

ブロッタホール 磁気圏での 粒子加速領域

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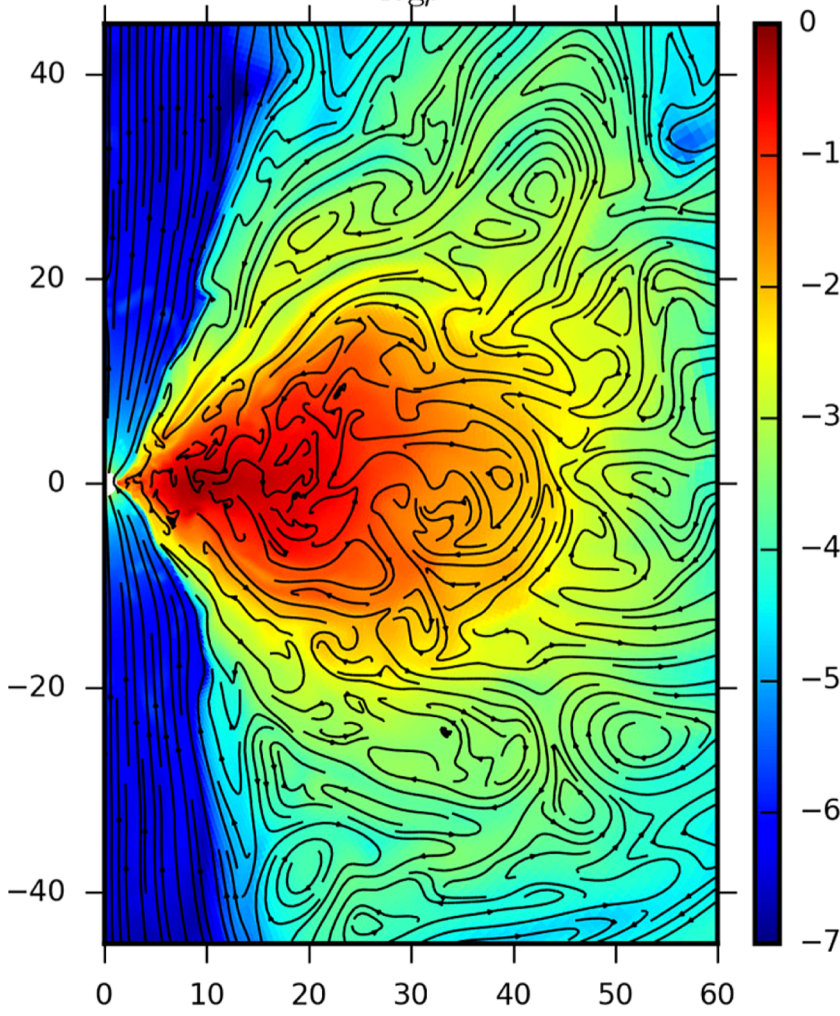
GRMHD Numerical Simulations

Artificial mass supply
in the jet region.

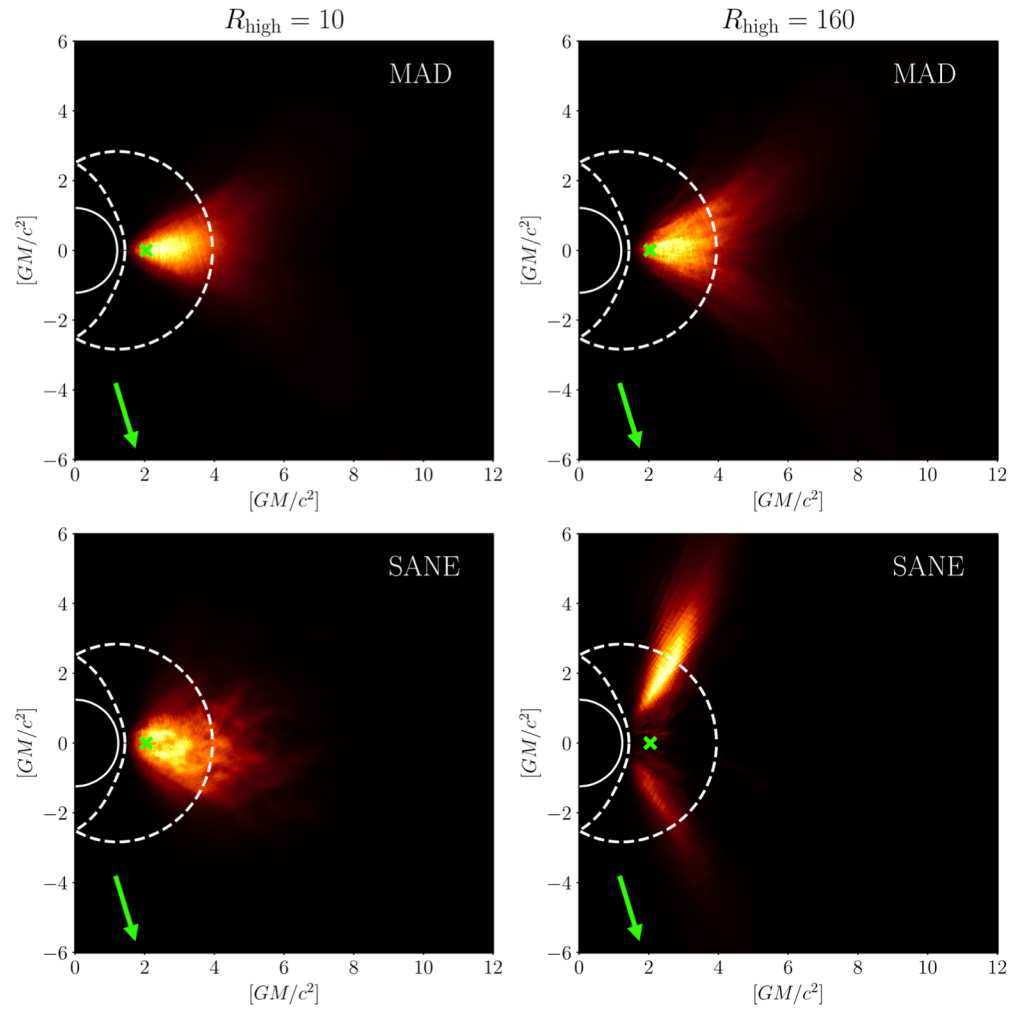
Predictions of the jet emission
are highly uncertain.

