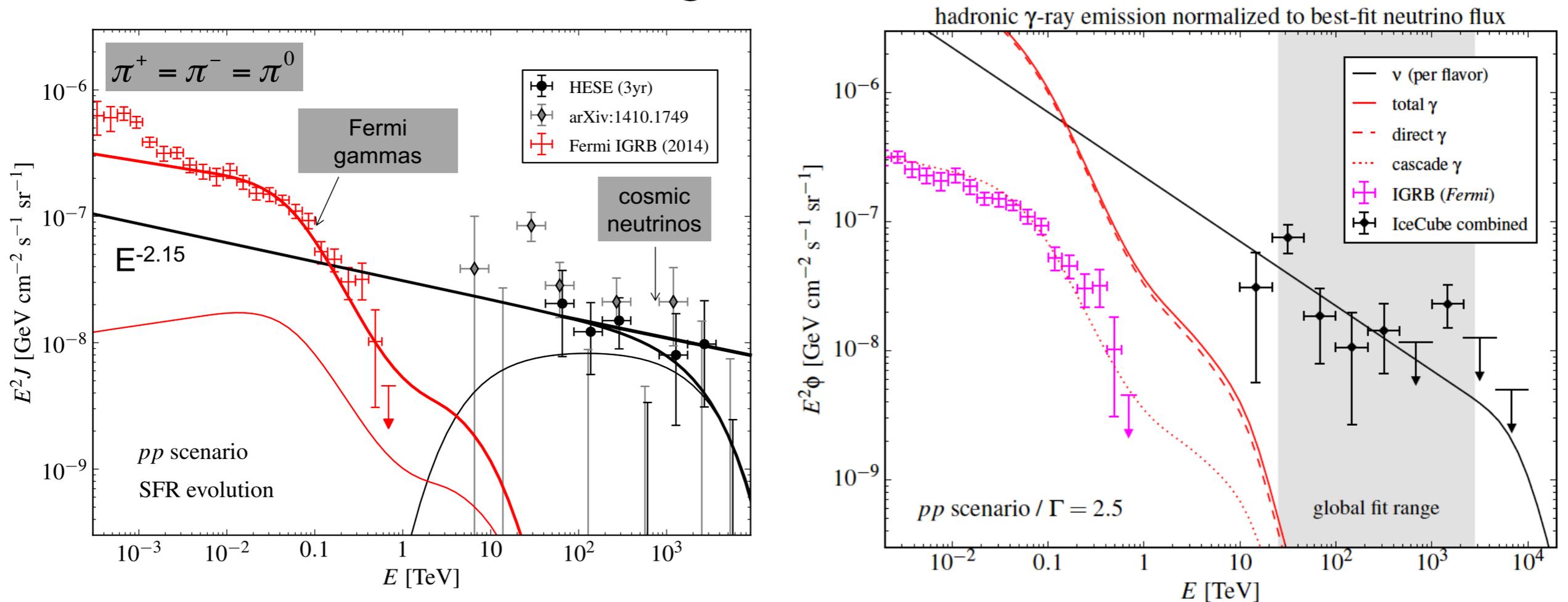
# Point Source Constraint

see Murase & Waxman 16



- No point-source detection  
→ High number density of neutrino sources
- IceCube already disfavors luminous sources (GRBs, Blazars, Jetted TDEs)

# Gamma-ray Constraints



- Astrophysical  $\nu$ s are accompanied with  $\gamma$  rays
- $\nu$  intensity at 10 TeV  $>$   $\gamma$ -ray intensity at 100 GeV
  - accompanying  $\gamma$  rays overshoot Fermi data
  - $\nu$  sources should be opaque to  $\gamma$  rays

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# Active Galactic Nuclei (AGNs)

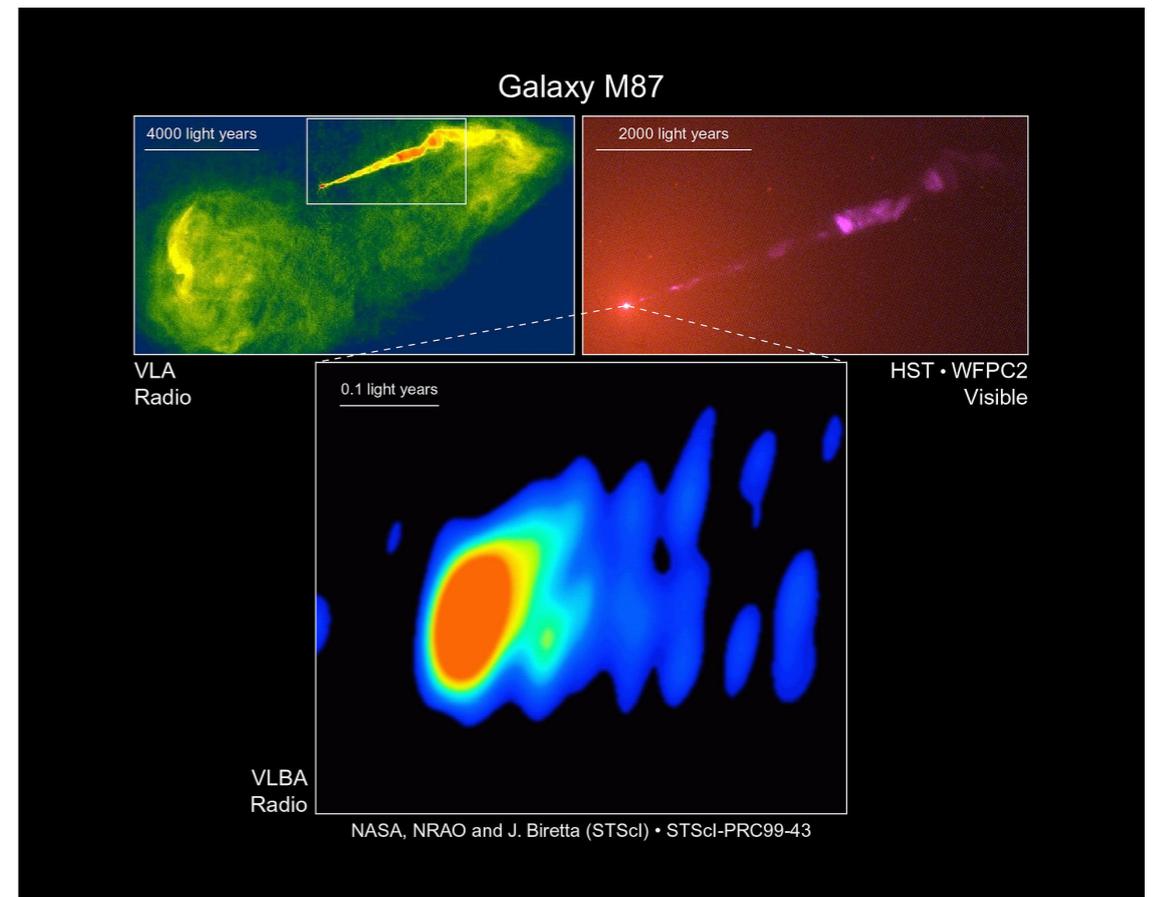
- Radio-quiet AGNs



M77 (NGC 1068): Wikipedia©

- No prominent jet
- 90% of AGNs

- Radio-loud AGNs

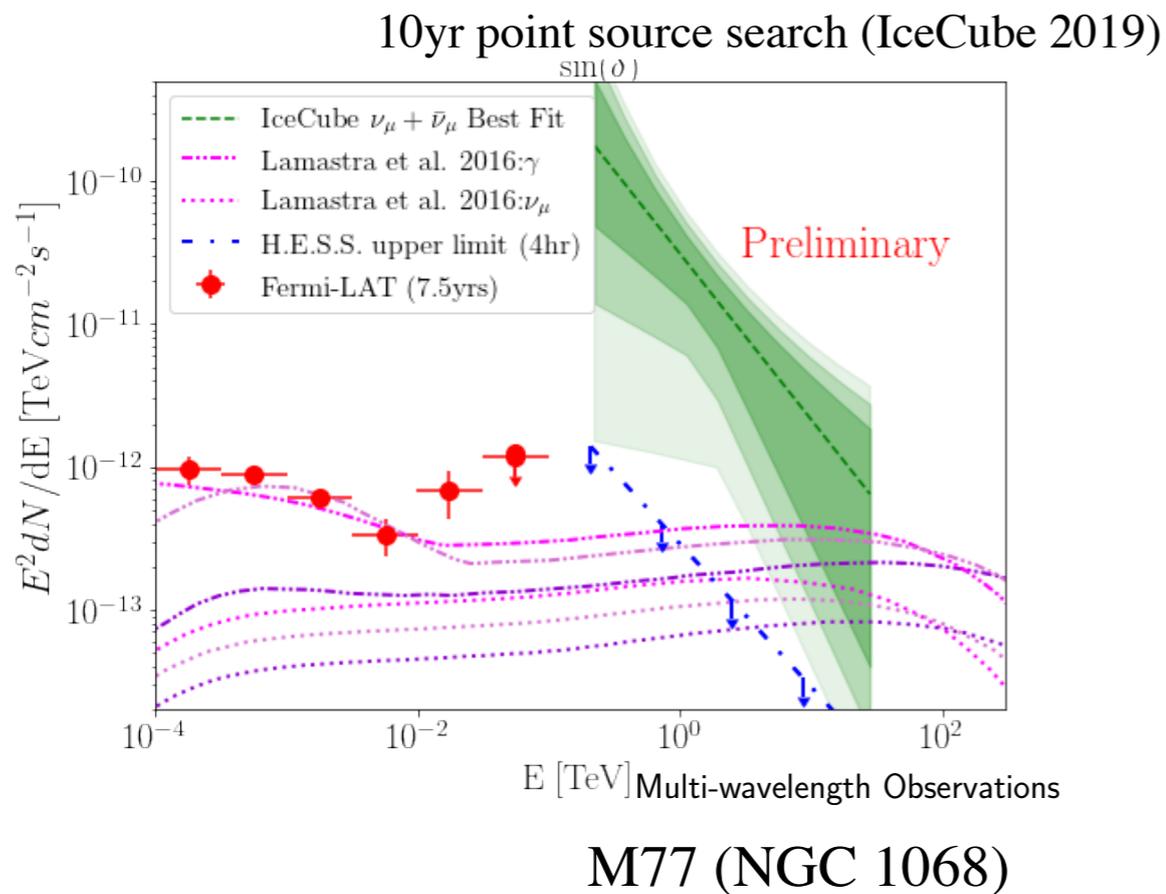


M87 (NGC 4486): Wikipedia©

- Powerful Jets
- 1-10 % of AGNs

# Active Galactic Nuclei (AGNs)

- Radio-quiet AGNs

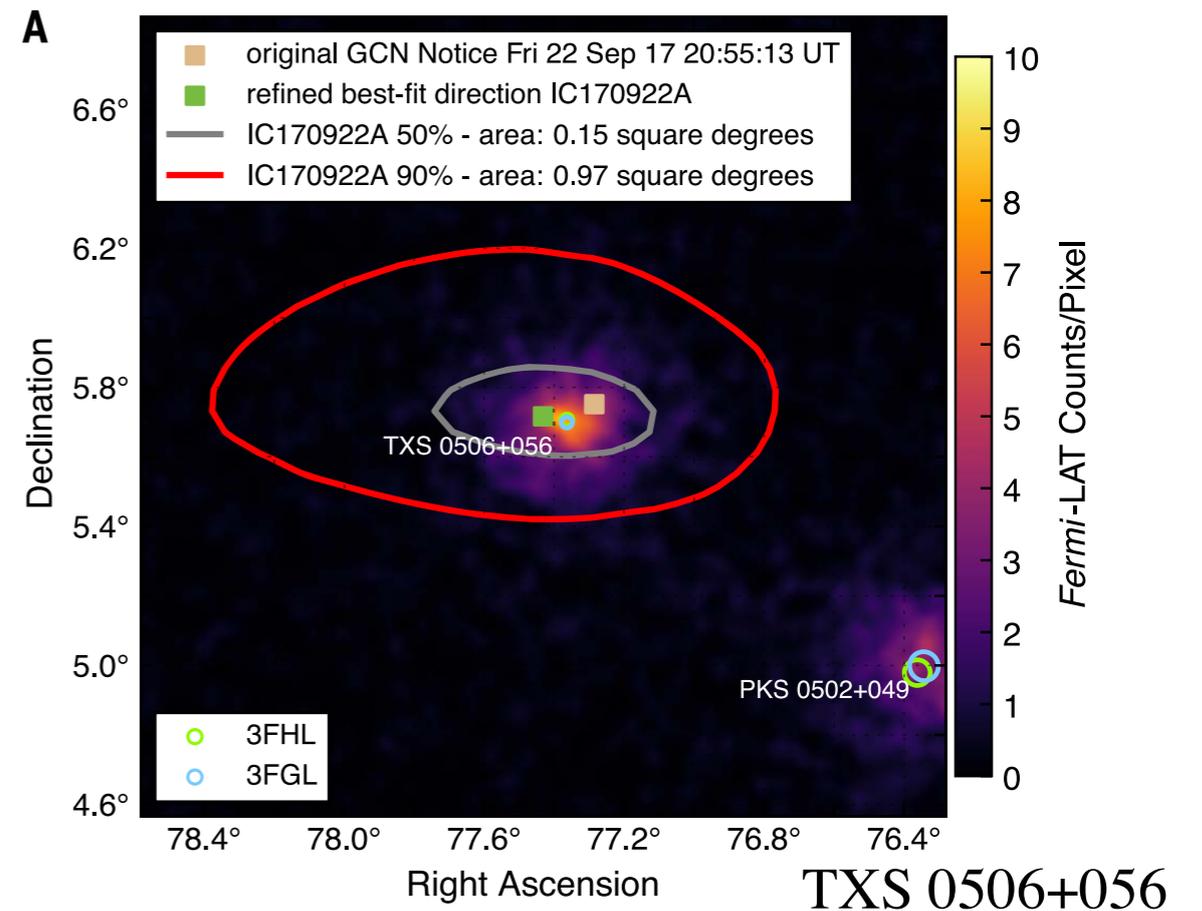


- Hottest Point in Northern Sky ( $2.9 \sigma$ )

- Radio-loud AGNs

Multi-messenger campaign

IceCube 2018

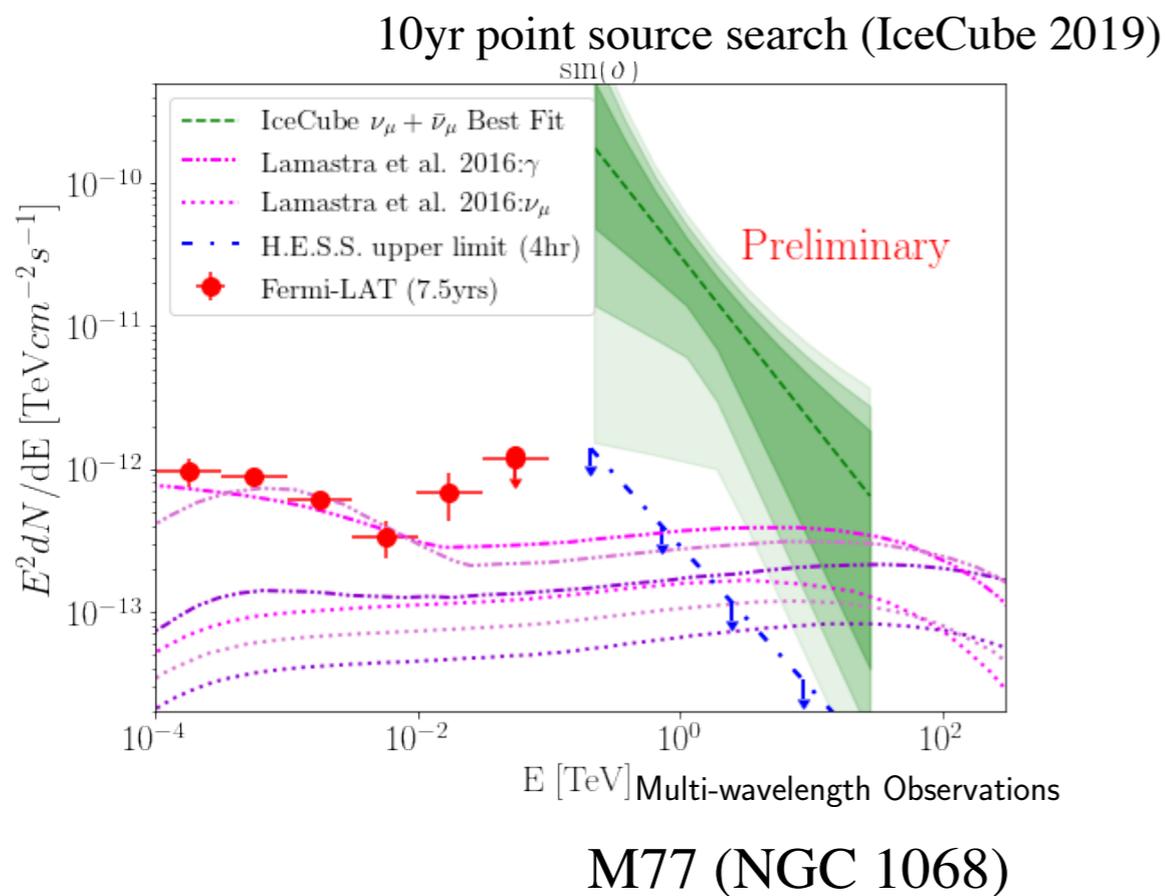


- IC 170922 ( $3 \sigma$ )

- 2014-2015 Neutrino flare ( $3.5 \sigma$ )

# Active Galactic Nuclei (AGNs)

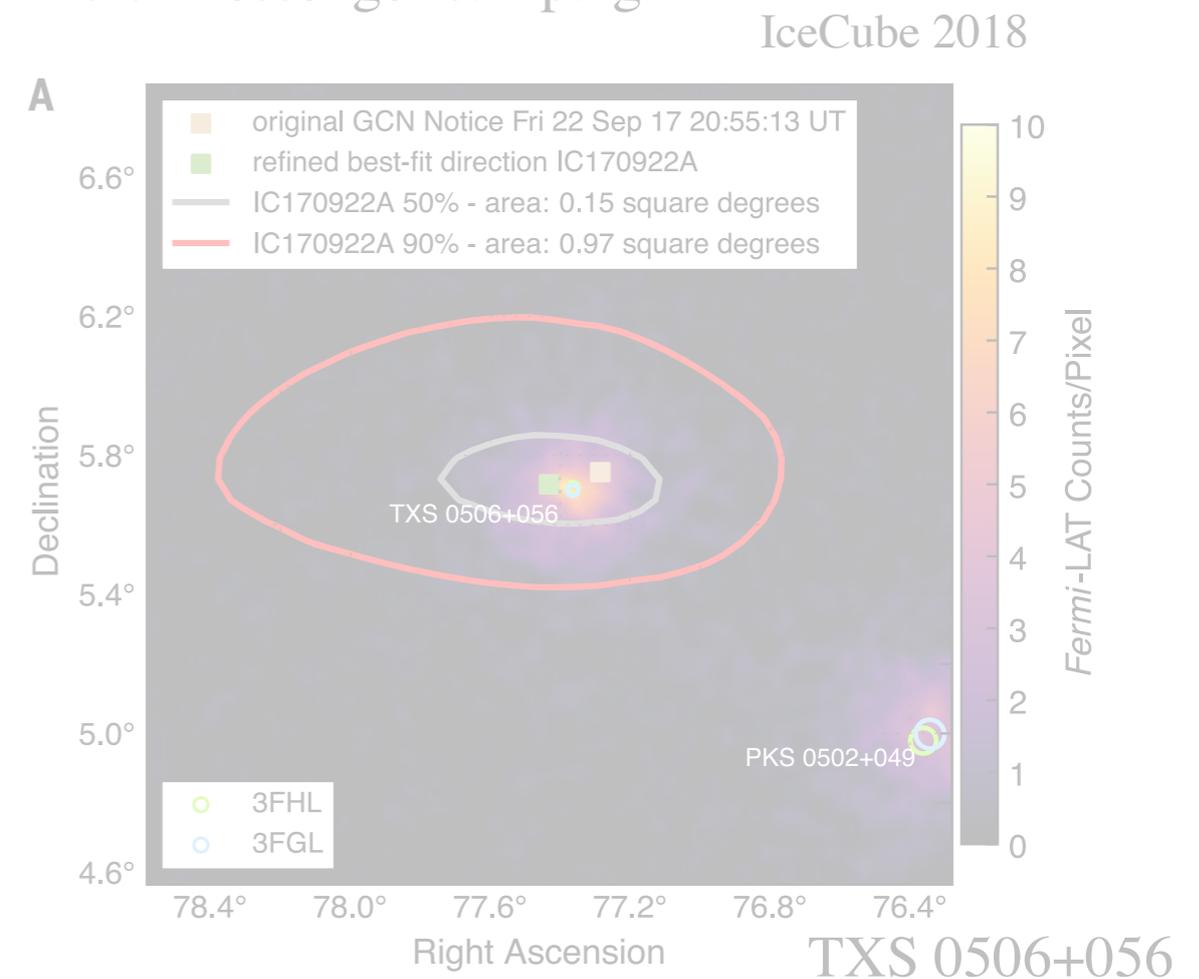
- Radio-quiet AGNs



- Hottest Point in Northern Sky ( $2.9 \sigma$ )

- Radio-loud AGNs

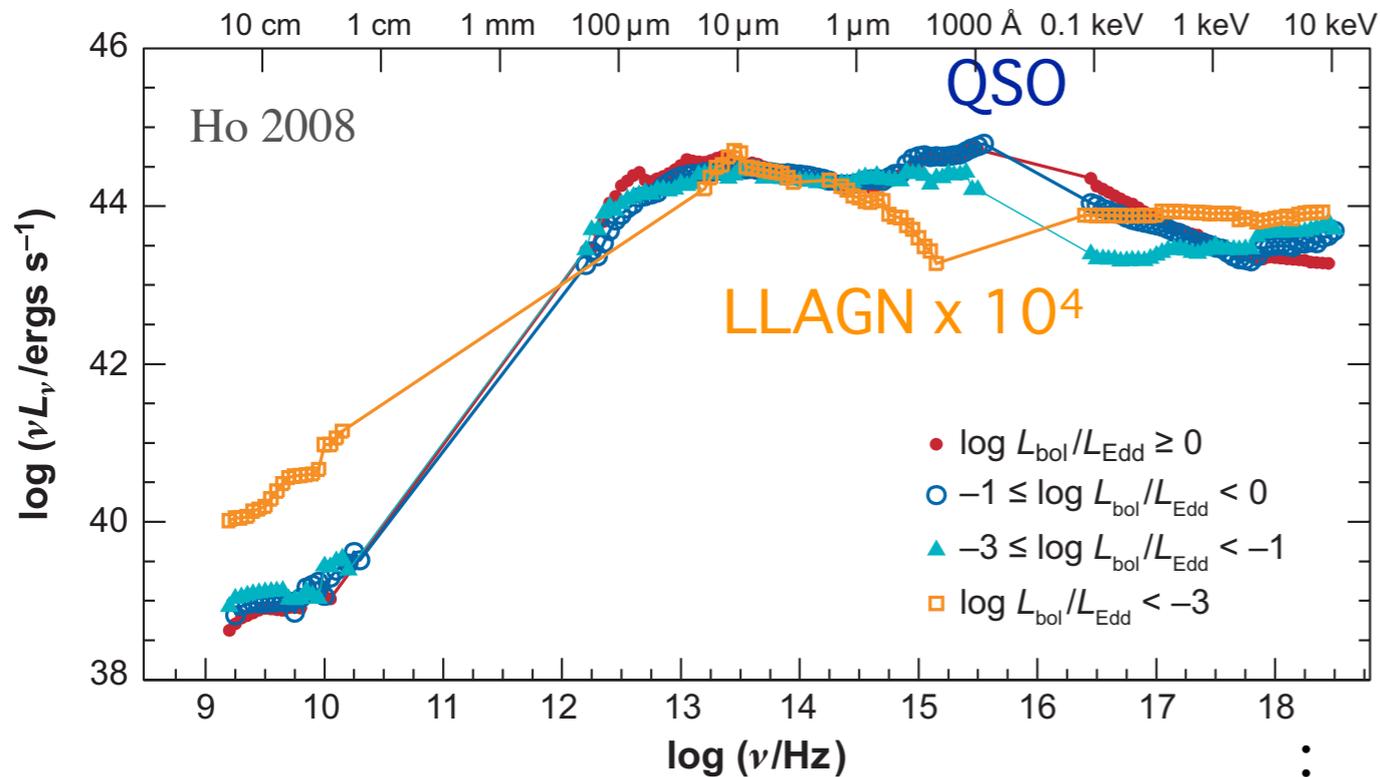
Multi-messenger campaign



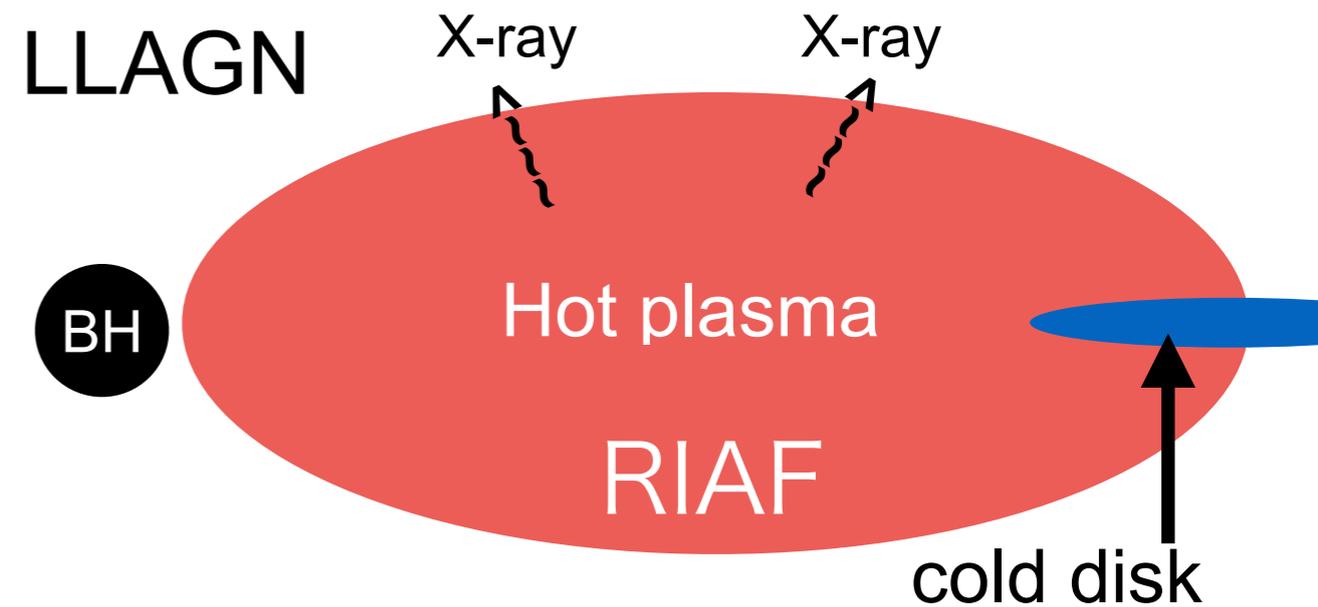
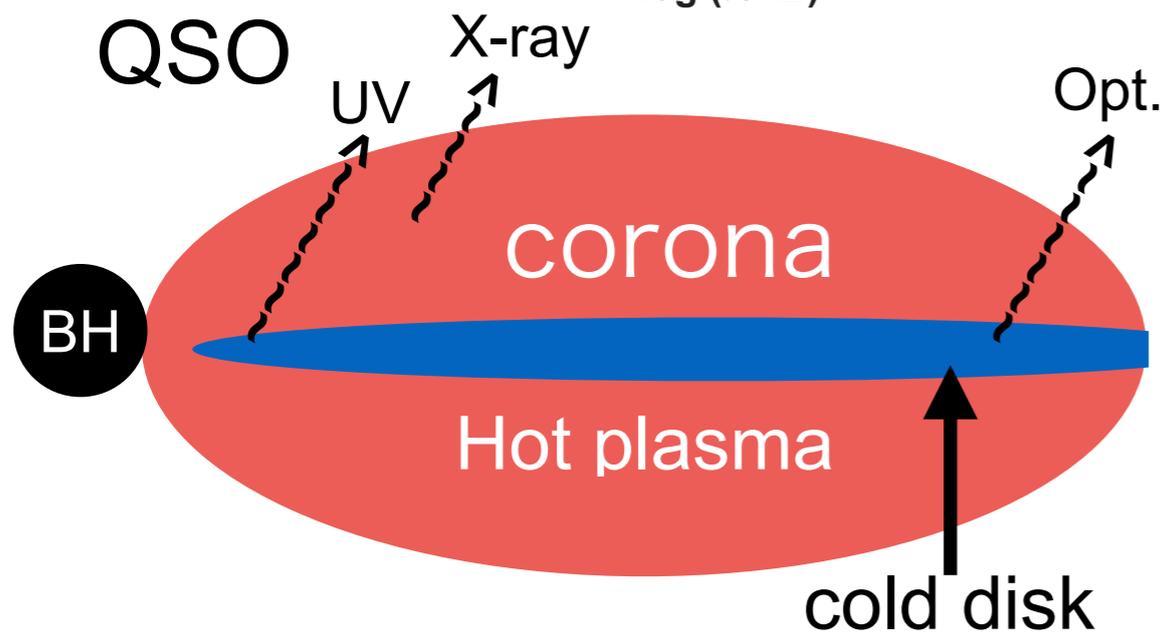
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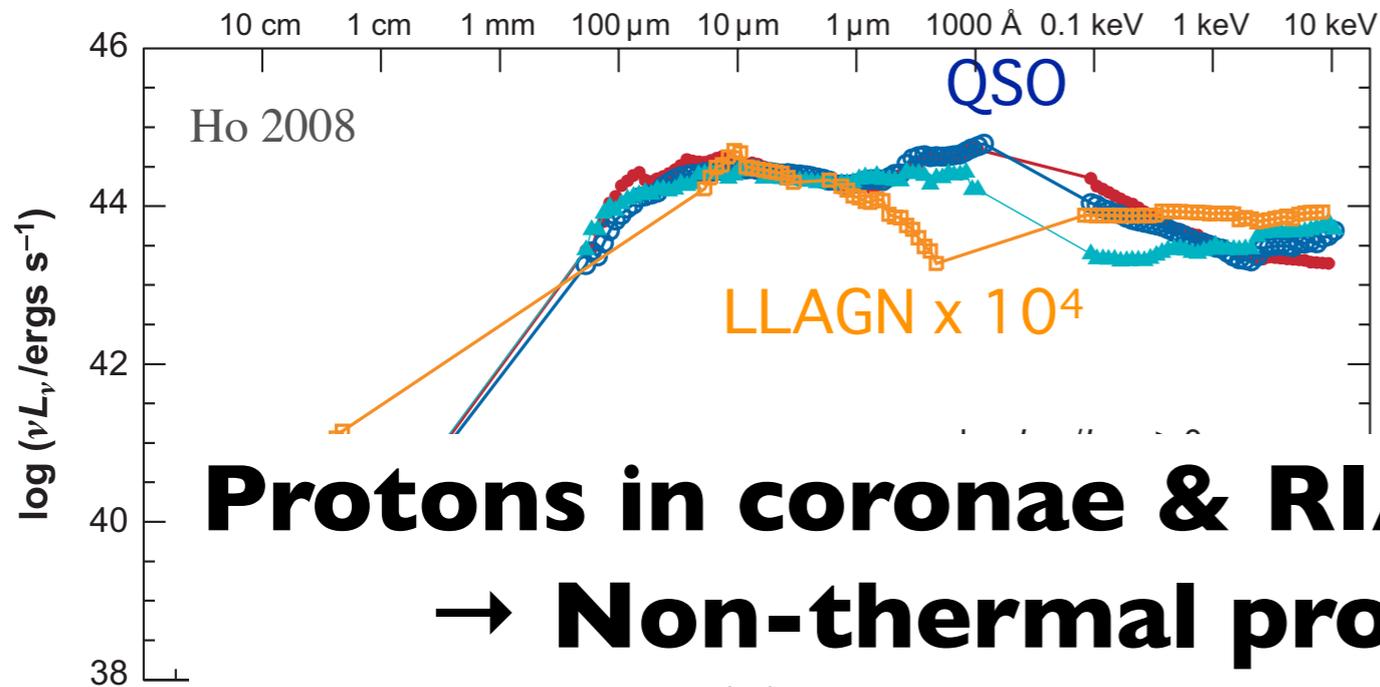
# Radio-quiet AGNs



- QSO: Blue bump & X-ray  
→ Optically thick disk + coroneae
- LLAGN: No blue bump & X-ray  
→ Optically thin flow  
[Radiatively Inefficient Accretion Flow (RIAF)]