FCT, $t = 2000.0 M$ Porth+ 17

$\log \rho$

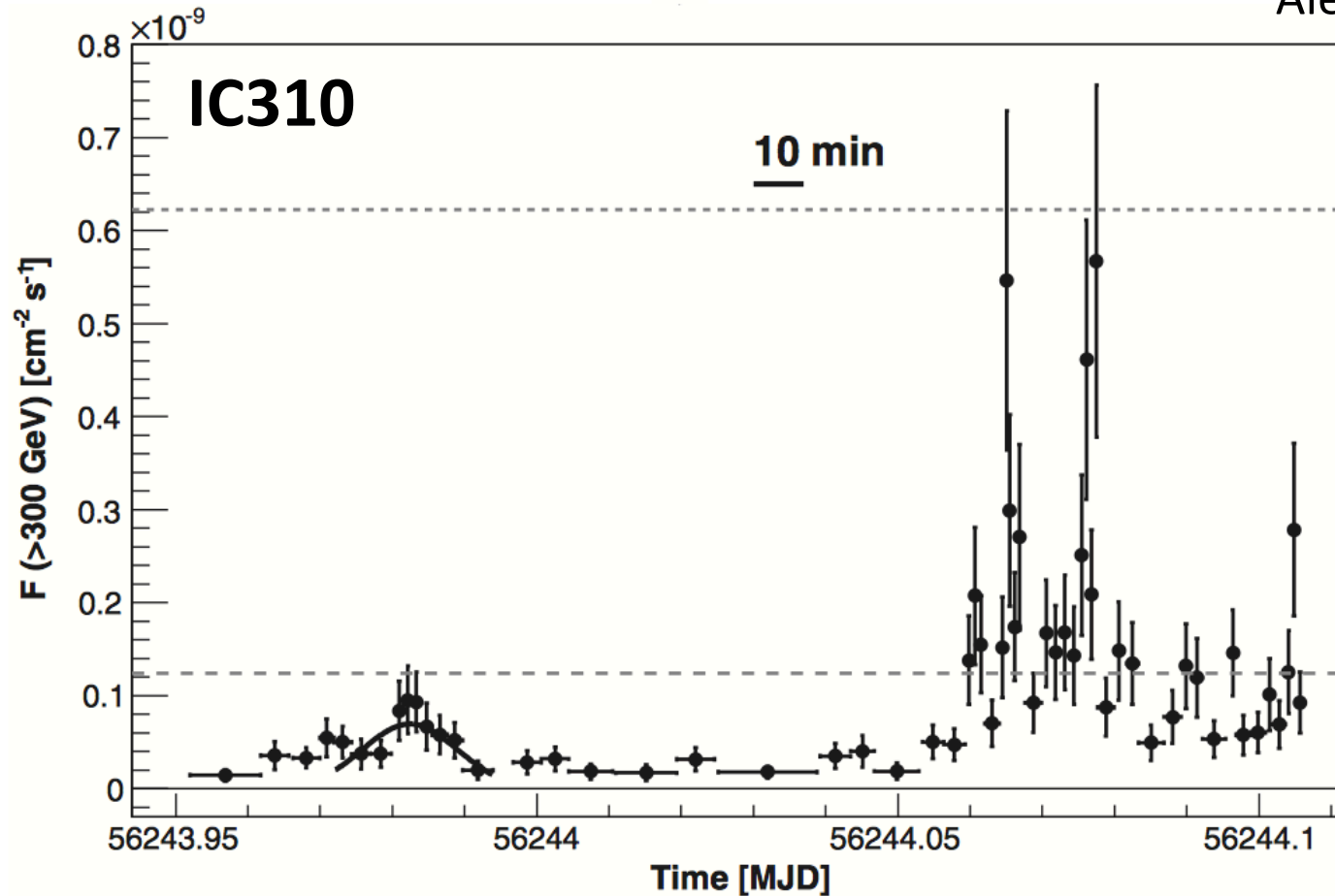


EHT Coll. 19



TeV flare from radio galaxy

Aleksić+ 14

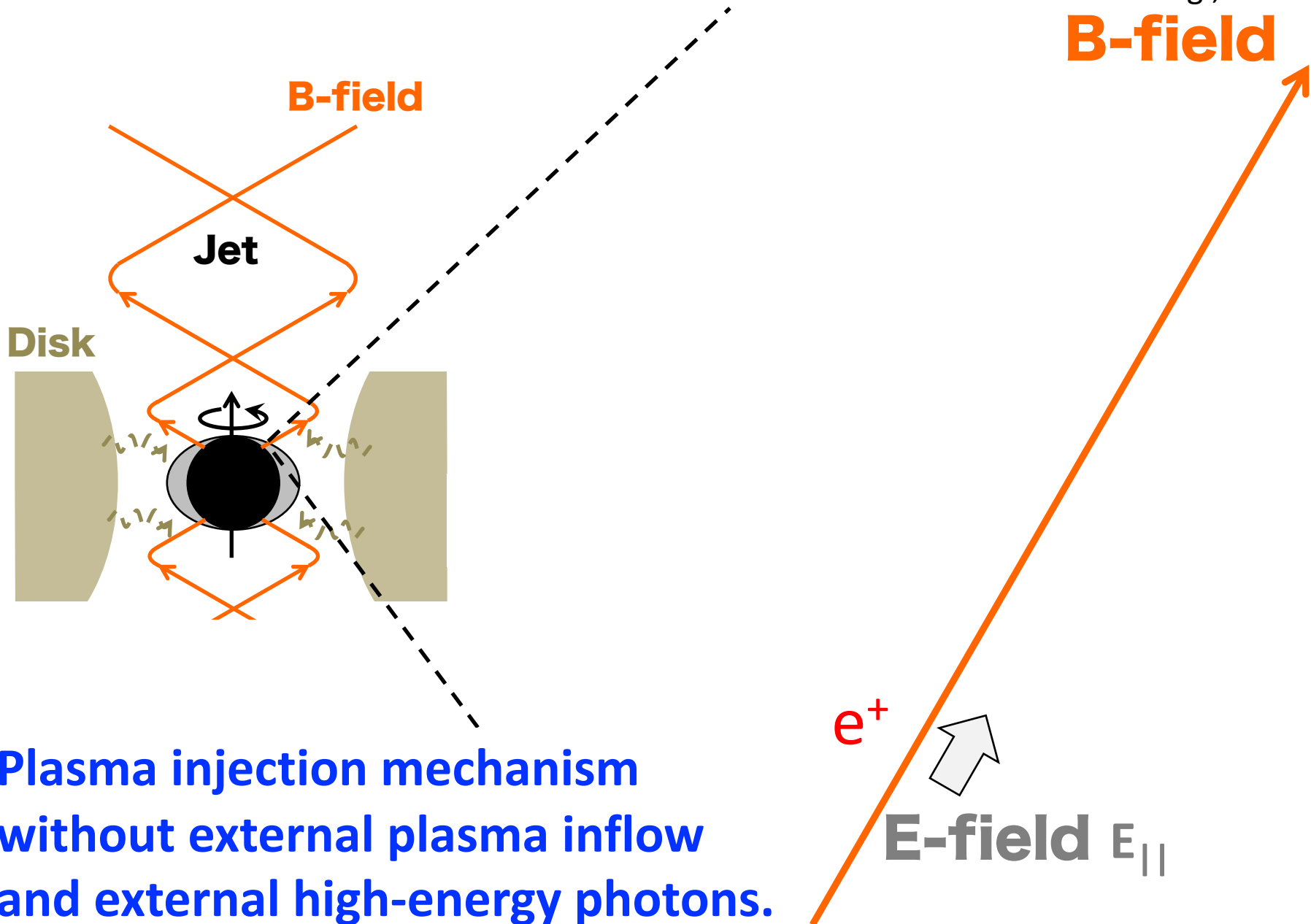


Flux doubling timescale < 4.8 min at 95% C.L.
corresponds to $\sim 20\%$ of the timescale r_g/c .

→ Particle acceleration at sub-horizon scale?

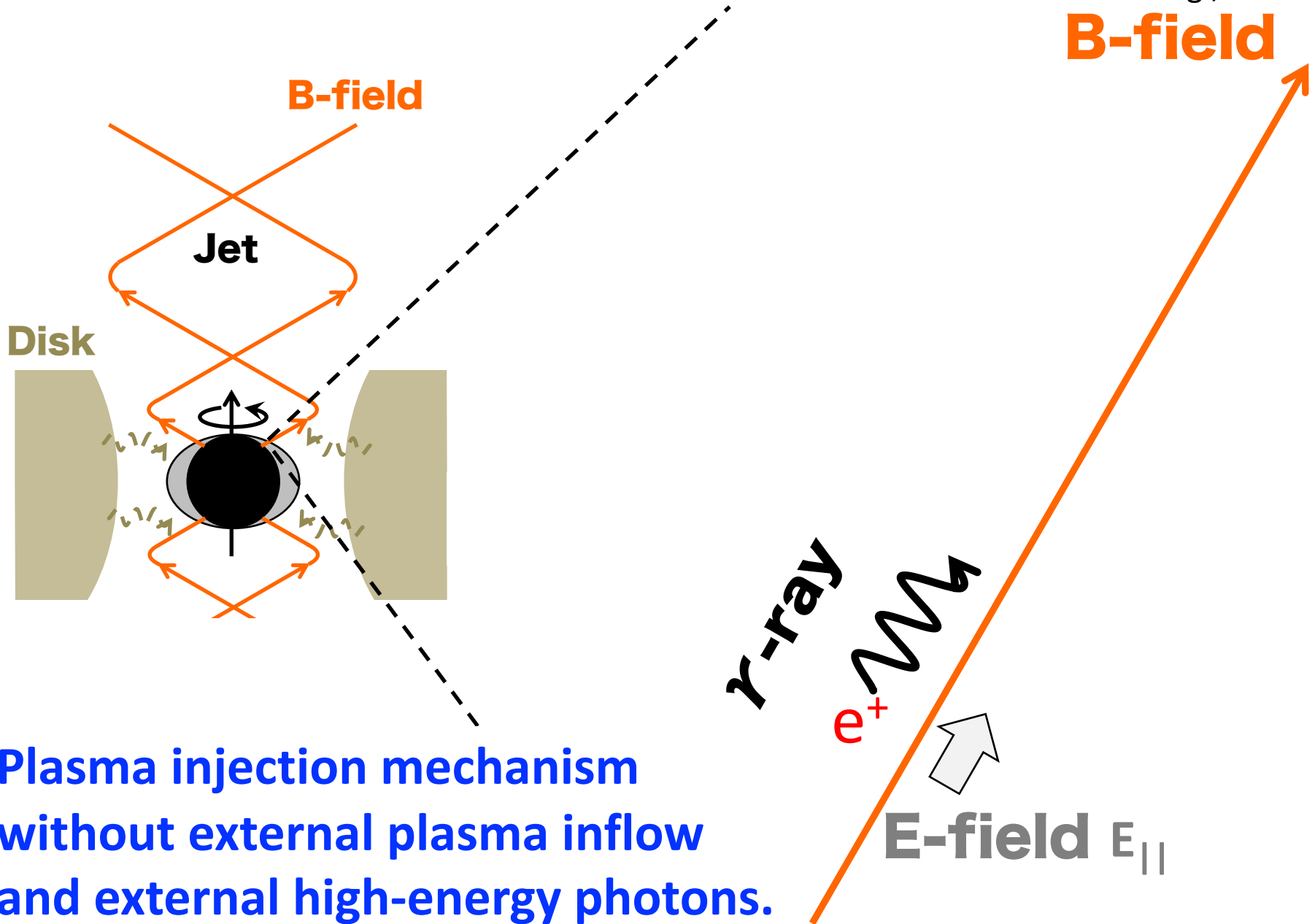
Electromagnetic Cascade

e.g., Beskin+ 92



Electromagnetic Cascade

e.g., Beskin+ 92



B-field

Jet

Disk

r-ray

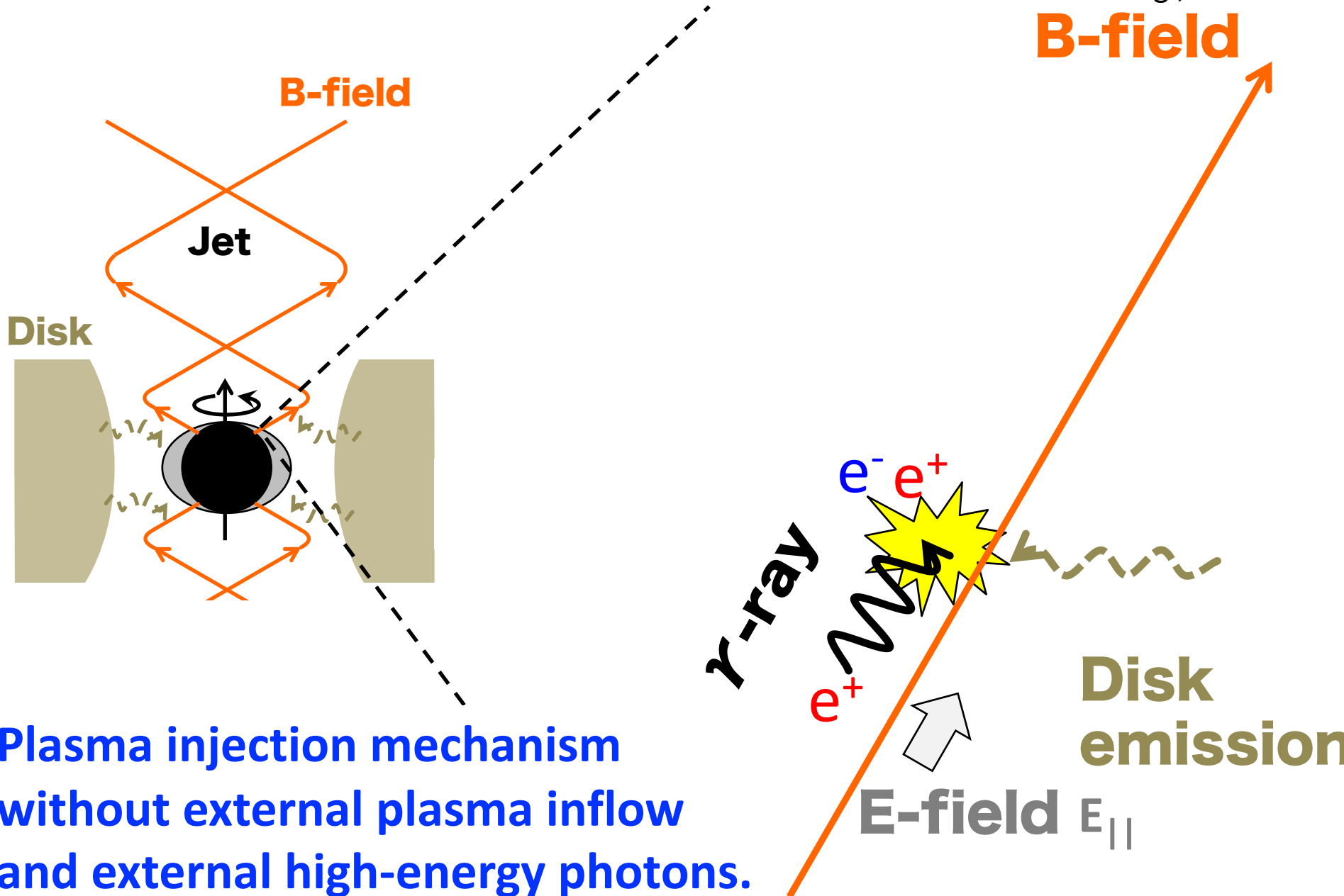
e^+ M

E-field $E_{||}$

**Plasma injection mechanism
without external plasma inflow
and external high-energy photons.**

Electromagnetic Cascade

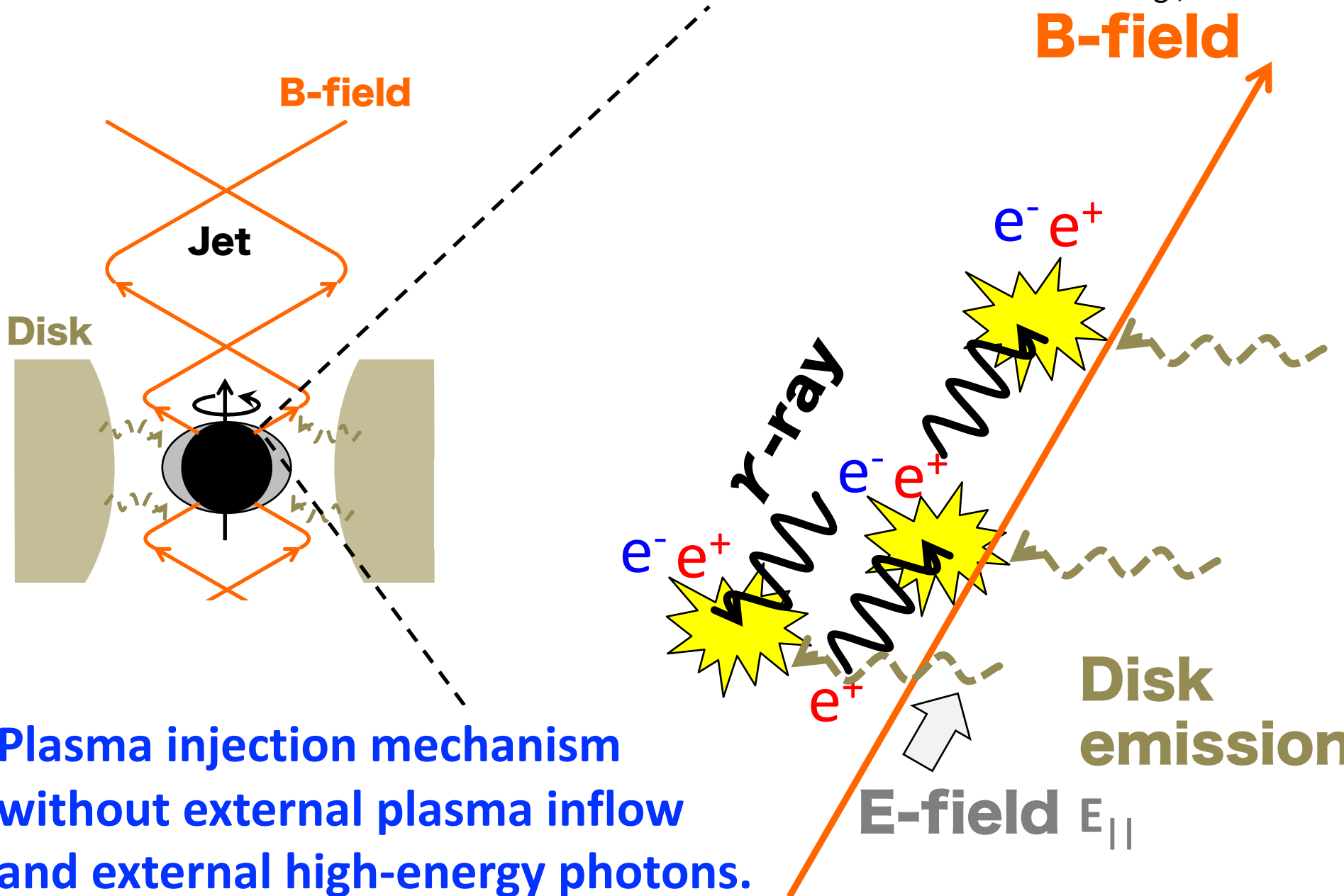
e.g., Beskin+ 92



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Electromagnetic Cascade

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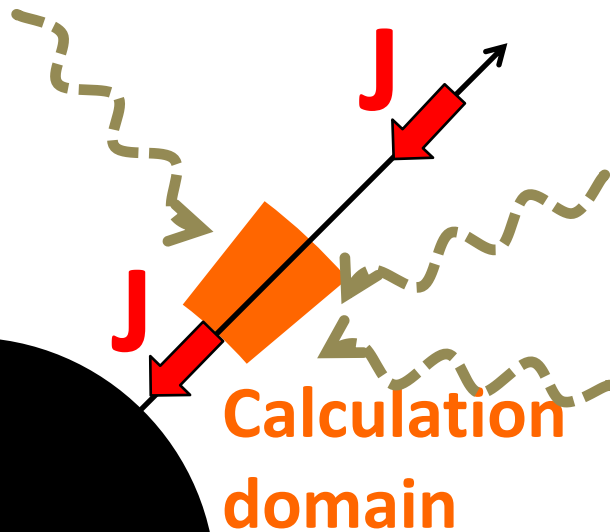


1D PIC Model

Levinson & Cerutti 18

- 1-dimensional structure: the gap extends along a poloidal magnetic surface as a function of θ .
 - Ignoring any MHD waves, considering only plasma oscillations.
- The gap constitutes a small disturbance.
 - The activity does not significantly affect the global structure (the B-field geometry and the angular velocity).
- Isotropic radiation field (from accretion disk) for seed photons.

$$I_s(x^\mu, \epsilon_s, \Omega_s) = I_0(\epsilon_s/\epsilon_{s,\min})^{-p}, \quad \epsilon_{s,\min} < \epsilon_s < \epsilon_{s,\max}$$



- No external plasma source.
- The global current is a free parameter.
- A split monopole geometry for the global B-field.
- The angular velocity of magnetic surface $\Omega = 0.5\omega_H$.