# Radio-quiet AGNs

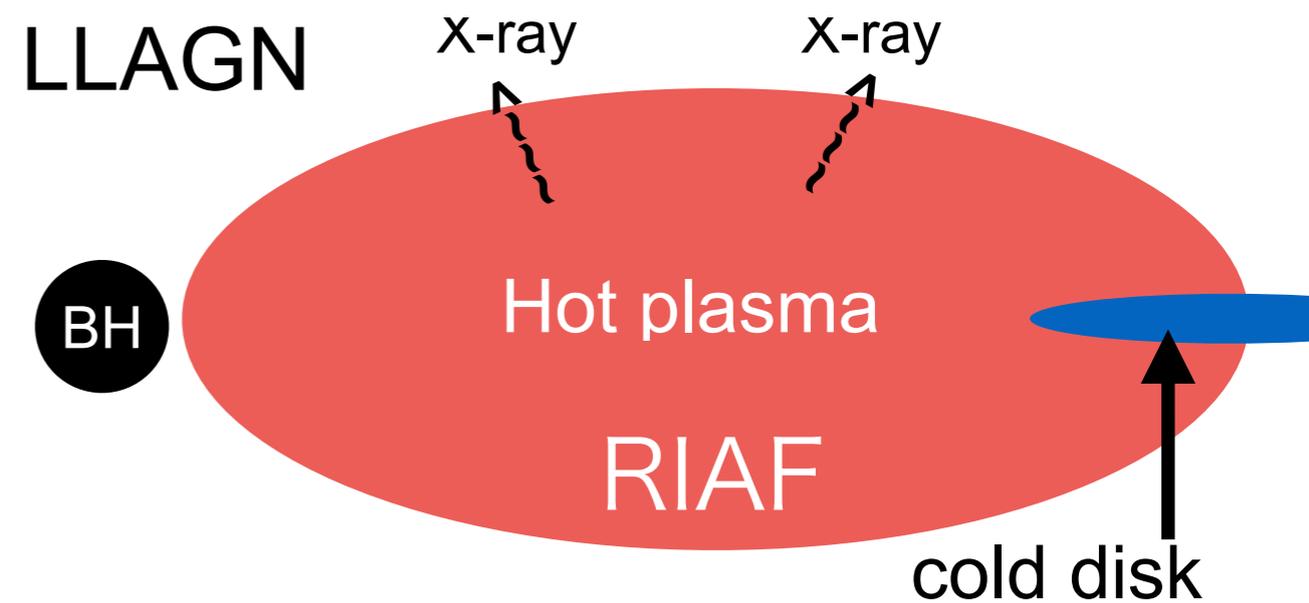
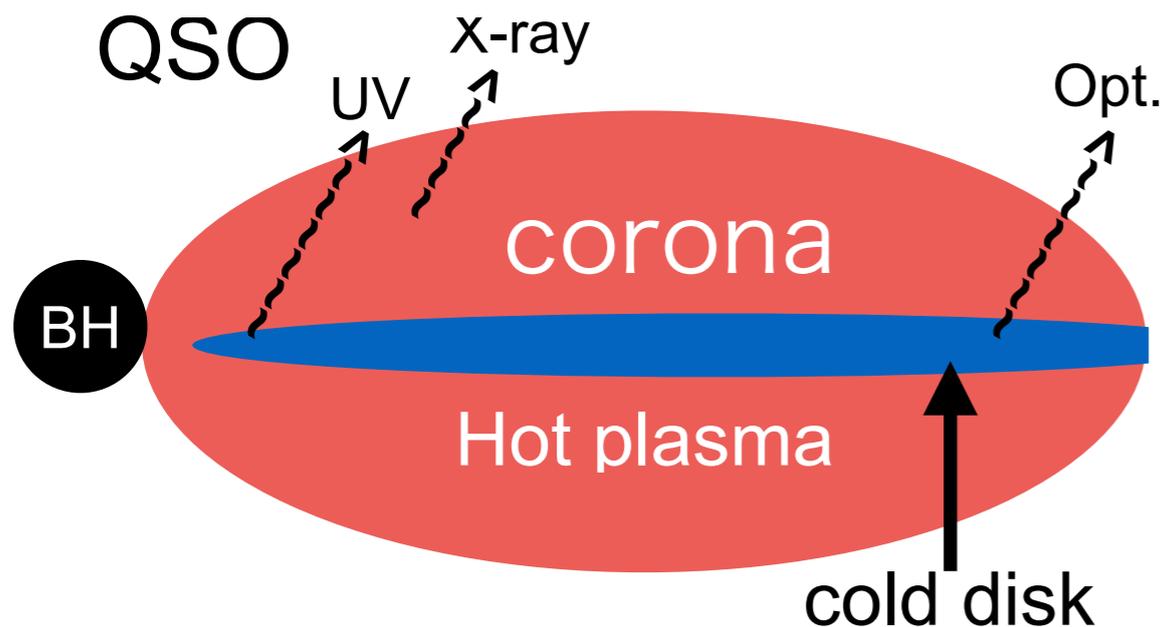


- QSO: Blue bump & X-ray  
→ Optically thick disk + coroneae
- LLAGN: No blue bump & X-ray  
→ Optically thin flow

**Protons in coroneae & RIAFs are collisionless**

→ **Non-thermal proton production**

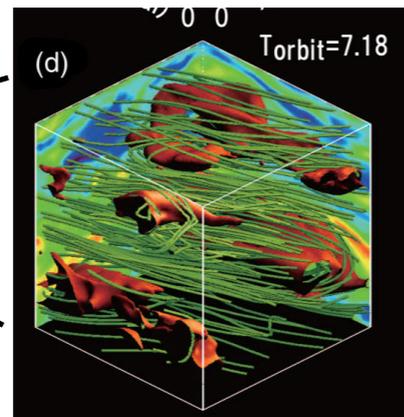
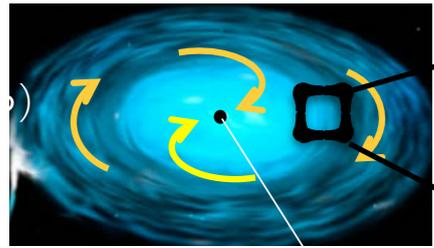
\*\* electrons are collisional \*\*



# Particle Acceleration in Accretion Flows

Particle-In-Cell Simulations

MHD + Test Particle Simulations

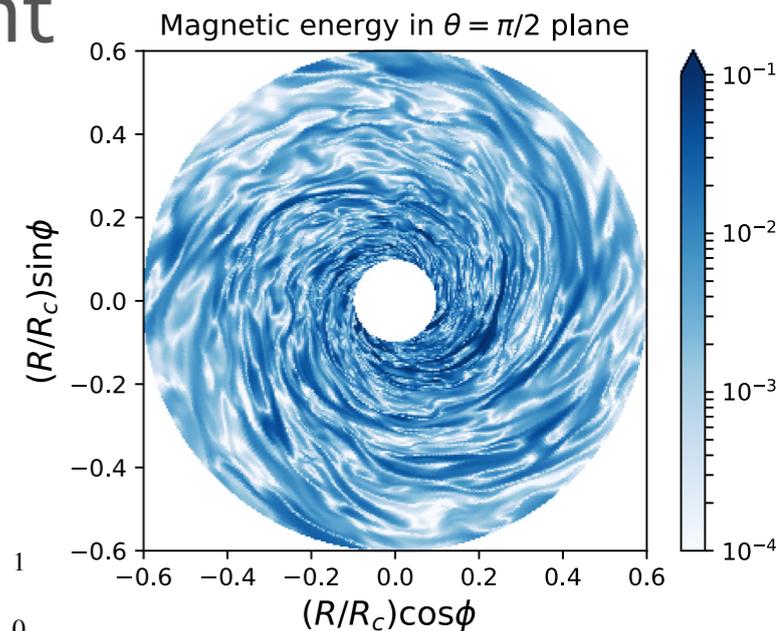


Numerical Simulations

turbulent

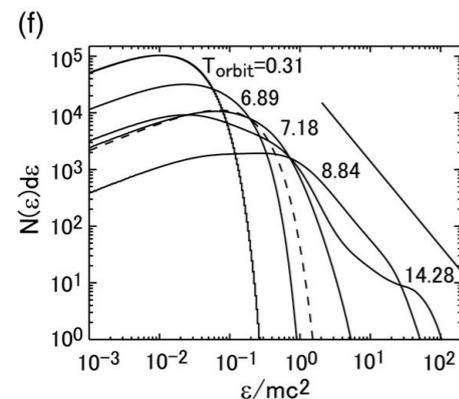
SSK et al. 2016; SSK et al. 2019

see also Lynn et al. 2014

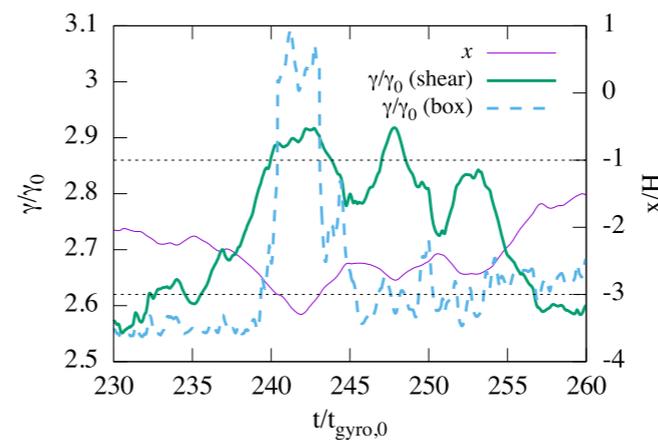


turbulent

Hoshino 2013; Hoshino 2015  
Riquelme et al. 2012; Kuntz et al. 2016



Non-thermal



Energy change

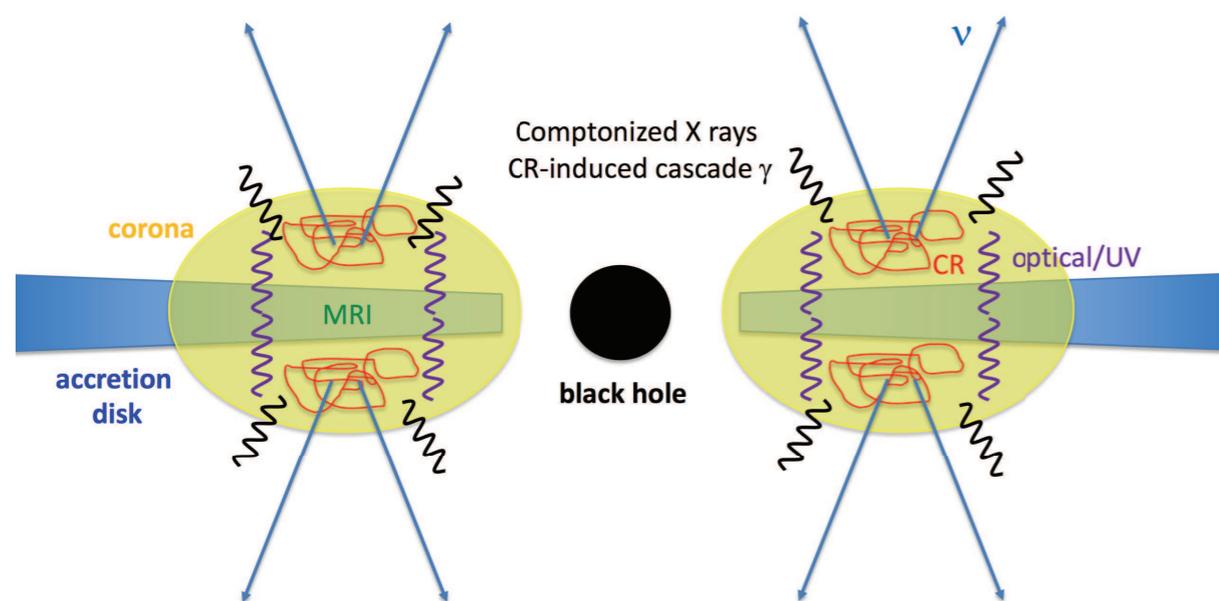
Magnetic reconnection or wave-particle interaction accelerates CRs,

# High-Energy Emission

Mahadevan et al. 1997; SSK et al. 2015

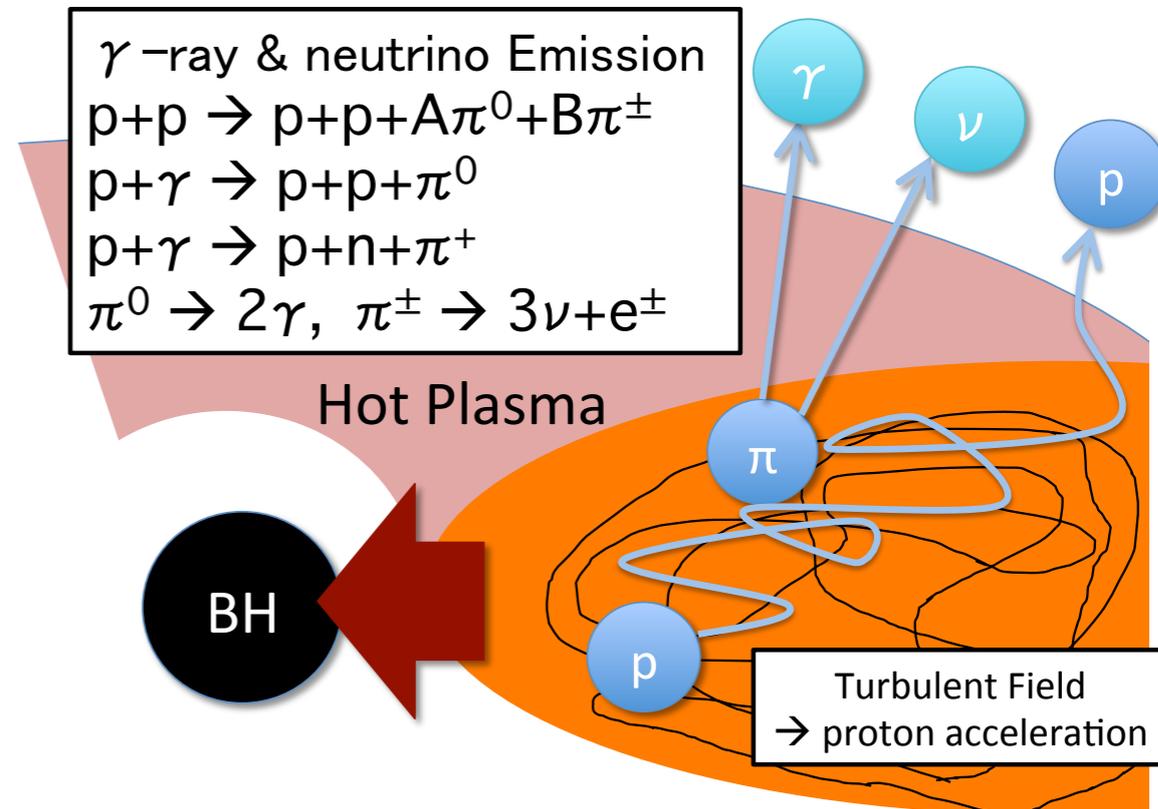
Murase, SSK et al. 2019

## Coronae in QSOs



## RIAFs in LLAGNs

$\gamma$ -ray & neutrino Emission  
 $p+p \rightarrow p+p+A\pi^0+B\pi^\pm$   
 $p+\gamma \rightarrow p+p+\pi^0$   
 $p+\gamma \rightarrow p+n+\pi^+$   
 $\pi^0 \rightarrow 2\gamma, \pi^\pm \rightarrow 3\nu+e^\pm$



- Fairly high photon & proton densities in Coronae & RIAFs
- Interaction between CRs and matter/photons  
 $\rightarrow$  neutrino & gamma-ray emission
- TeV-PeV  $\gamma$ -rays are reprocessed to MeV-GeV  $\gamma$  rays

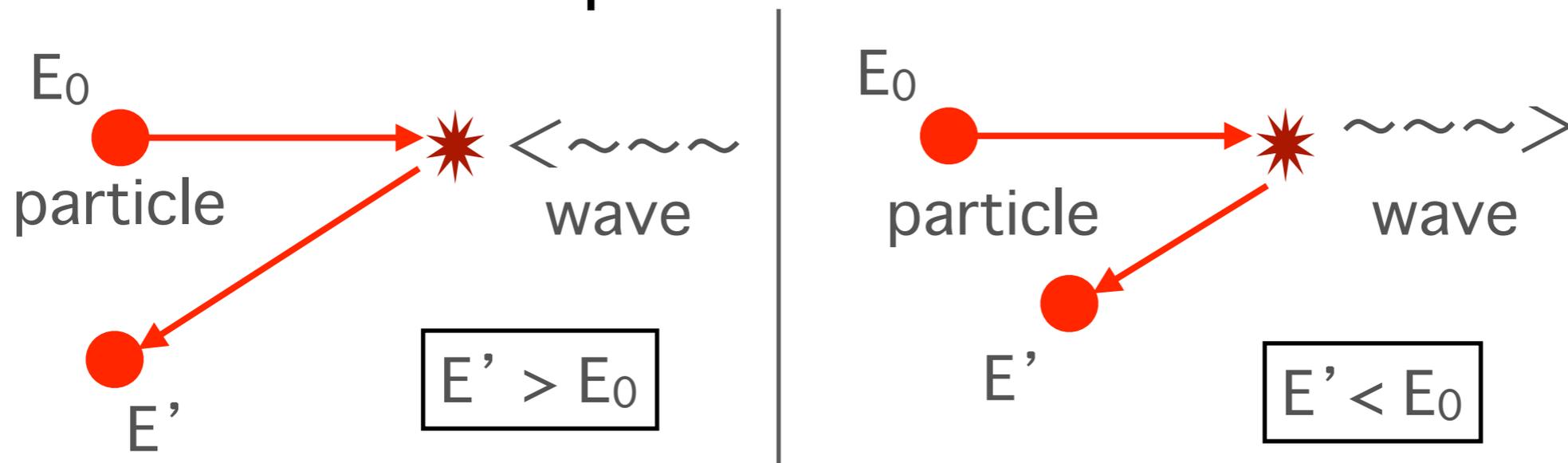
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# Stochastic Acceleration

e.g.) Fermi 1949, Stawarz & Petrosian 2008, SSK et al. 2015

- Consider plasma with turbulent fields



Some gain  $E$ , others lose  $E \rightarrow$  diffusion in  $E$  space

$$\frac{\partial f}{\partial t} = \frac{1}{p^2} \frac{\partial}{\partial p} \left( p^2 D_p \frac{\partial f}{\partial p} \right)$$