Accretion disk

Background spacetime

Kerr metric given in BL coordinates

$$ds^2 = -\alpha^2 dt^2 + g_{\varphi\varphi}(d\varphi - \omega dt)^2 + g_{rr}dr^2 + g_{\theta\theta}d\theta^2$$

$$\alpha^2 = \frac{\Sigma\Delta}{A}; \quad \omega = \frac{2ar_g r}{A}; \quad g_{rr} = \frac{\Sigma}{\Delta};$$

$$g_{\theta\theta} = \Sigma; \quad g_{\varphi\varphi} = \frac{A}{\Sigma} \sin^2 \theta,$$

$$\Delta = r^2 + a^2 - 2r_g r, \quad \Sigma = r^2 + a^2 \cos^2 \theta, \quad A = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta.$$

Tortoise coordinate

$$\xi(r) = \frac{1}{r_+ - r_-} \ln \left(\frac{r - r_+}{r - r_-} \right)$$

$$\xi \rightarrow -\infty \text{ as } r \rightarrow r_H = r_+$$

$$\xi \rightarrow 0 \text{ as } r \rightarrow \infty$$

$$r_{\pm} = 1 \pm \sqrt{1 - \tilde{a}^2}$$

Basic equations

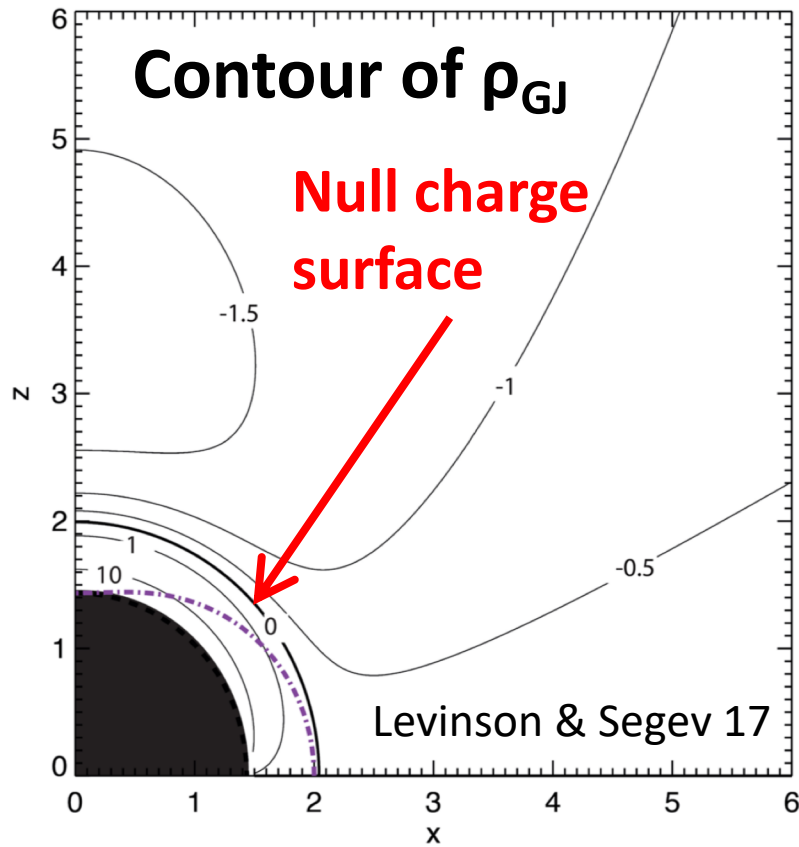
Levinson & Cerutti 18

Gauss's law

$$\partial_{\mu}(\sqrt{-g}F^{t\mu}) = (\sqrt{-g}j^t)$$

$$\rightarrow \partial_{\xi}(\sqrt{A}E_r) = 4\pi\Delta\Sigma(j^t - \rho_{\text{GJ}})$$

$$\rho_{\text{GJ}} = \frac{B_{\text{H}} \sqrt{A_{\text{H}}}}{4\pi \sqrt{-g}} \left[\frac{\sin^2 \theta}{\alpha^2} (\omega - \Omega) \right]_{,\theta}$$



Basic equations

Levinson & Cerutti 18

Gauss's law

$$\partial_{\mu}(\sqrt{-g}F^{t\mu}) = (\sqrt{-g}j^t)$$

$$\rightarrow \partial_{\xi}(\sqrt{A}E_r) = 4\pi\Delta\Sigma(j^t - \rho_{\text{GJ}})$$

Ampère's law

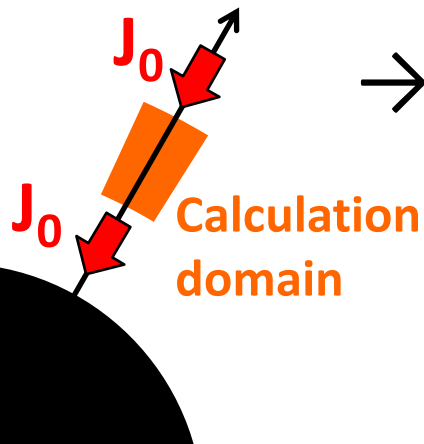
(radial component)

$$\rho_{\text{GJ}} = \frac{B_H \sqrt{A_H}}{4\pi \sqrt{-g}} \left[\frac{\sin^2 \theta}{\alpha^2} (\omega - \Omega) \right]_{,\theta}$$

$$\partial_{\mu}(\sqrt{-g}F^{r\mu}) = (\sqrt{-g}j^r)$$

$$\rightarrow \partial_t(\sqrt{A}E_r) = -4\pi(\Sigma j^r - J_0)$$

$$J_0 = \frac{1}{4\pi \sin \theta} \left(\frac{\Delta \sin \theta}{\Sigma} F_{r\theta} \right)_{,\theta}$$



Basic equations

Levinson & Cerutti 18

Equation of motion for i -th particle

Curvature
loss term

$$\frac{du_i^\mu}{d\tau_i} = -\Gamma^\mu_{\alpha\beta} u_i^\alpha u_i^\beta + \frac{q_i}{m_e} F^\mu_{\alpha} u_i^\alpha + \underline{s_i^\mu}$$

$$\frac{du_{i\mu}}{d\tau_i} = \Gamma_{\alpha\mu\beta} u_i^\alpha u_i^\beta + \frac{q_i}{m_e} F_{\mu\alpha} u_i^\alpha + \underline{s_{i\mu}}$$

$$\begin{aligned} \rightarrow \frac{du_i}{dt} &= \sqrt{g^{\text{rr}}} \left[-\gamma_i \partial_r(\alpha) + \frac{q_i}{m_e} F_{\text{rt}} \right] - \frac{s_{it}}{u_i} \\ &= -\sqrt{g^{\text{rr}}} \gamma_i \partial_r(\alpha) + \alpha \left(\frac{q_i}{m_e} E_r - \frac{P_{\text{cur}}(\gamma_i)}{m_e v_i} \right) \end{aligned}$$

$$\begin{aligned} \rightarrow \frac{d\xi_i}{dt} &= \frac{1}{\Delta} \frac{dr_i}{dt} = \frac{1}{\Delta} \frac{u_i^r}{u_i^t} = \frac{v_i}{\sqrt{A}} & \partial_r(\alpha) &= \frac{\alpha}{A} \left(\frac{2r^2 \tilde{a}^2 \sin^2 \theta}{\Sigma} + \frac{r^4 - \tilde{a}^4}{\Delta} \right) \\ & & P_{\text{cur}}(\gamma) &= \frac{2}{3} \frac{e^2 \gamma^4 v^4}{R_c^2} \end{aligned}$$

Basic equations

Levinson & Cerutti 18

Equation of motion for i -th photon

$$\frac{d\tilde{p}_k^r}{dt} = - \sqrt{g^{rr}} \tilde{p}_k^t \partial_r (\alpha)$$

$$\frac{d\xi_k}{dt} = \frac{1}{\sqrt{A}} \frac{\tilde{p}_k^r}{\tilde{p}_k^t}$$

IC scattering

$$\delta\tau_{\text{sc}} = \int_{r(t)}^{r(t+\delta t)} \kappa_{\text{KN}} \sqrt{g_{rr}} dr$$

Pair production

$$\delta\tau_{\text{pp}} = \int_{r(t)}^{r(t+\delta t)} \kappa_{\text{pp}} \sqrt{g_{rr}} dr$$

Parameters

Fiducial optical depth

$$\tau_0 = 4\pi r_g \sigma_T I_0 / hc$$

Global current density

$$\dot{j}_0$$

BH mass

$$M_{\text{BH}} = 10^9 M_\odot$$

Dimensionless spin parameter

$$a_* = 0.9$$

B-field on the horizon

$$B_{\text{H}} = 2\pi \times 10^3 \text{G}$$

Inclination angle of magnetic surface $\theta = 30^\circ$

Minimum energy of seed photon

$$\epsilon_{s,\text{min}}$$

Slope of seed photon spectrum

$$p = 2$$

Curvature radius

$$R_{\text{cur}} = r_g$$

Number of cell

$$N = 32768$$
$$\gtrsim \frac{r_g}{l_p} \sim 10^3 \sqrt{\frac{\kappa M_9 B_{\text{H},3}}{\langle \gamma_8 \rangle}}$$

We neglect the scattering of particles with $\gamma < 10^7$.

Pair multiplicity in gap

Pairs quickly accelerate to the terminal Lorentz factor in the gap.

$$eE_{\parallel} = P_{\text{rad}}/c \quad \rightarrow \quad \gamma_{\text{max}} \sim 10^{10}$$

$$\tau_0 = 4\pi r_g \sigma_T I_0 / hc$$

Minimum energy
of soft photons

$$\epsilon_{\text{min}} = 10^{-9} m_e c^2$$

Scattering optical
depth (KN)

$$\tau_{\text{IC}}/\tau_0 \sim O(10^{-1})$$

Pair creation
optical depth

$$\tau_{\gamma\gamma}/\tau_0 \sim O(10^{-1})$$

Pair multiplicity
in the gap

$$O(10^{-2}) \times \tau_0$$

cf. Levinson & Cerutti 18

$$\epsilon_{\text{min}} = 10^{-8} m_e c^2$$

$$\tau_0 = 10$$

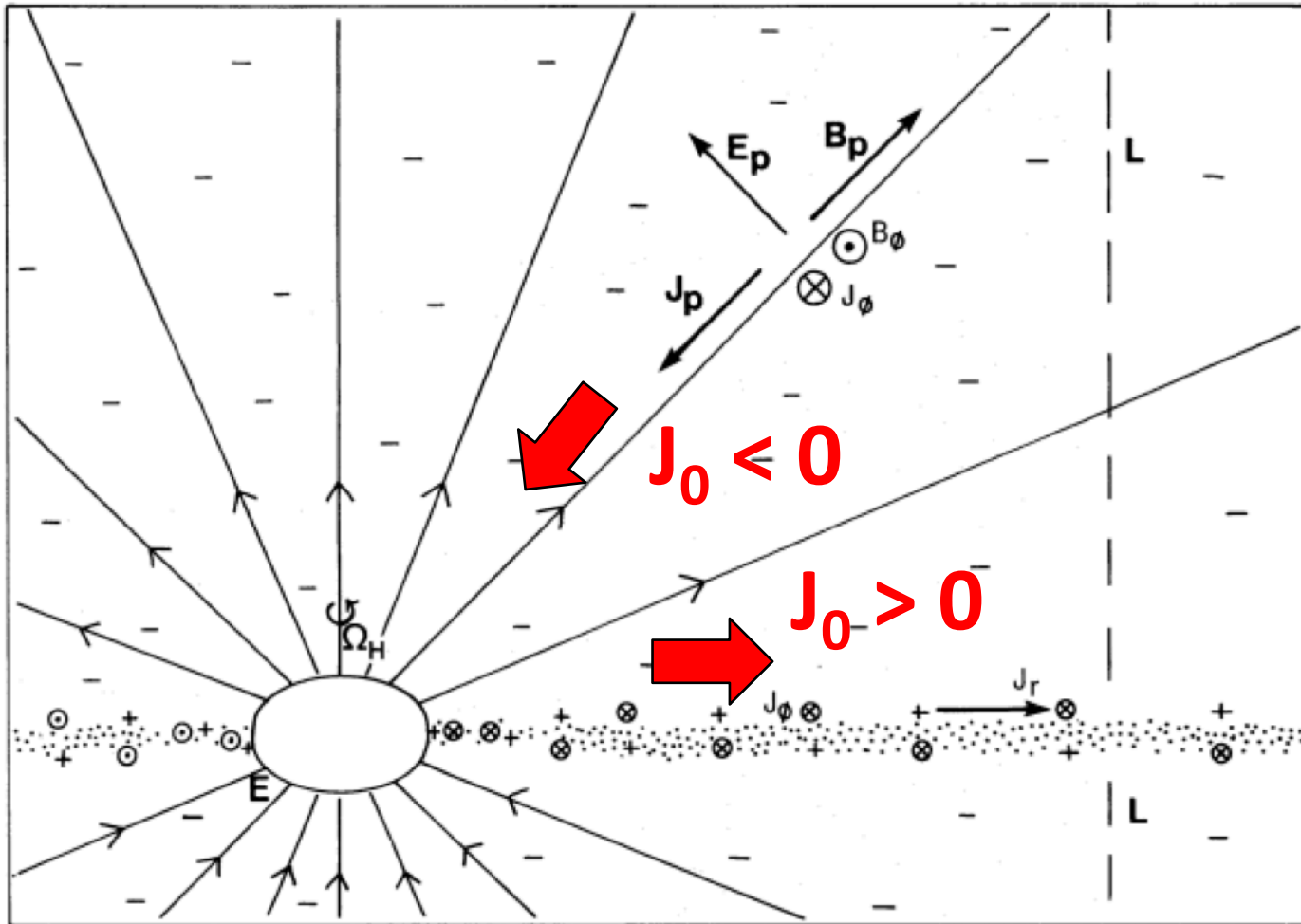
Multiplicity $O(10^{-3})$

$$\epsilon_{\text{min}} = 10^{-9} m_e c^2 \quad \rightarrow \quad \tau_0 \gtrsim 100$$