According to quasi-linear theory,  
gyro-resonant scattering results in

$$D_p \propto p^q \quad (\text{power-spectrum } P_k = P_0 k^{-q})$$

# Basic Equations

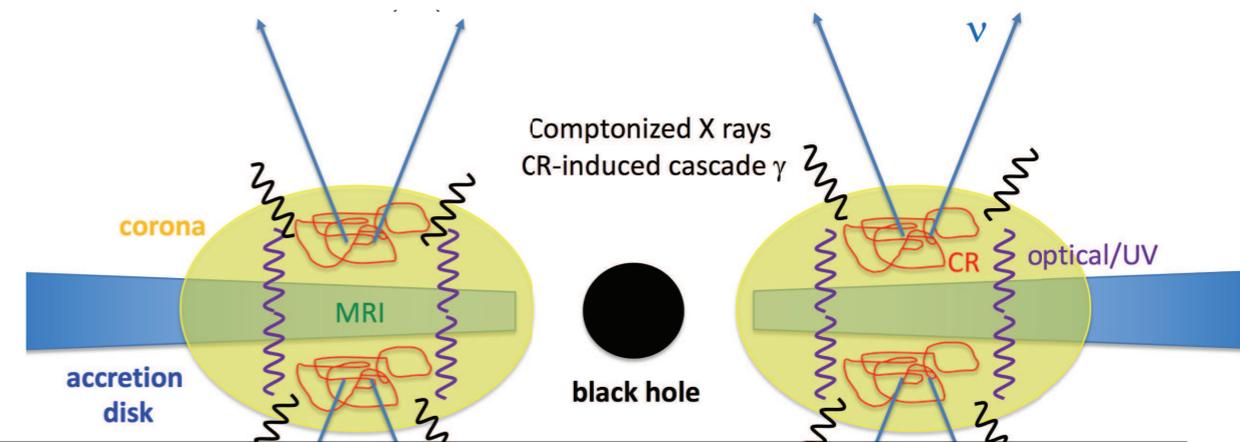
- Stochastic Acceleration (SA)

cf.) Dermer et al. 2012, SSK et al. 2015

$$\frac{\partial F_p}{\partial t} = \frac{1}{\varepsilon_p^2} \frac{\partial}{\partial \varepsilon_p} \left( \varepsilon_p^2 D_{\varepsilon_p} \frac{\partial F_p}{\partial \varepsilon_p} + \frac{\varepsilon_p^3}{t_{p\text{-cool}}} F_p \right) - \frac{F_p}{t_{\text{esc}}} + \dot{F}_{p,\text{inj}}$$

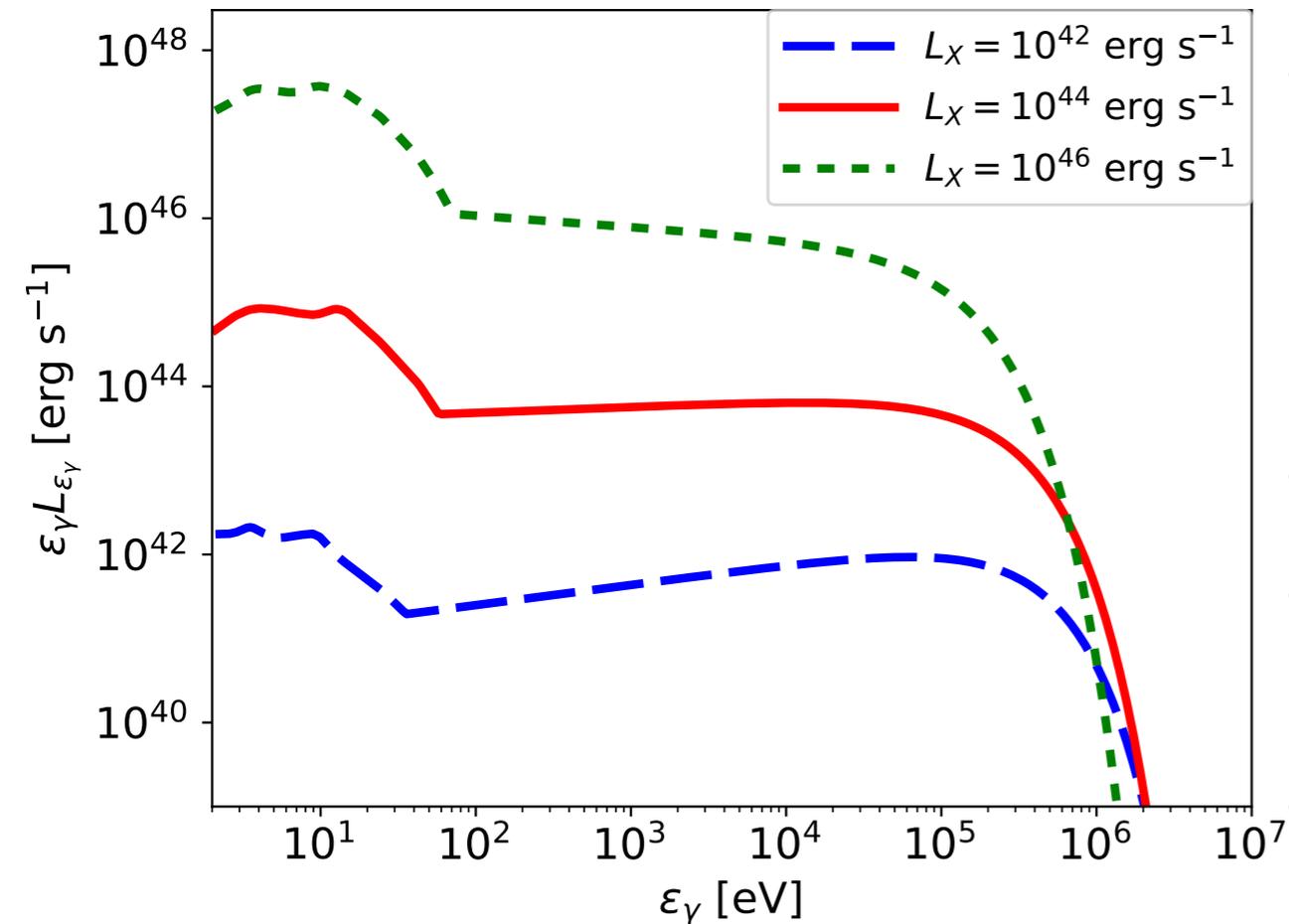
$$D_{\varepsilon_p} \approx \frac{\zeta c}{H} \left( \frac{V_A}{c} \right)^2 \left( \frac{r_L}{H} \right)^{q-2} \varepsilon_p^2,$$

$$\dot{F}_{p,\text{inj}} = \dot{F}_0 \delta(\varepsilon_p - \varepsilon_{p,\text{inj}})$$



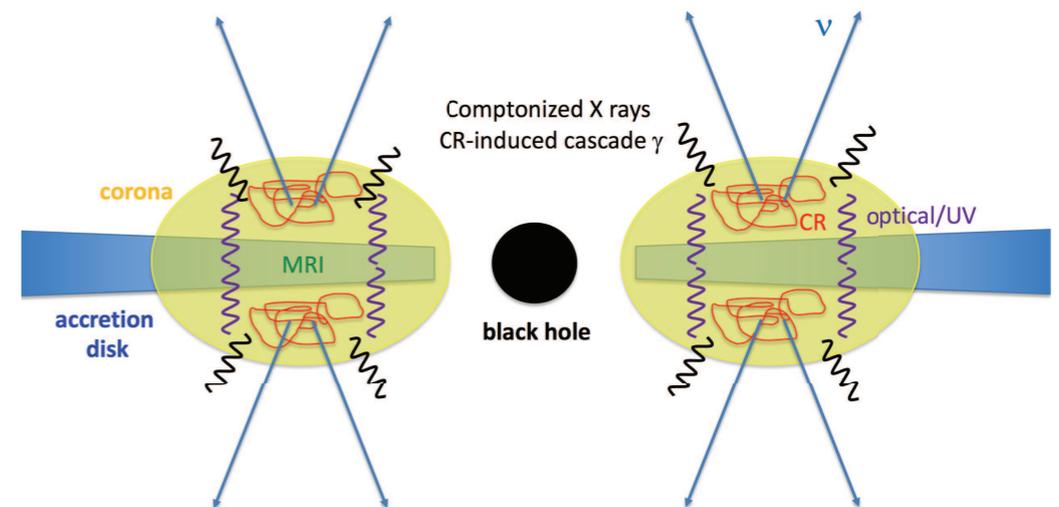
- Escape : Diffusive escape & infall to SMBH
- Coolings: pp inelastic collision, photomeson production  
proton synchrotron, Bethe-Heitler process ( $p+\gamma \rightarrow p+e^++e^-$ )
- Muon & Pion Coolings are negligibly inefficient
- HE  $\gamma$ -rays are absorbed by target photons ( $\gamma+\gamma \rightarrow e^++e^-$ )  
→ electron & positron emit high-energy gamma-rays  
→ Calculate electromagnetic cascades

# Target Photon Field

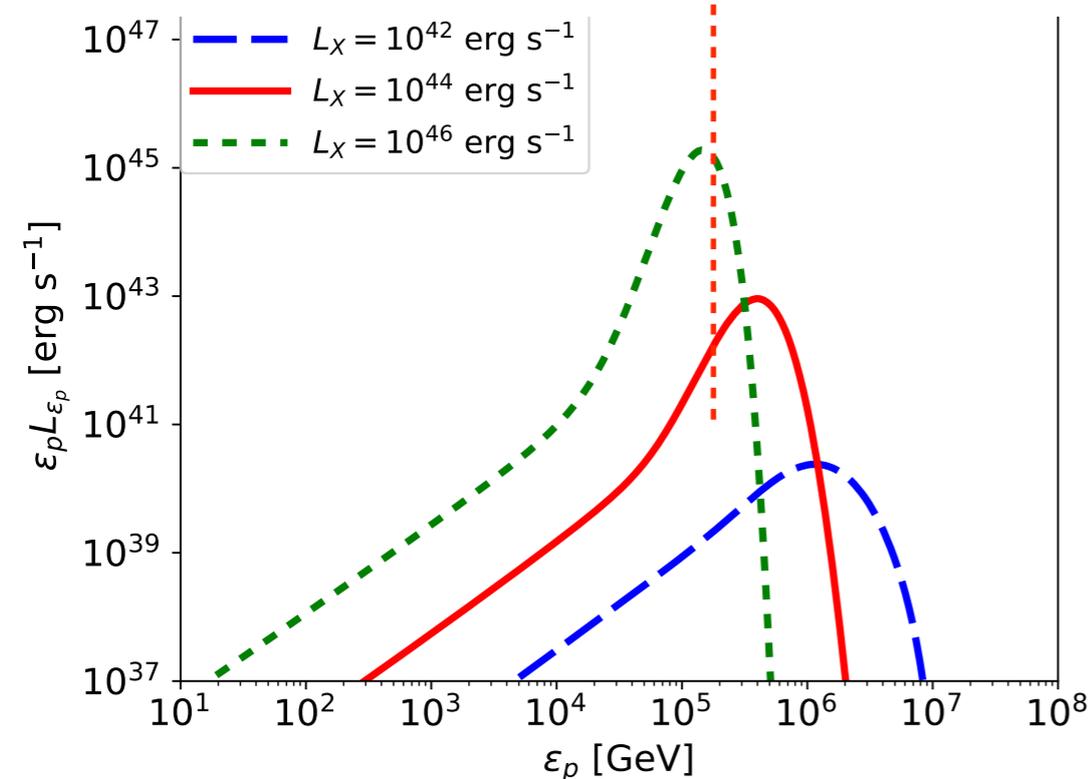
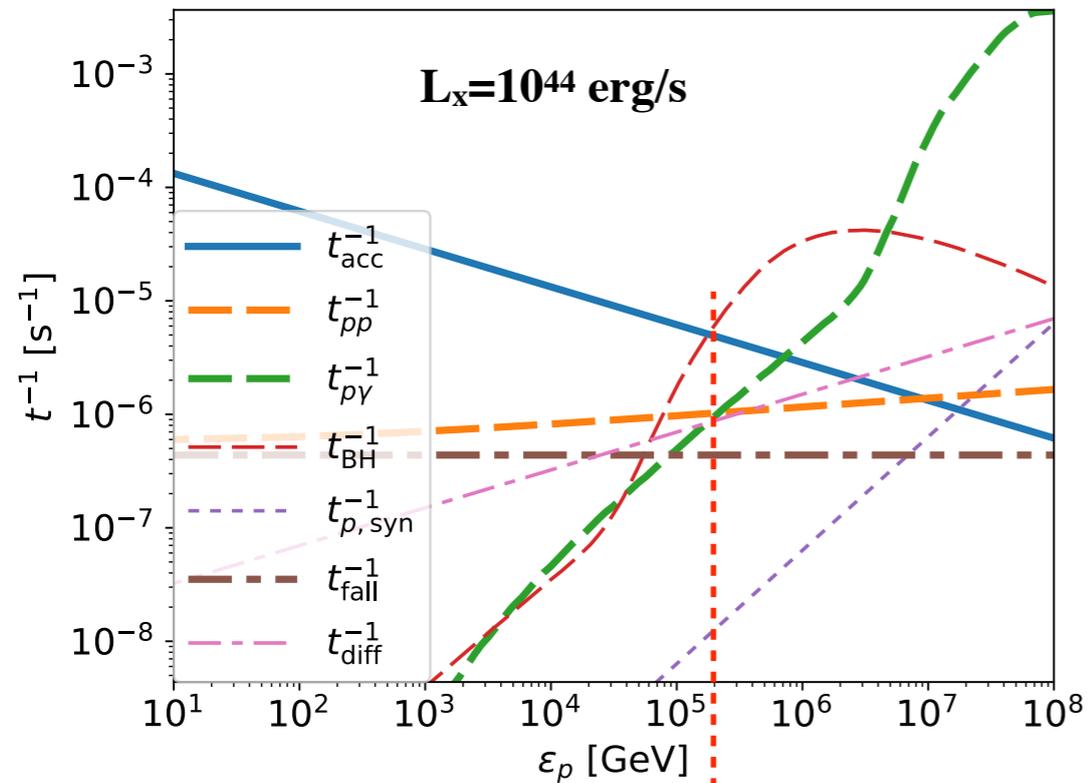


- Luminous objects
  - Rich observational data
  - **We can use empirical relation based on observations**
- Opt-UV photons from accretion disk
- X-rays from hot coronae above thin disk
- Higher  $L_{\text{opt}}/L_X$  for higher  $L_X$  AGNs
- Softer spectra for higher  $L_X$  AGNs

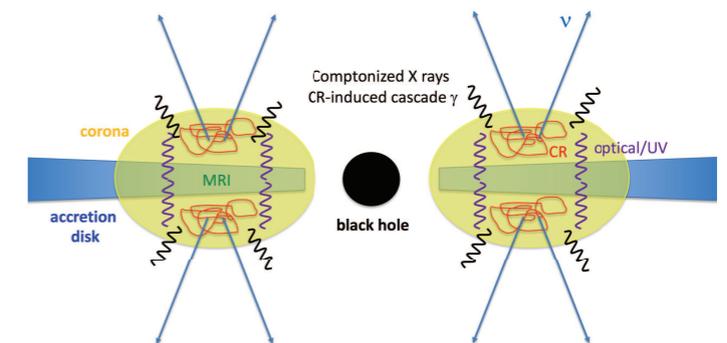
Pringle 1981, Ho 2008, Hopkins 2007  
 Bat AGN Spectroscopic Survey 2017, 2018,  
 Mayers et al. 2018



# Rates & CR Spectrum

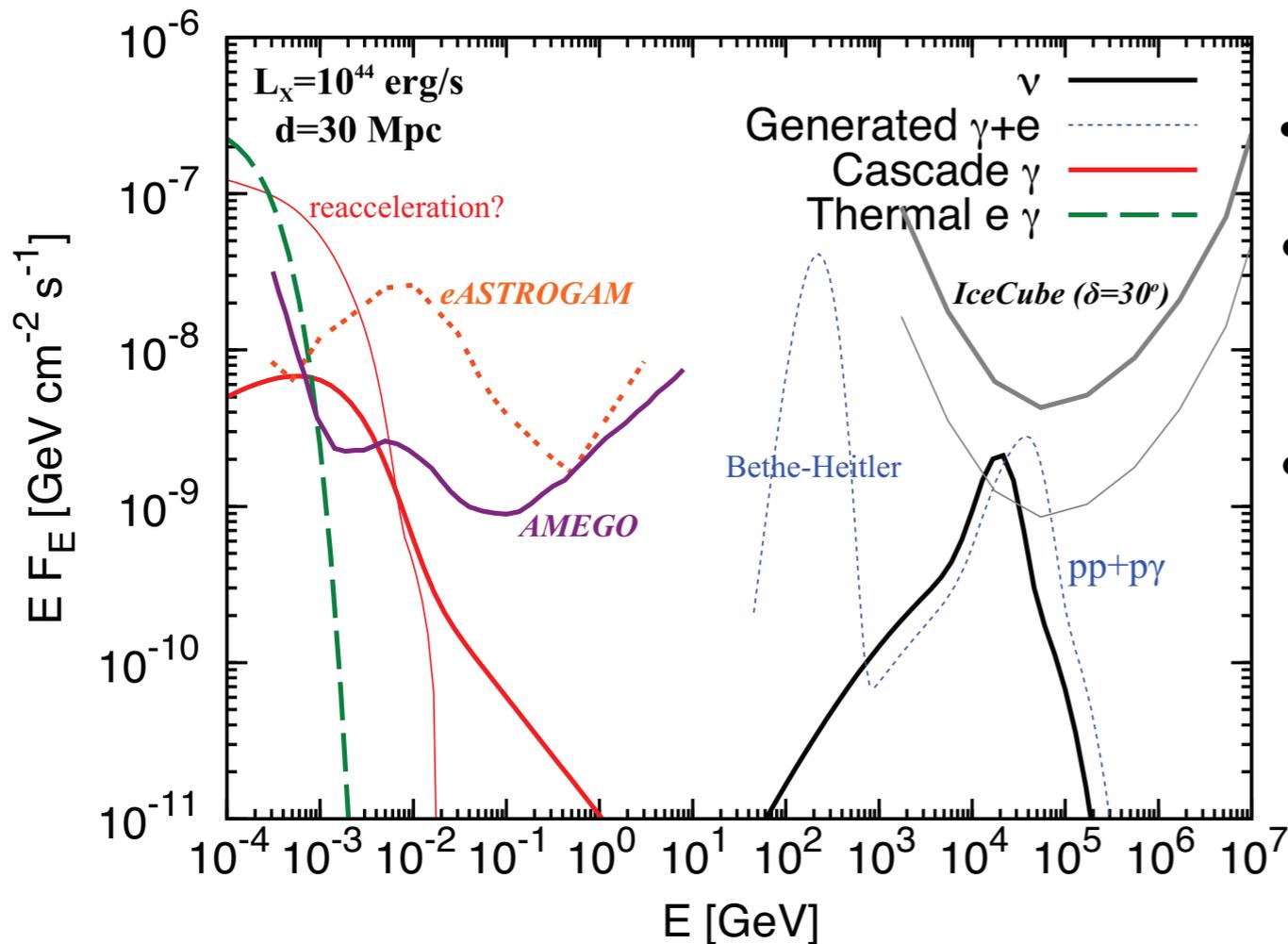


- $E_{p,\text{max}} \sim 10^5$  GeV by  $t_{\text{acc}} = t_{\text{BH}}$
- BH suppresses  $\nu$  production at  $E_p \sim 3 \times 10^4 - 3 \times 10^6$  GeV
- Escape is inefficient

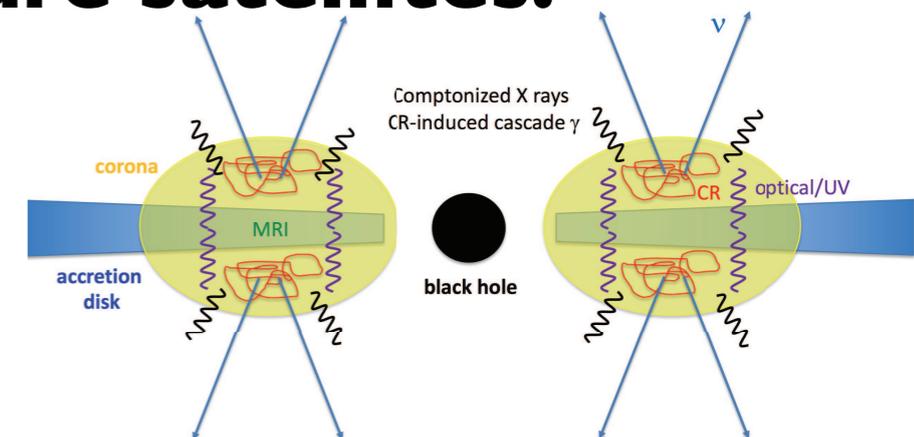


- Hard spectrum due to SA
- Pile up around  $E_{\text{max}}$
- Higher  $L_x \rightarrow$  lower  $E_{\text{max}}$  because of efficient cooling

# HE particles from Nearby Seyfert Galaxies



- A typical Seyfert at 100 Mpc
- **p $\gamma$  neutrinos are detectable by IceCube-Gen2**
- **MeV  $\gamma$ -rays can be detected by future satellites.**

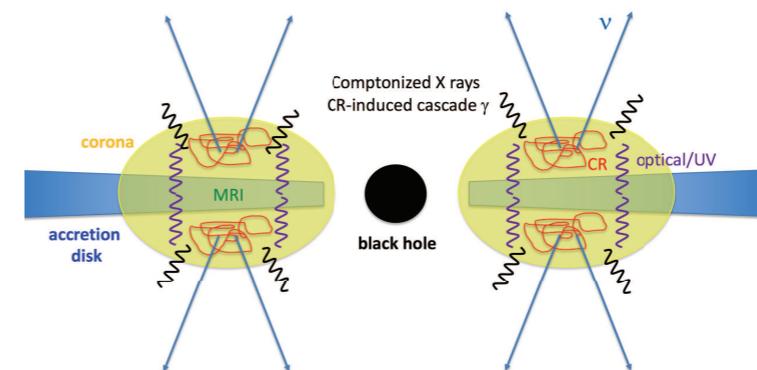
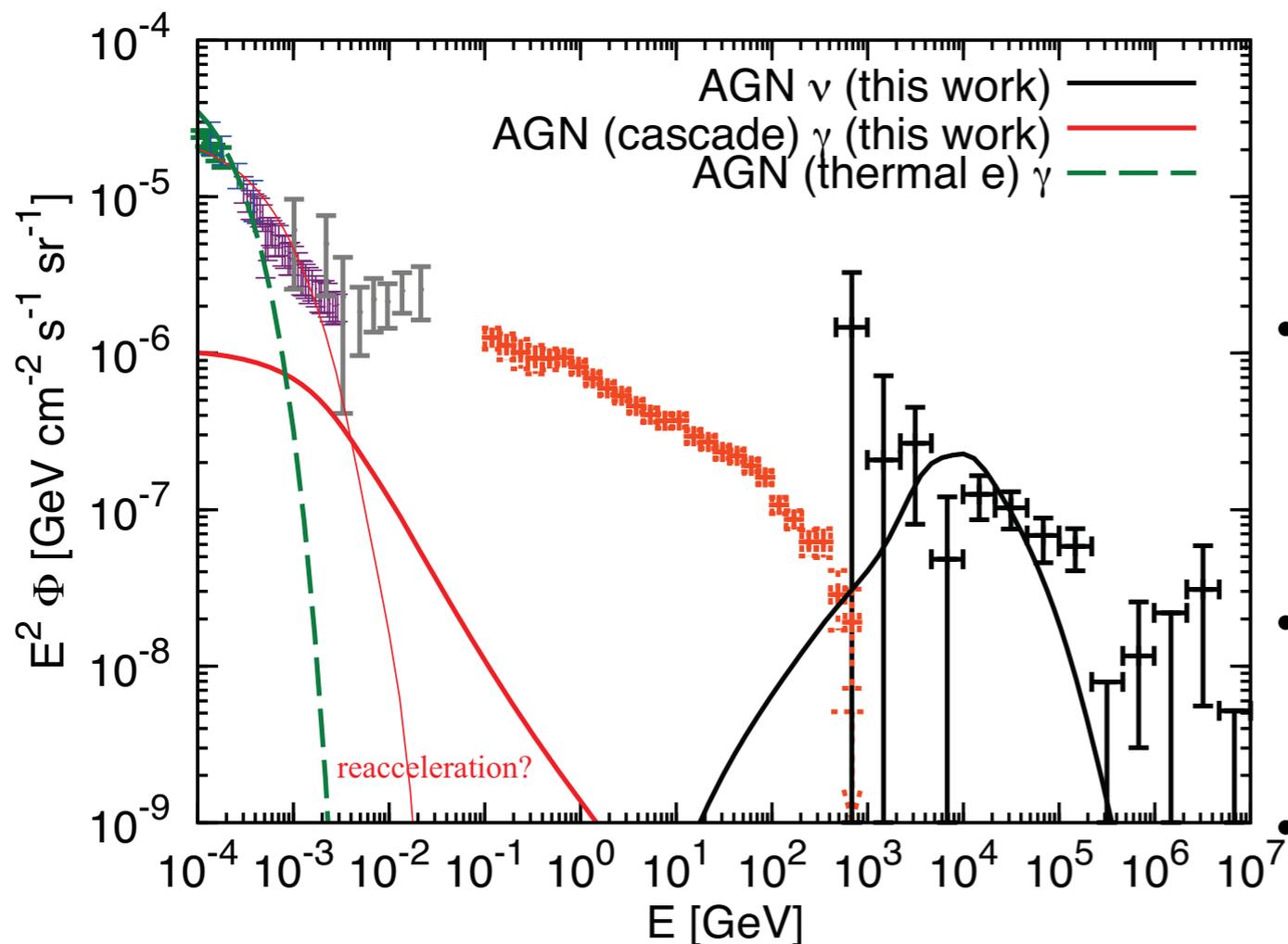


- MeV  $\gamma$ -ray luminosity is determined by B-H pair production  
 → Ratio of  $\gamma$  to  $\nu$  flux is fixed by the observed photon field  
 → We can robustly test our model by future experiments

# Extragalactic $\gamma$ & $\nu$ Backgrounds

$$\Phi_{\nu, \text{ob}}^{\text{diff}}(E_{\nu, \text{ob}}) = \frac{1}{4\pi} \int_{L_{\text{min}}}^{L_{\text{max}}} dL_X \int_0^{z_{\text{max}}} dz \frac{dn_0}{dL_X} f(z) \frac{dV}{dz} \Phi_{\nu, \text{ob}},$$

- AGNs with  $L_X \sim 10^{44}$  erg/s provide the dominant contribution  
e.g., Ueda et al. 2014

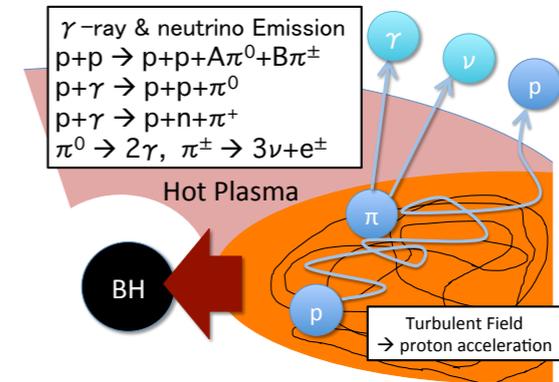


- We choose the injection efficiency so that our model can explain the MESE excess.
- Energetically reasonable:**  
 $P_{\text{CR}}/P_{\text{th}} \sim P_{\text{CR}}/P_{\text{B}} \sim 0.01$
- Cascade emission provides 10 - 30 % of MeV  $\gamma$ -ray background

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- **Stochastic Acceleration (SA)**

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$$\dot{F}_{p,\text{inj}} = \dot{F}_0 \delta(\varepsilon_p - \varepsilon_{p,\text{inj}})$$

- **Power-law Injection (PL)**

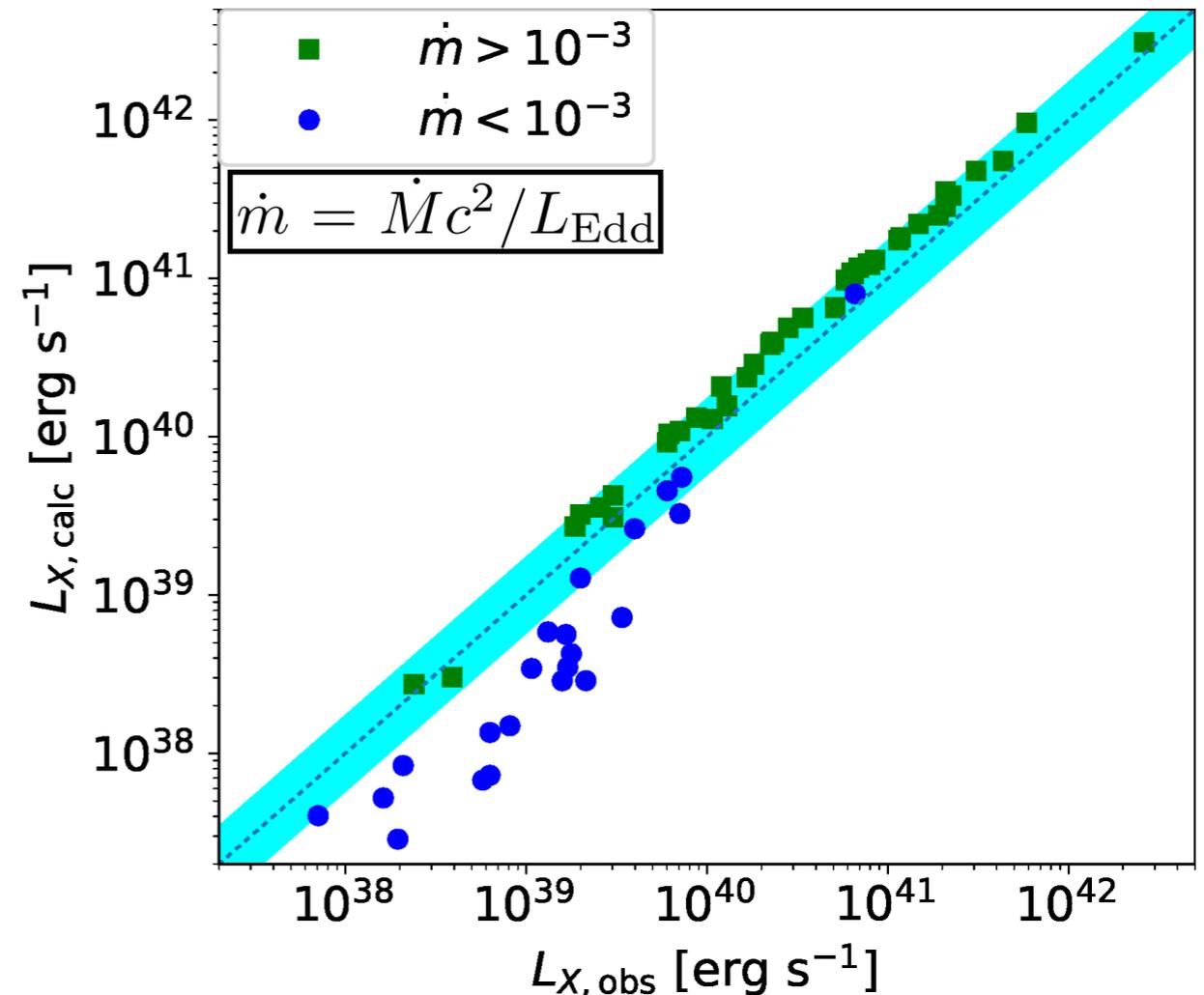
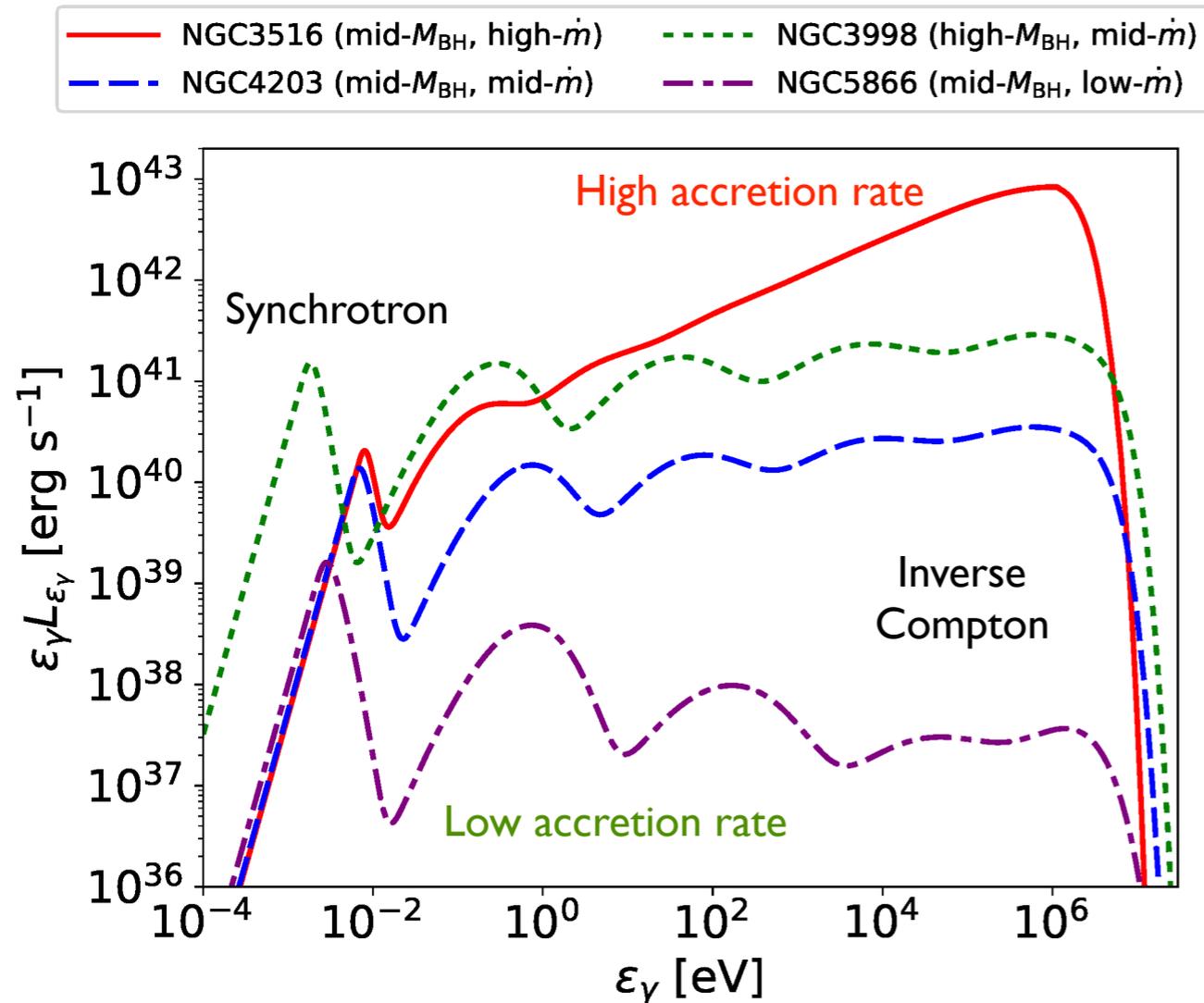
$$\frac{d}{d\varepsilon_p} \left( -\frac{\varepsilon_p}{t_{\text{cool}}} N_{\varepsilon_p} \right) = \dot{N}_{\varepsilon_p,\text{inj}} - \frac{N_{\varepsilon_p}}{t_{\text{esc}}}$$

$$\dot{N}_{\varepsilon_p,\text{inj}} = \dot{N}_0 \left( \frac{\varepsilon_p}{\varepsilon_{p,\text{cut}}} \right)^{-s_{\text{inj}}} \exp \left( -\frac{\varepsilon_p}{\varepsilon_{p,\text{cut}}} \right)$$