Electric Currents

Polar region : $J_0 < 0$

Equatorial region : $J_0 > 0$



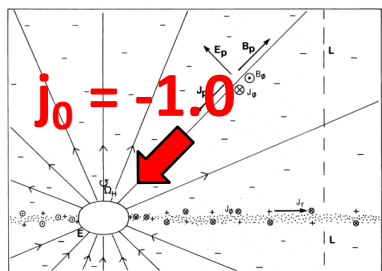
Blandford & Znajek 77

$$j_0 = -1.0, 1.0$$

Results

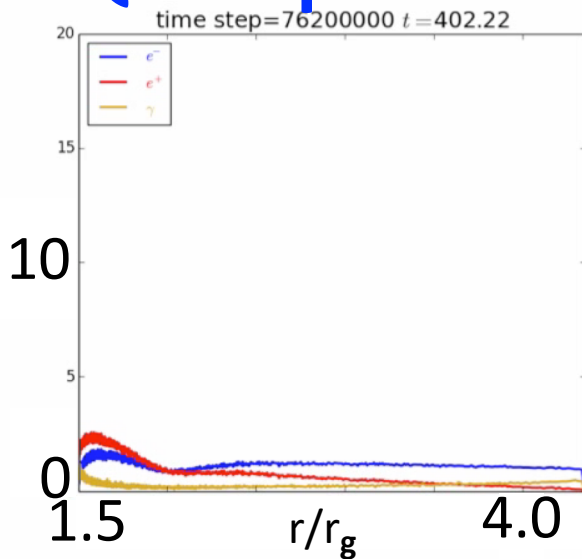
$\tau_0 = 300$

Quasi-periodic oscillation



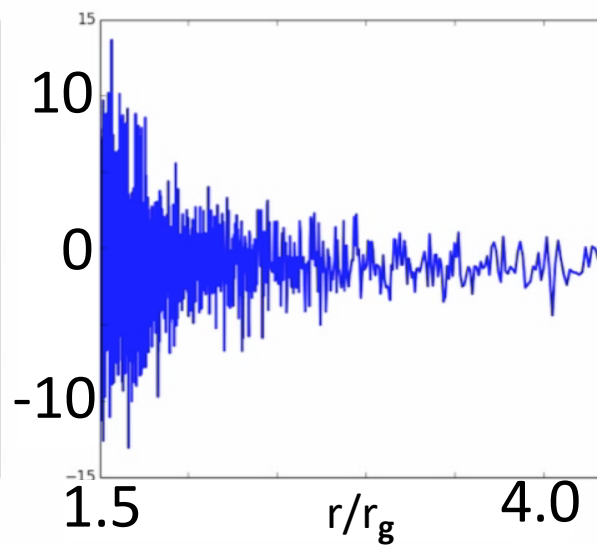
Number density

$$\Delta(n/n_{\text{GJ}})$$



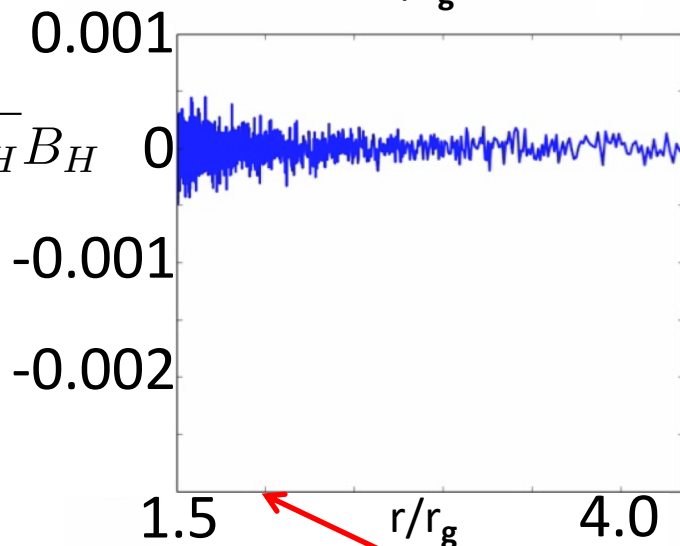
Current density

$$\Sigma j^r / |J_0|$$



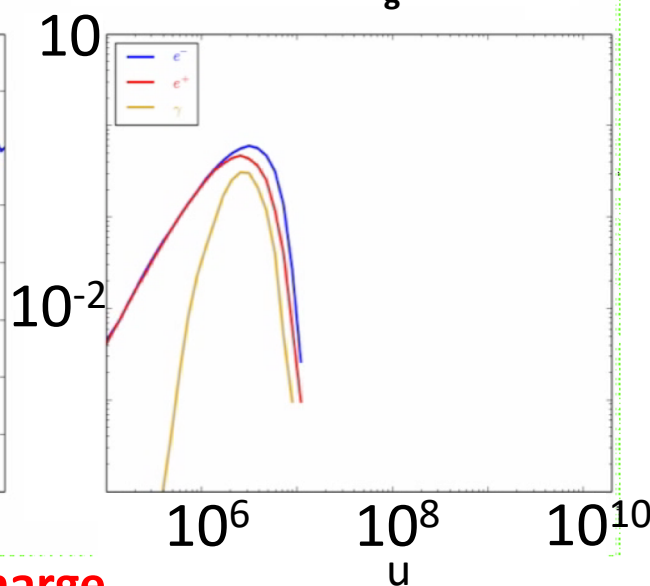
E-field

$$\sqrt{A} E_r / \sqrt{A_H} B_H$$



Energy spectrum

$$u^2 \frac{dN}{du}$$



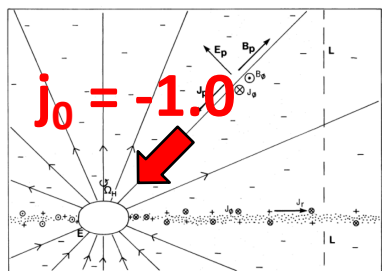
$$\rho_{\text{GJ}} = \frac{B_H \sqrt{A_H}}{4\pi \sqrt{-g}} \left[\frac{\sin^2 \theta}{\alpha^2} (\omega - \Omega) \right]_{,\theta}$$

Null charge surface ($\rho_{\text{GJ}}=0$)

Results

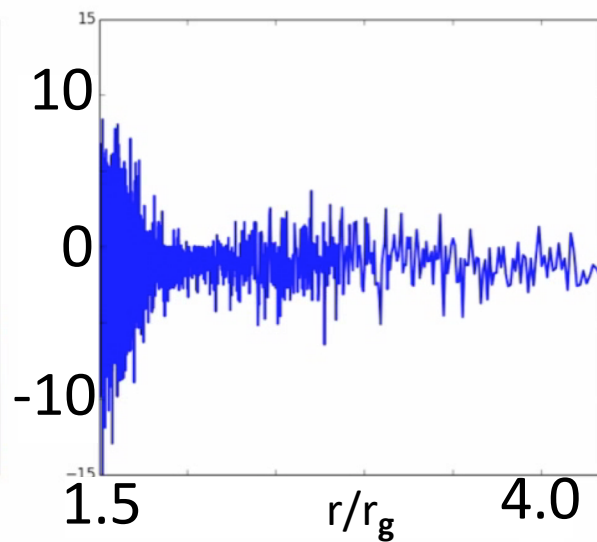
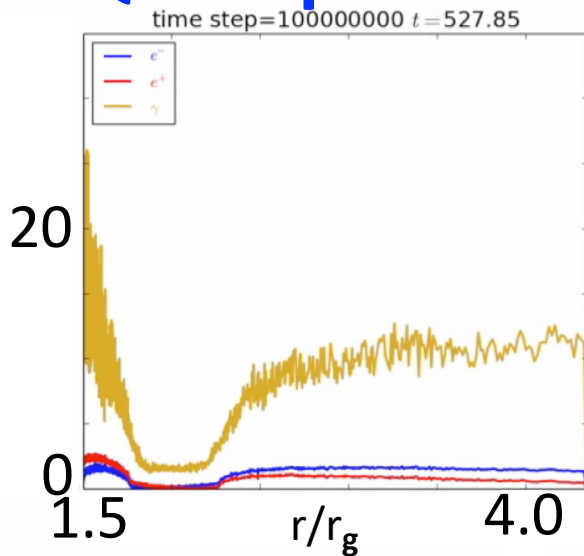
$\tau_0 = 100$