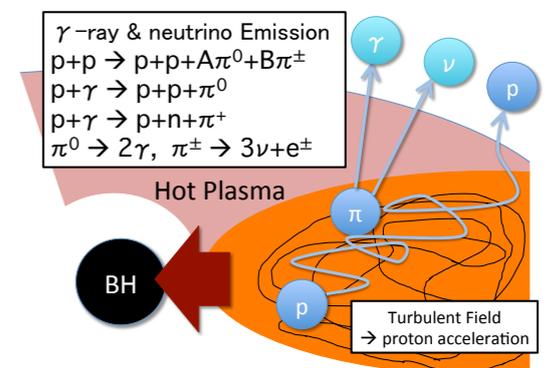
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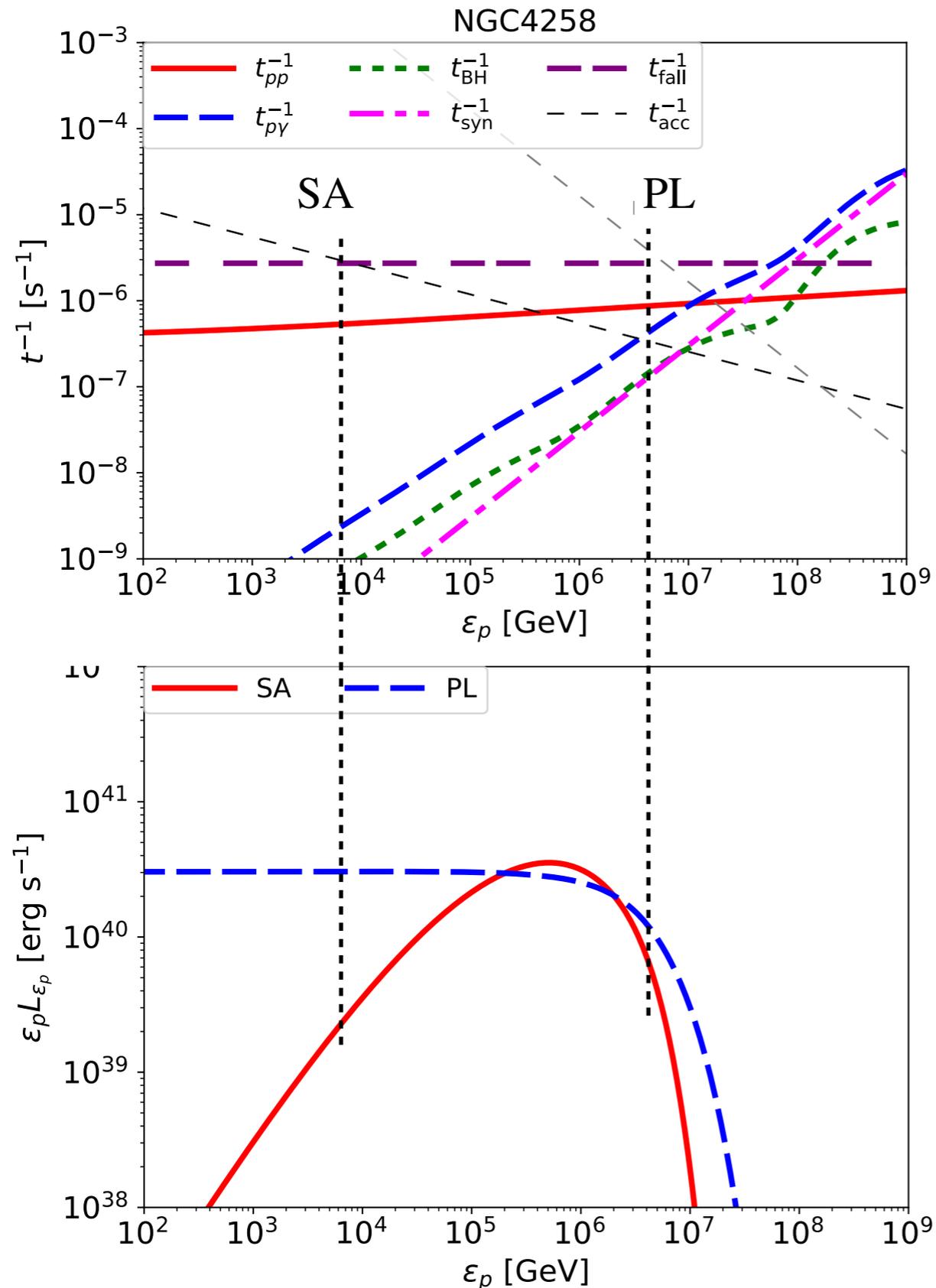
cf. SSK et al. 2015



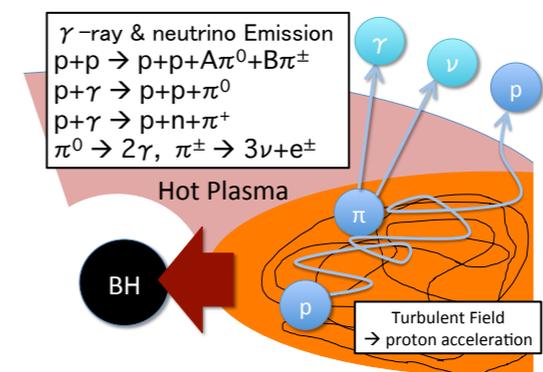
- Low-luminosity  $\rightarrow$  Poor observational data  
 $\rightarrow$  **Formulation based on theory**
- Thermal electrons in RIAFs emit seed photons
- Our results are consistent with X-ray observations



# Rates & CR Spectrum



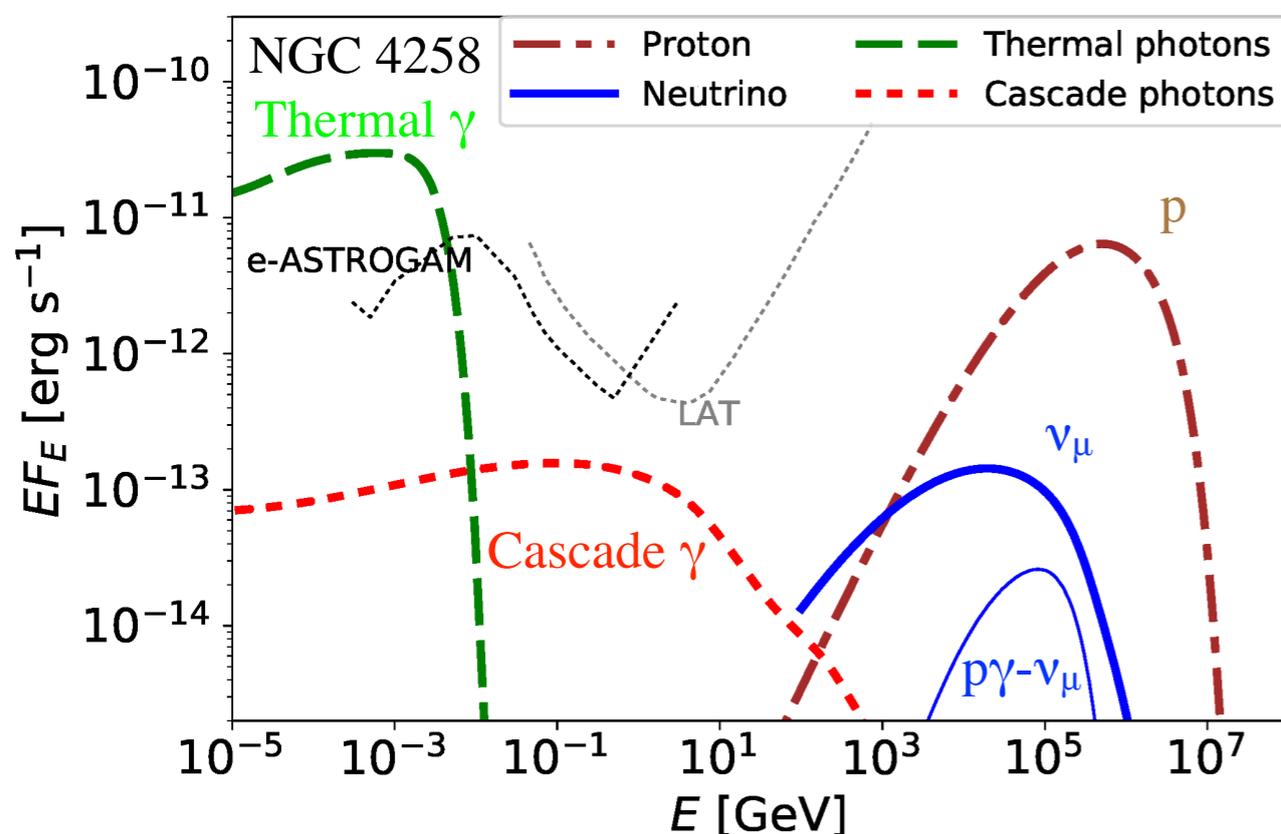
- **Infall is dominant**
- **Neutrino is mainly produced by pp**
- For SA, acceleration rate is much lower than PL model



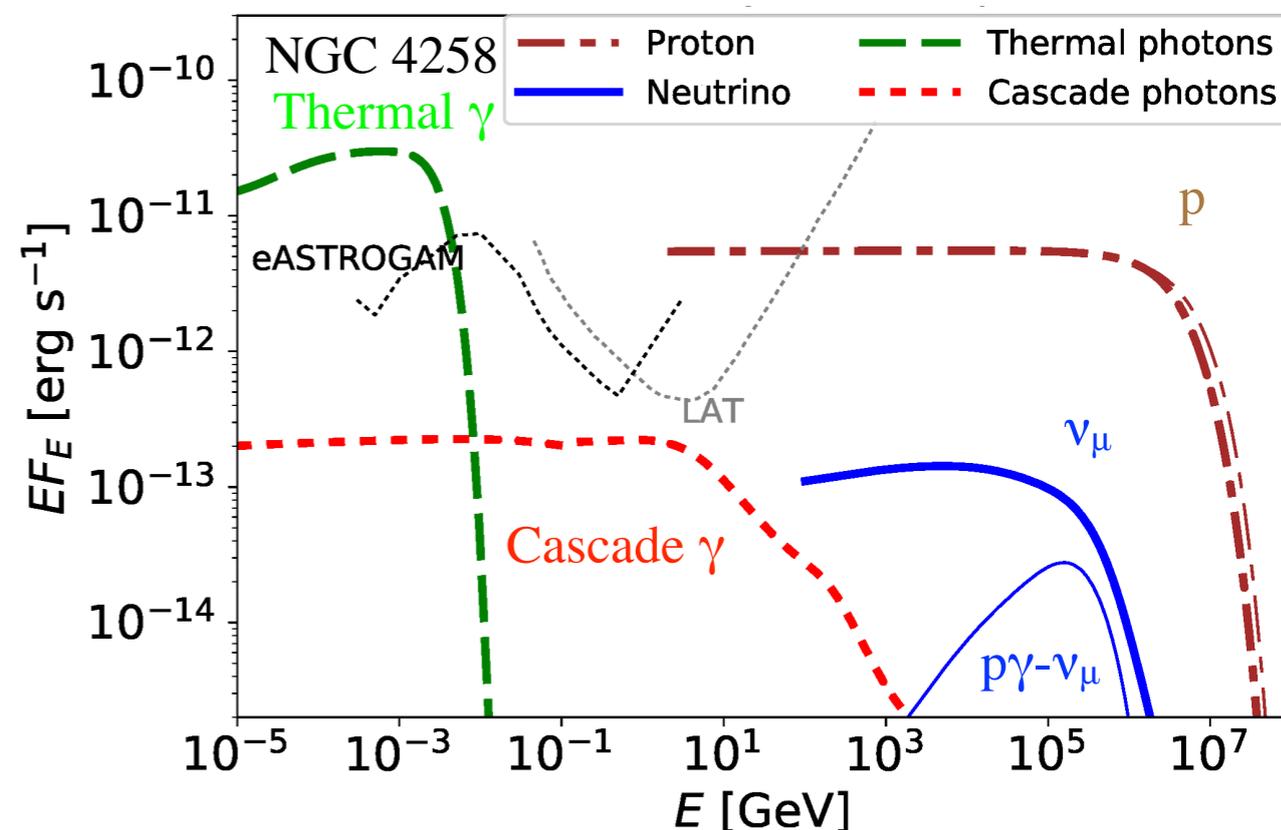
- Hard proton spectrum for SA
- Cutoff energy for SA is much higher than the critical energy ( $t_{acc} = t_{fall}$ ) due to hard spectrum & gradual cutoff

# HE Particle Spectrum

- Stochastic Acceleration (SA)



- Power-law Injection (PL)

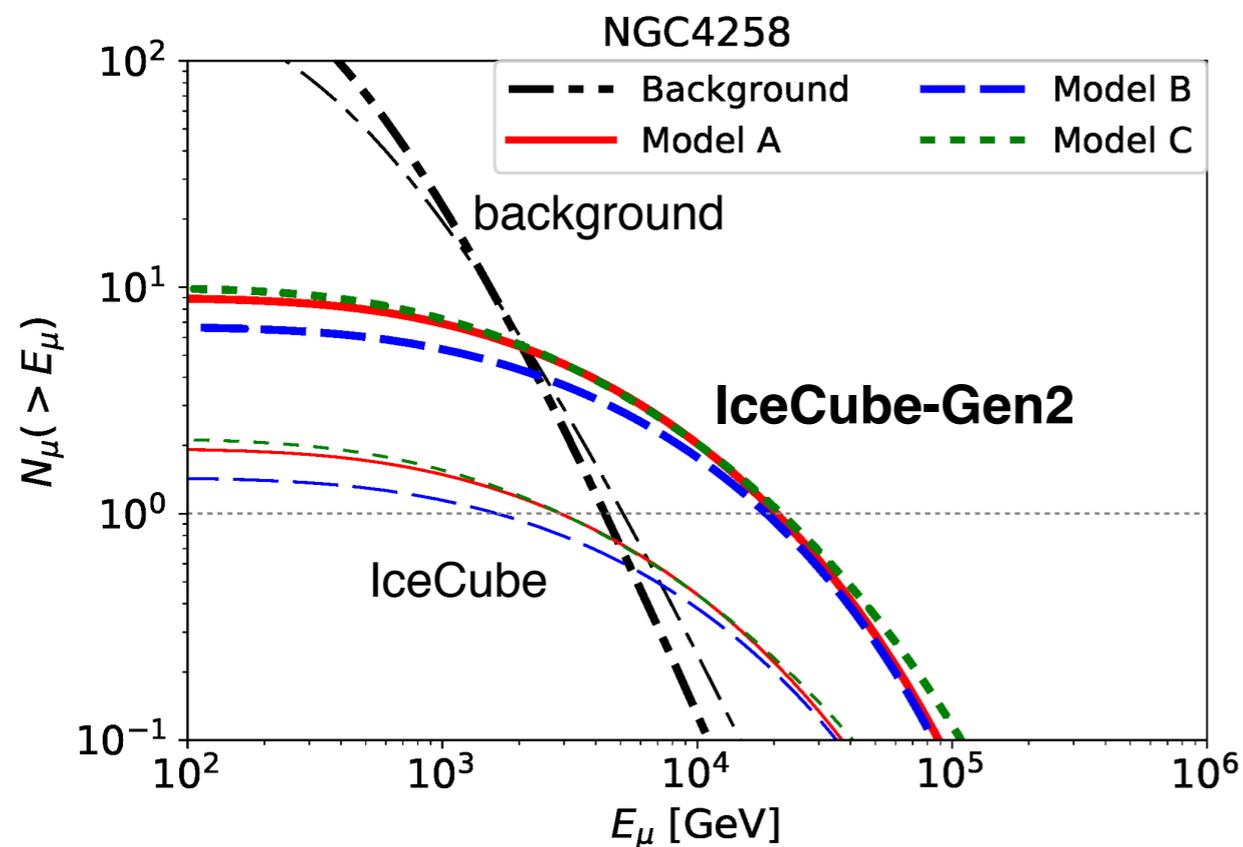


- pp-neutrino dominant,  $p\gamma$  contribution is comparable at high-energy
- Cascade  $\gamma$ -rays @  $E < 1$  GeV: **too faint to detect** by LAT & CTA
- Thermal  $\gamma$ -rays @  $E < \text{MeV}$ : **detectable** by future MeV satellites

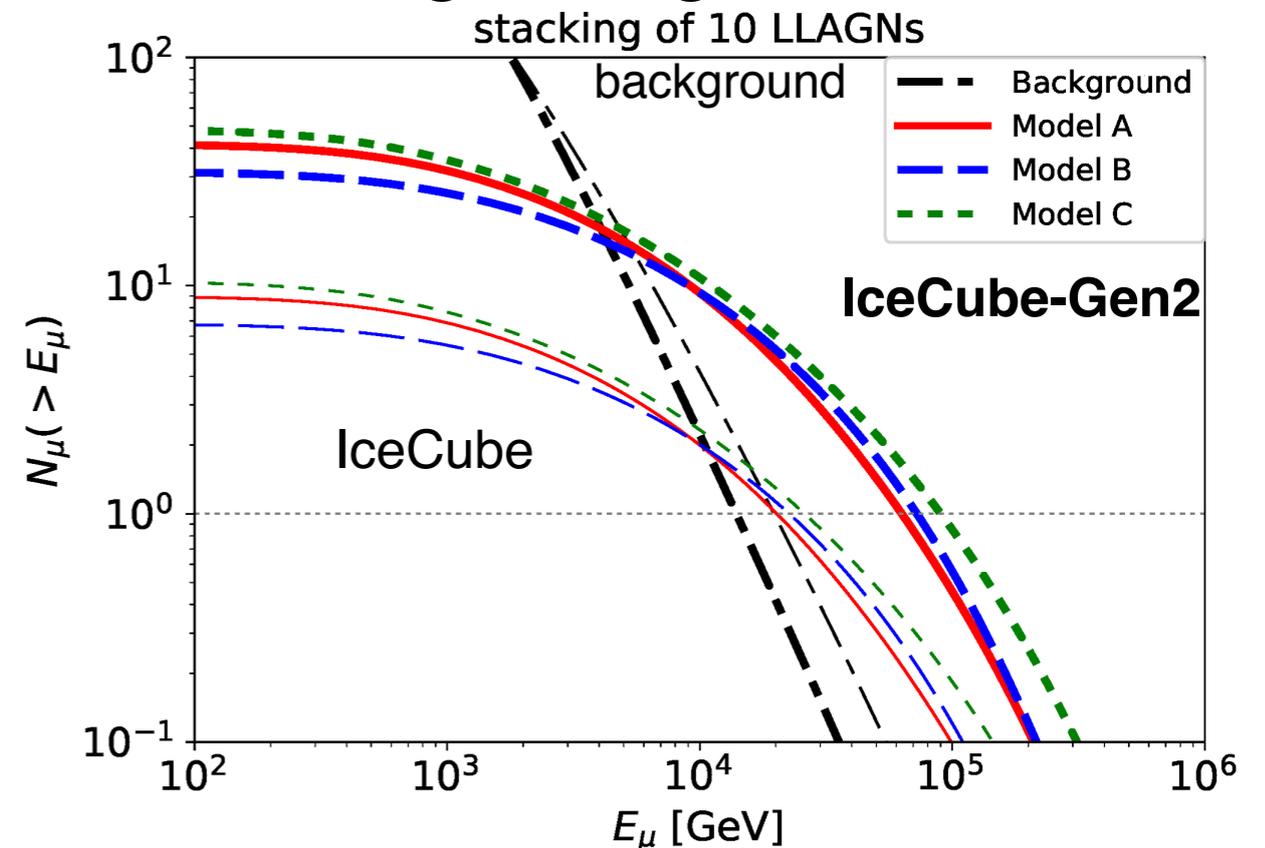
# Number of Muon Tracks

Murase & Waxman 2016

## A Bright LLAGN (NGC 4258)



## Stacking 10 brightest LLAGNs

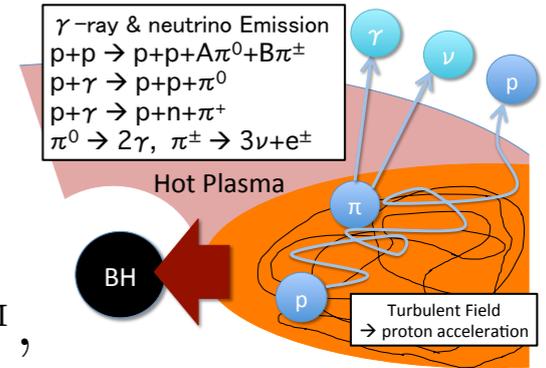


- IceCube cannot detect a neutrino
- **IceCube-Gen2 can detect** a few neutrinos of  $E > 10$  TeV

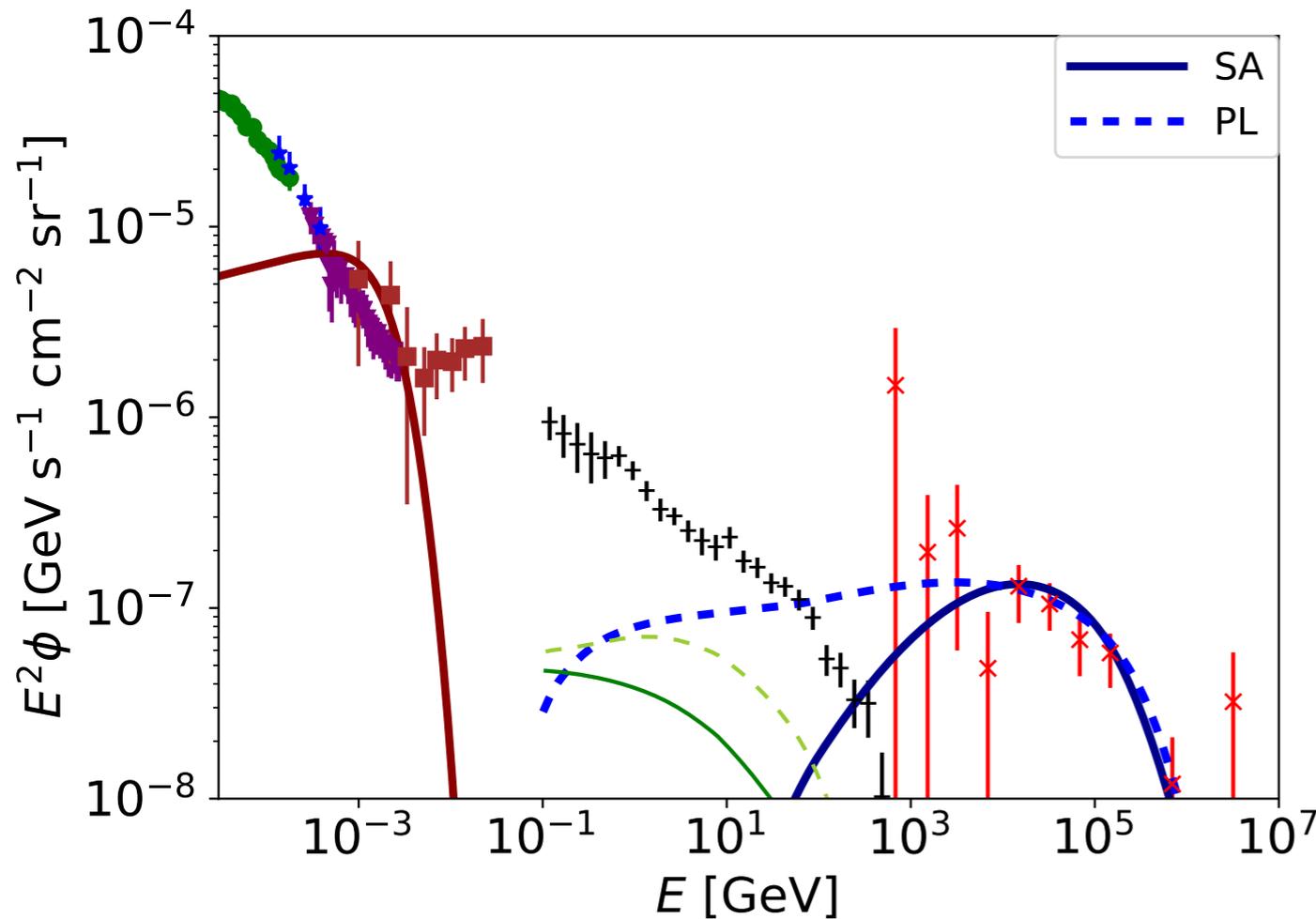
- IceCube cannot distinguish signals from the background
- **IceCube-Gen2 can detect** several neutrinos of  $E > 30$  TeV

# Extragalactic $\gamma$ & $\nu$ Backgrounds

$$\Phi_i = \frac{c}{4\pi H_0} \int \frac{dz}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} \int dL_{\text{H}\alpha} \rho_{\text{H}\alpha} \frac{L_{\epsilon_i}}{\epsilon_i} e^{-\tau_{i,\text{IGM}}},$$



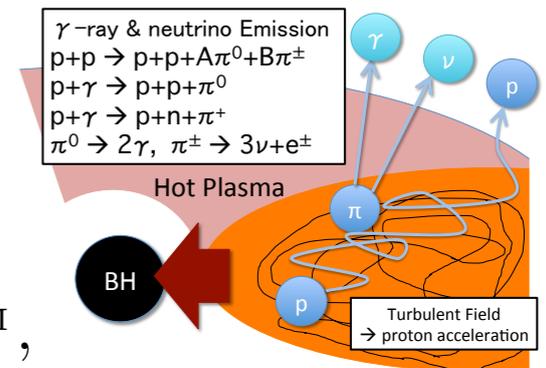
- AGNs with  $L_{\text{H}\alpha} < 4 \times 10^{41}$  erg/s equally contribute to  $\nu$
- LLAGNs with  $L_{\text{H}\alpha} \sim 4 \times 10^{41}$  erg/s mainly contribute to MeV  $\gamma$



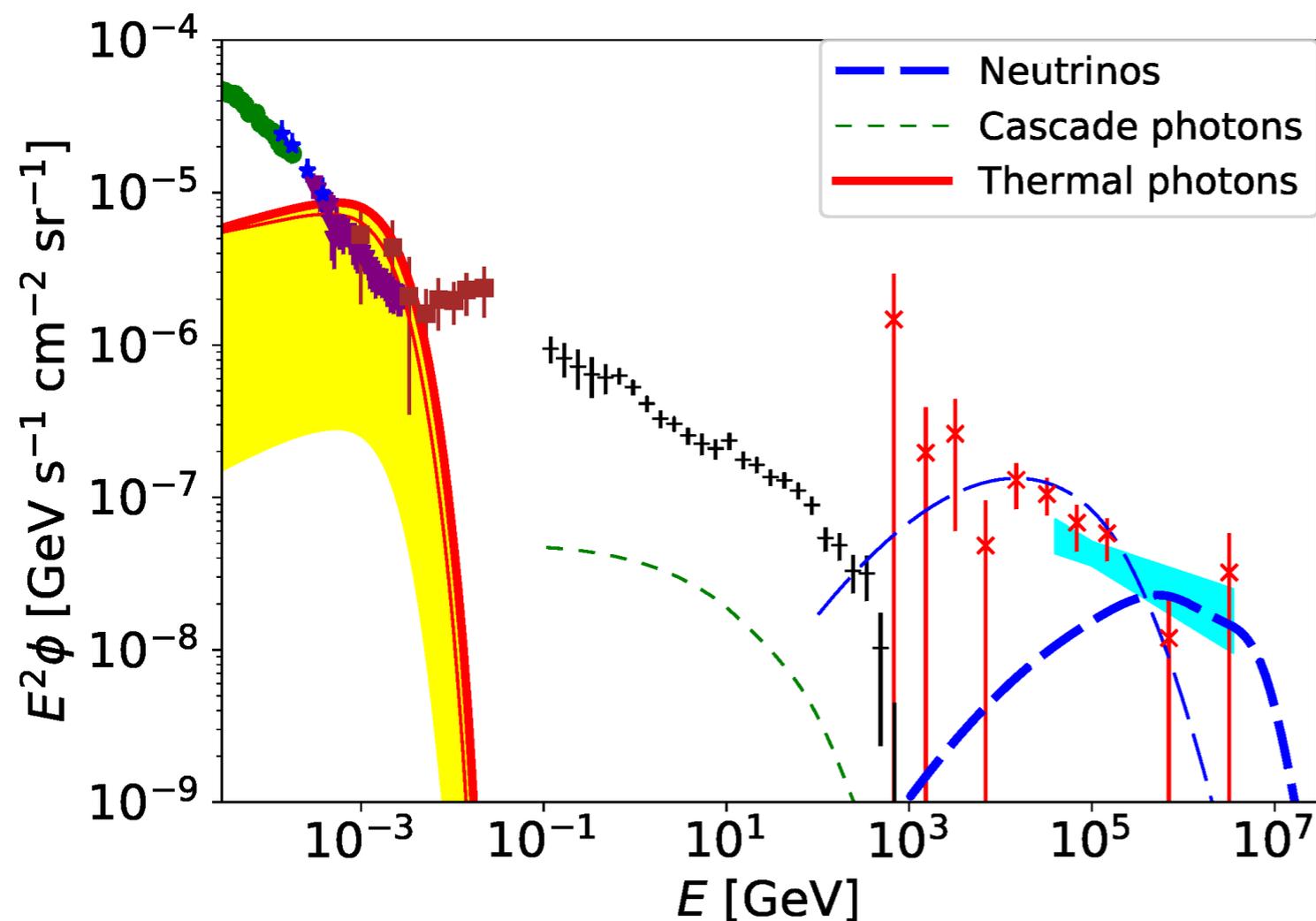
- **LLAGN can explain TeV-PeV  $\nu$  and MeV  $\gamma$  bkgnds simultaneously**
- GeV  $\gamma$ s are attenuated at RIAFs in LLAGNs  
 → consistent with Fermi data
- $P_{\text{CR}} \sim 0.1 P_{\text{th}}$  for SA,  
 $P_{\text{CR}} \sim 0.4 P_{\text{th}}$  for PL  
 → **Need hard spectrum**

# Extragalactic $\gamma$ & $\nu$ Backgrounds

$$\Phi_i = \frac{c}{4\pi H_0} \int \frac{dz}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} \int dL_{\text{H}\alpha} \rho_{\text{H}\alpha} \frac{L_{\epsilon_i}}{\epsilon_i} e^{-\tau_{i,\text{IGM}}},$$



- LLAGNs with  $L_{\text{H}\alpha} < 4 \times 10^{41}$  erg/s equally contribute to  $\nu$
- LLAGNs with  $L_{\text{H}\alpha} \sim 4 \times 10^{41}$  erg/s mainly contribute to MeV  $\gamma$



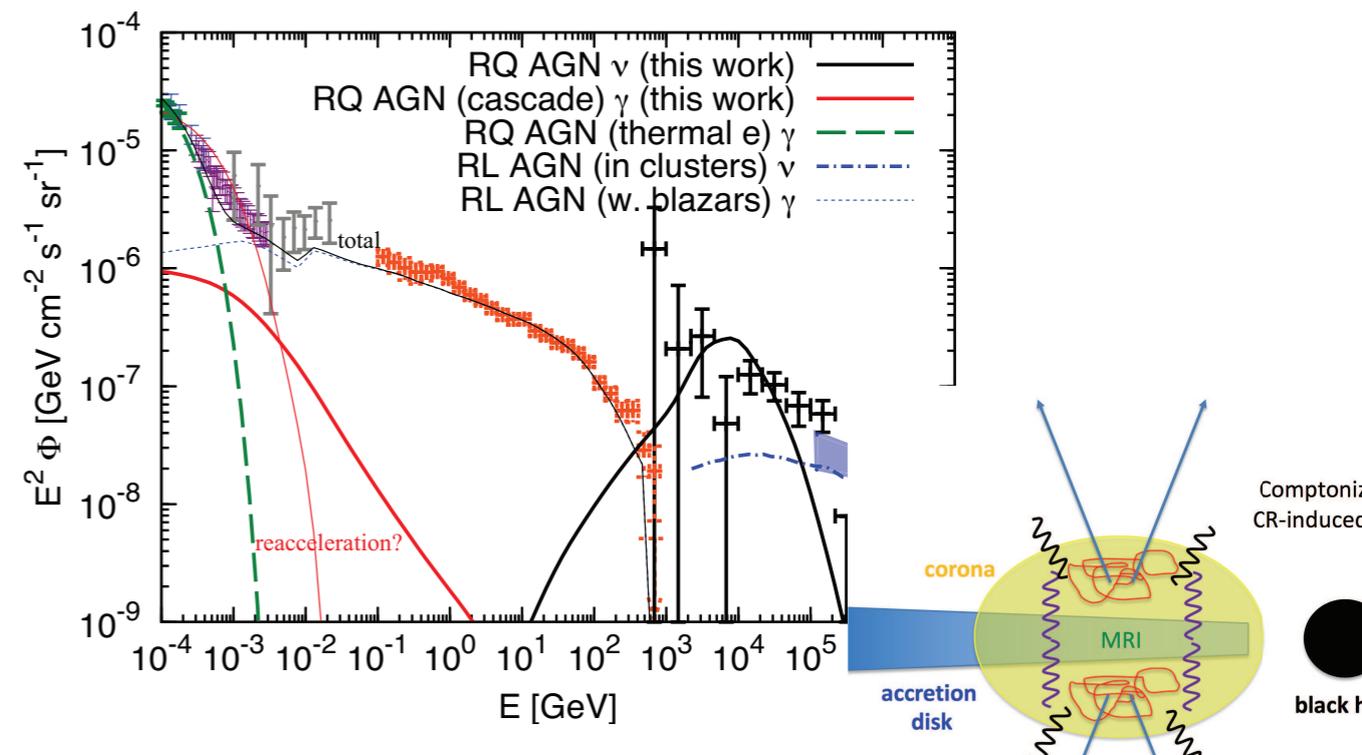
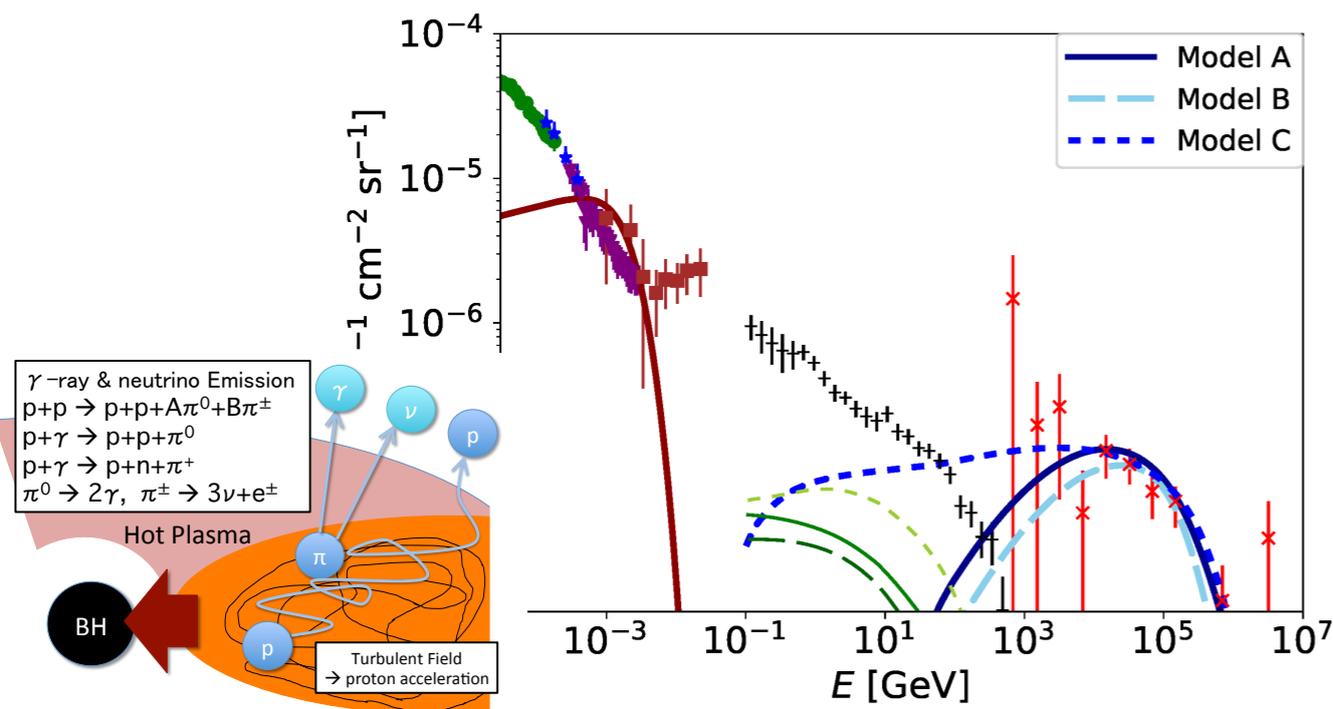
- **LLAGN can explain TeV-PeV  $\nu$  and MeV  $\gamma$  obs. simultaneously**
- GeV  $\gamma$ s are attenuated at RIAFs in LLAGNs  $\rightarrow$  consistent with Fermi data
- $P_{\text{CR}} \sim 0.01 P_{\text{th}}$  for PeV  $\nu$ ,  
 $P_{\text{CR}} \sim 0.1 P_{\text{th}}$  for TeV  $\nu$