Quasi-periodic oscillation



Number density

$$\Delta(n/n_{\text{GJ}})$$

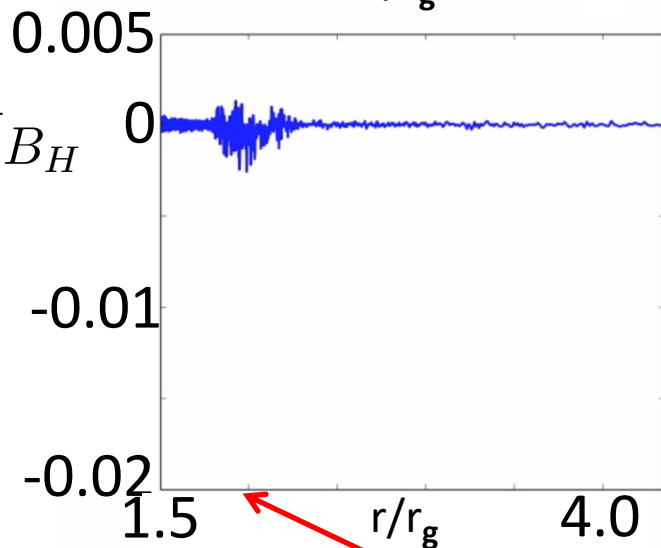


Current density

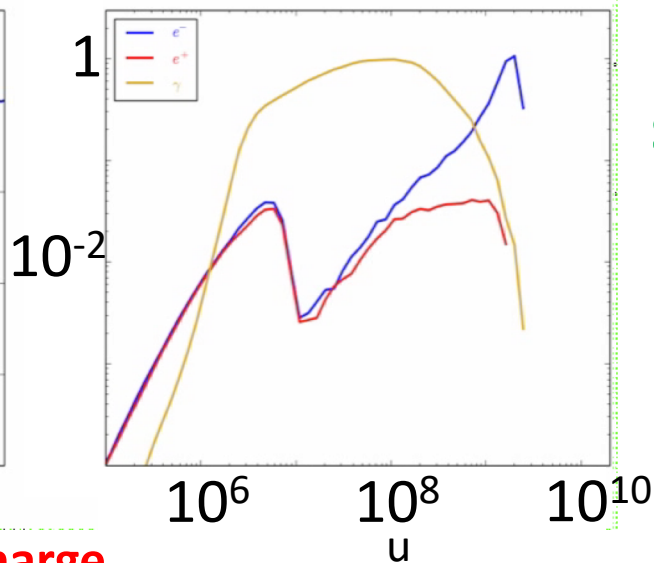
$$\Sigma j^r / |J_0|$$

E-field

$$\sqrt{A} E_r / \sqrt{A_H} B_H$$



Energy spectrum



$$u^2 \frac{dN}{du}$$

$$\rho_{\text{GJ}} = \frac{B_H \sqrt{A_H}}{4\pi \sqrt{-g}} \left[\frac{\sin^2 \theta}{\alpha^2} (\omega - \Omega) \right]_{,\theta}$$

Null charge surface ($\rho_{\text{GJ}}=0$)

Acceleration model

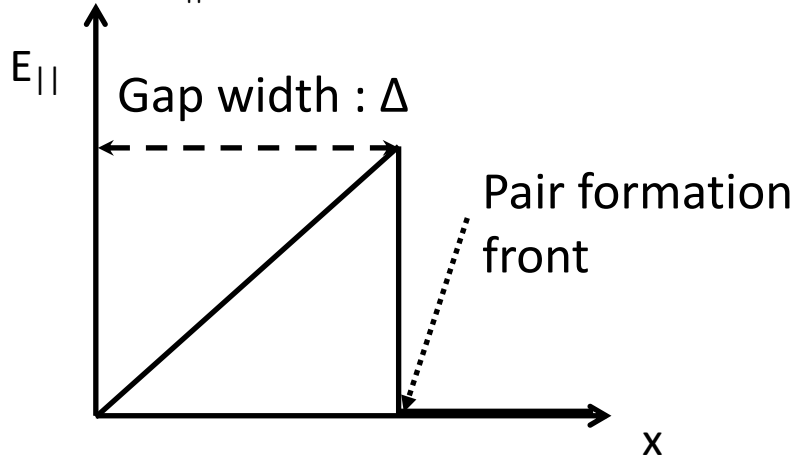
For simplicity, GR effects are neglected.

Considering the gap until the pair formation front.

Electric field

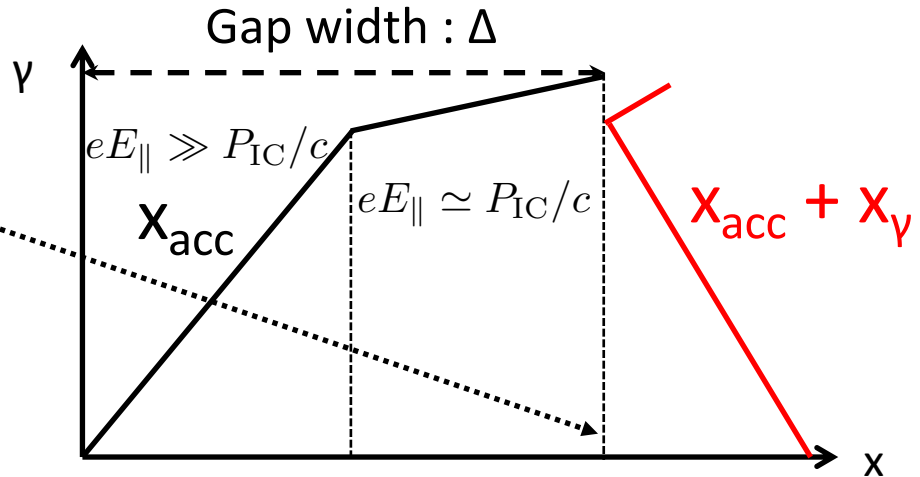
$$\nabla \cdot E_{\parallel} = -4\pi(\rho_e - \rho_{GJ})$$

$$\rightarrow E_{\parallel} = 2\Omega B \xi x / c$$



Lorentz factor

$$\frac{d(\gamma m_e c)}{dt} = eE_{\parallel} - \frac{P_{IC}}{c} - \frac{P_{cur}}{c}$$



$$\xi \equiv (\rho_{GJ} - \rho_e) / \rho_{GJ}, \quad \rho_{GJ} \sim \Omega B / (2\pi c)$$

Gap width Δ is the minimum value of the sum of the acceleration length $x_{acc}(\gamma)$ and the mean free path $x_{\gamma}(\gamma)$.

$$\rightarrow \Delta / r_g \sim 9 \times 10^{-3} \tau_2^{-\frac{1}{3}} \epsilon_{\min, -9}^{-\frac{1}{3}} a_{0.9}^{-\frac{1}{3}} B_4^{-\frac{1}{3}} M_9 \xi^{-\frac{1}{3}} \quad B_2 \equiv B / 10^2 \text{ G}$$

$$\rightarrow E_{\parallel} / B \sim 2 \xi a \Delta / r_g \sim 2 \times 10^{-2} \tau_2^{-\frac{1}{3}} \epsilon_{\min, -9}^{-\frac{1}{3}} a_{0.9}^{\frac{2}{3}} B_2^{-\frac{1}{3}} M_9 \xi^{\frac{2}{3}}$$

Summary

We perform 1D GRPIC simulation for pair cascade in the black hole magnetosphere.