# Index

- IceCube Neutrinos
- Accretion Flow in AGN  
Murase, SSK, Meszaros, arXiv:1904.04226
- AGN Corona model
- LLAGN RIAF model  
SSK, Murase, Meszaros, 2019, PRD, 100, 083014  
SSK, Murase, Meszaros in prep.
- **Summary**

# Summary

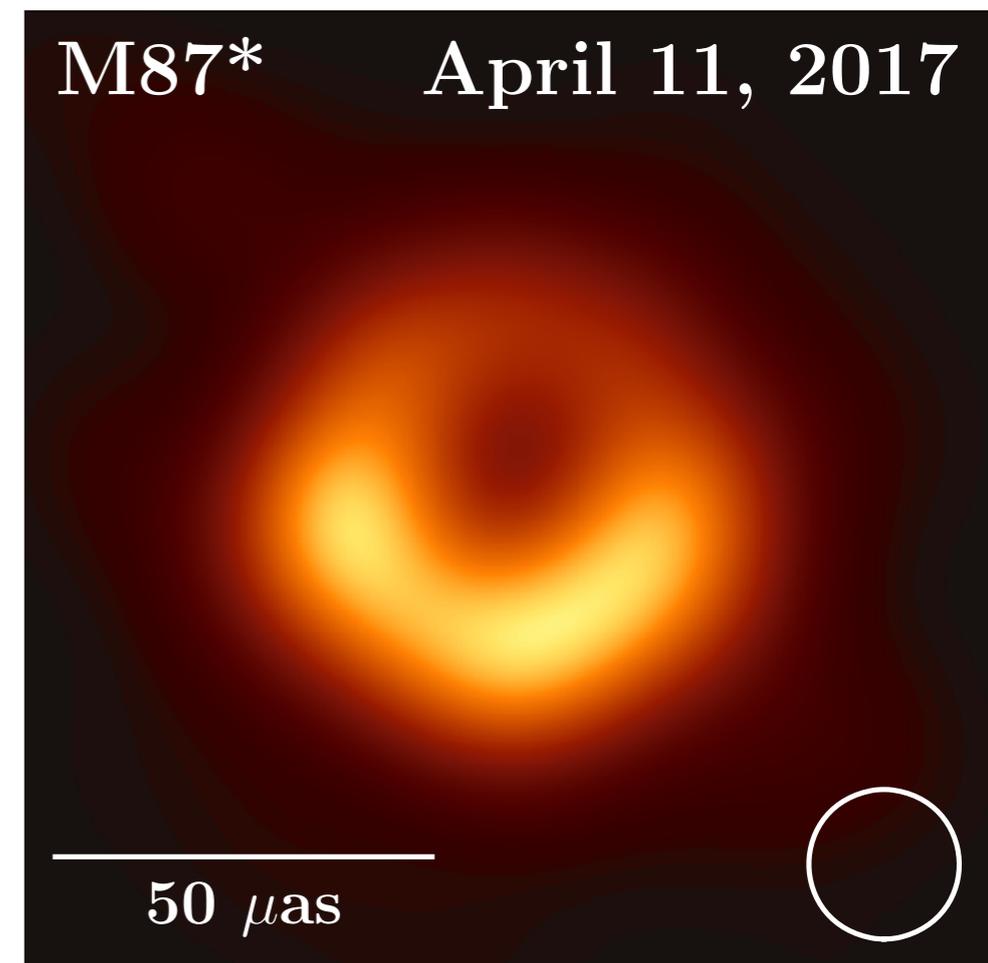
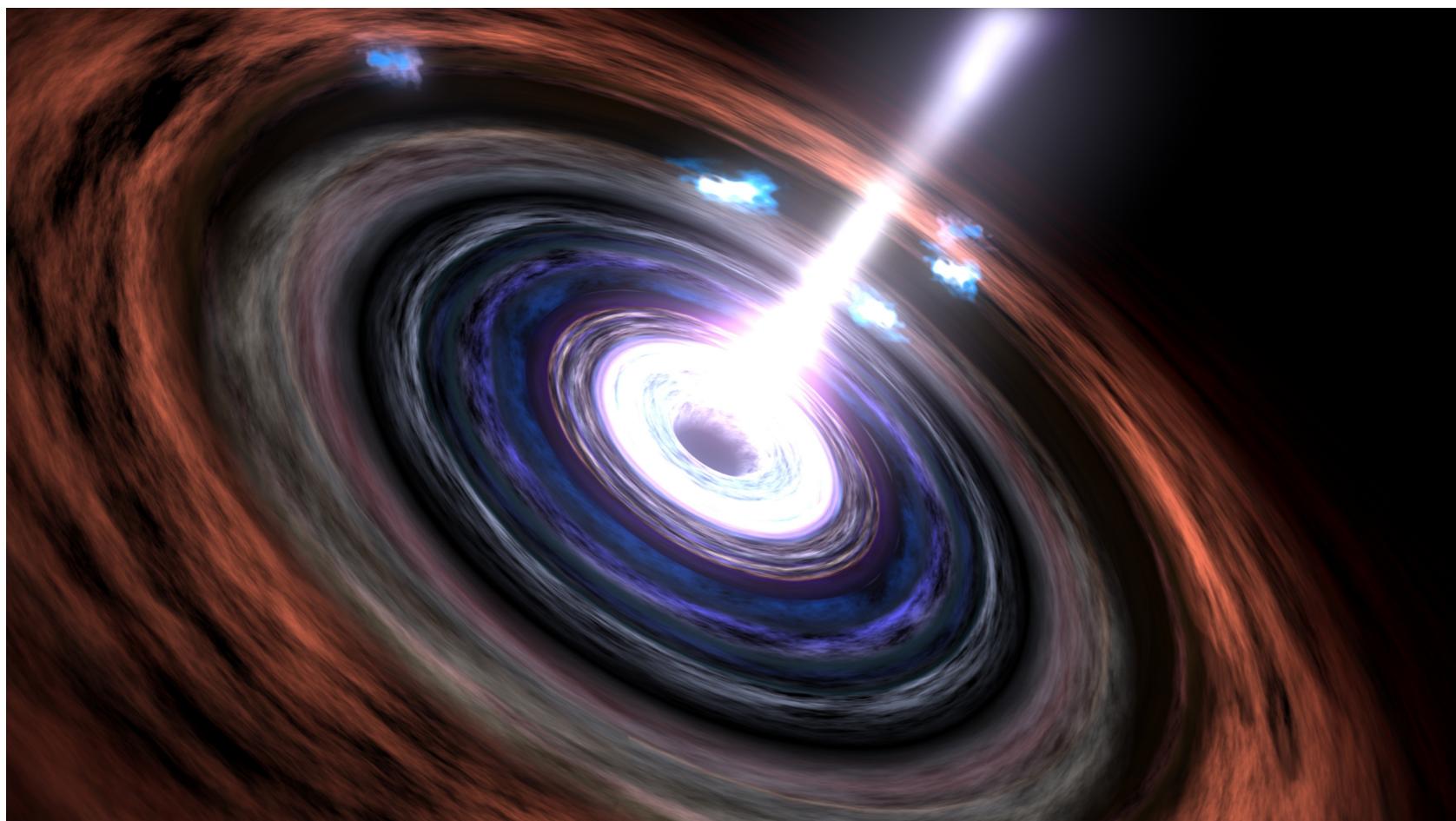
- Accretion flows in AGNs are a feasible neutrino source
- **Seyfert galaxies can reproduce TeV-PeV  $\nu$ s** without violating Fermi constraints with observation based parameters.
- **RIAFs in LLAGNs can explain MeV  $\gamma$  & TeV-PeV  $\nu$  backgrounds simultaneously** without violating Fermi constraints
- **Future multi-messenger observations can robustly test both models:**
  - IceCube-Gen2 can detect AGNs as point sources
  - AMEGO can detect MeV  $\gamma$  rays from AGNs



Thank you  
for your attention

# Active Galactic Nuclei (AGNs)

- Accretion onto Supermassive black hole ( $M_{\text{BH}} \sim 10^8 M_{\text{sun}}$ )  
gravitational energy  $\rightarrow$  radiation or thermal energy
- SMBH paradigm is proved by Event Horizon Telescope



# Plasma Conditions in Accretion Flow

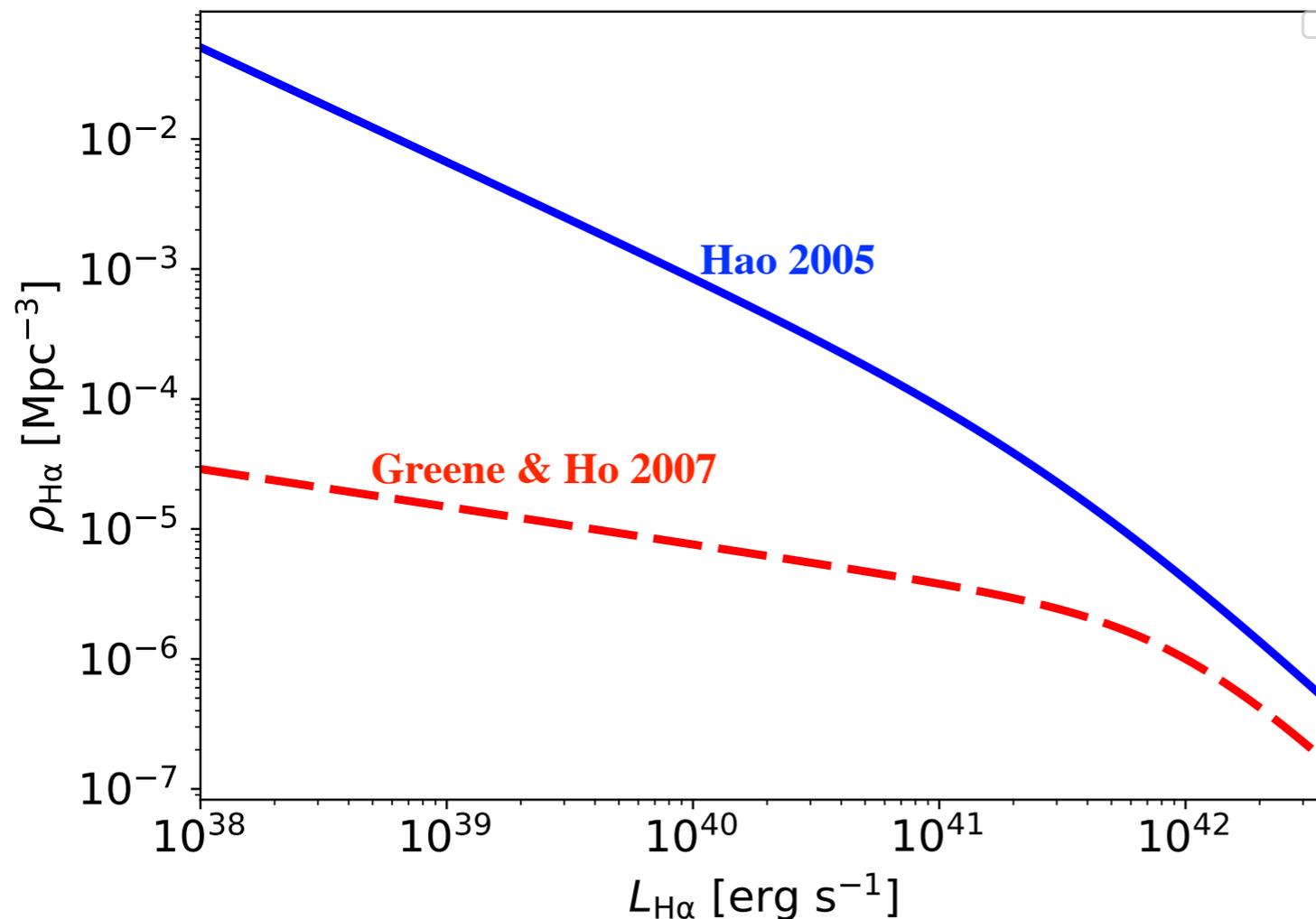
- To accelerate non-thermal particles, relaxation time  $>$  dissipation time
- For RIAFs in LLAGNs,  $t_{\text{dis}} \sim t_{\text{dyn}} \sim R/V_{\text{fall}}$
- For Coronae in QSO,  $t_{\text{dis}} \sim H/V_A$
- **Protons are collisionless** for both cases  
→ **Non-thermal Proton**
- **Electrons are collisional** for both cases  
→ **Thermal electrons only**

# Luminosity Function

- $\nu$  &  $\gamma$  intensities from LLAGNs

$$\Phi_i = \frac{c}{4\pi H_0} \int \frac{dz}{\sqrt{(1+z)^3 \Omega_m + \Omega_\Lambda}} \int dL_{\text{H}\alpha} \rho_{\text{H}\alpha} \frac{L_{\varepsilon_i}}{\varepsilon_i} e^{-\tau_{i,\text{IGM}}},$$

- $\rho_{\text{H}\alpha}$ : H $\alpha$  line Luminosity function



We use higher LF

# Target Photon Field

- Low-luminosity → Poor observational data  
→ **Theory driven formulation**
- Thermal electrons in RIAFs emit seed photons
- Provide X-ray luminosity by observation  
→ Bolometric correction based on AGN survey  
→ Estimate mass accretion rate  
→ Obtain physical quantities ( $\rho$ ,  $B$ ,  $n$ ,  $T_e$ ) in RIAFs  
→ Calculate target photon spectrum by one-zone approximation
- **We do not adjust X-ray luminosity**

# EM Cascades in IGM

- Cutoff energy by  $\gamma\gamma$  pair production in RIAFs  
 $E_{\text{cut}} \sim 0.1 - 100 \text{ GeV}$
- $\tau_{\gamma\gamma, \text{IGM}} \sim 1$  for  $E_\gamma = 100 \text{ GeV}$  @  $z=0.5$   
 → **We need to consider attenuation**
- $\gamma\gamma$  pair production:  $e^+e^-$  of  $\gamma_e \sim 10^5$  for  $E_\gamma \sim 100 \text{ GeV}$   
 →  $E_{\text{ic}} \sim 4\gamma_e^2 E_{\text{CMB}}/3 \sim 10 \text{ MeV}$   
 → **we can ignore EM cascade in IGM**
- $\tau \ll 1$  for  $\nu$  & MeV  $\gamma$   
 → we can ignore attenuations

# Implications & Caveats

- **Multi-messenger tests are promising:**  
Nearby LLAGNs are detectable by IC-Gen2 & e-ASTROGAM
- High source number density ( $\sim 10^{-3} \text{ Mpc}^{-3}$ )  
→ LLAGNs can **avoid the point-source constraints**
- Luminosity Function (LF) is very uncertain
  - LF by Hao et al. (2005)  $\gg$  Greene & Ho (2007)
  - If we use Greene & Ho (2007), neutrino flux becomes too low to explain TeV-PeV neutrinos