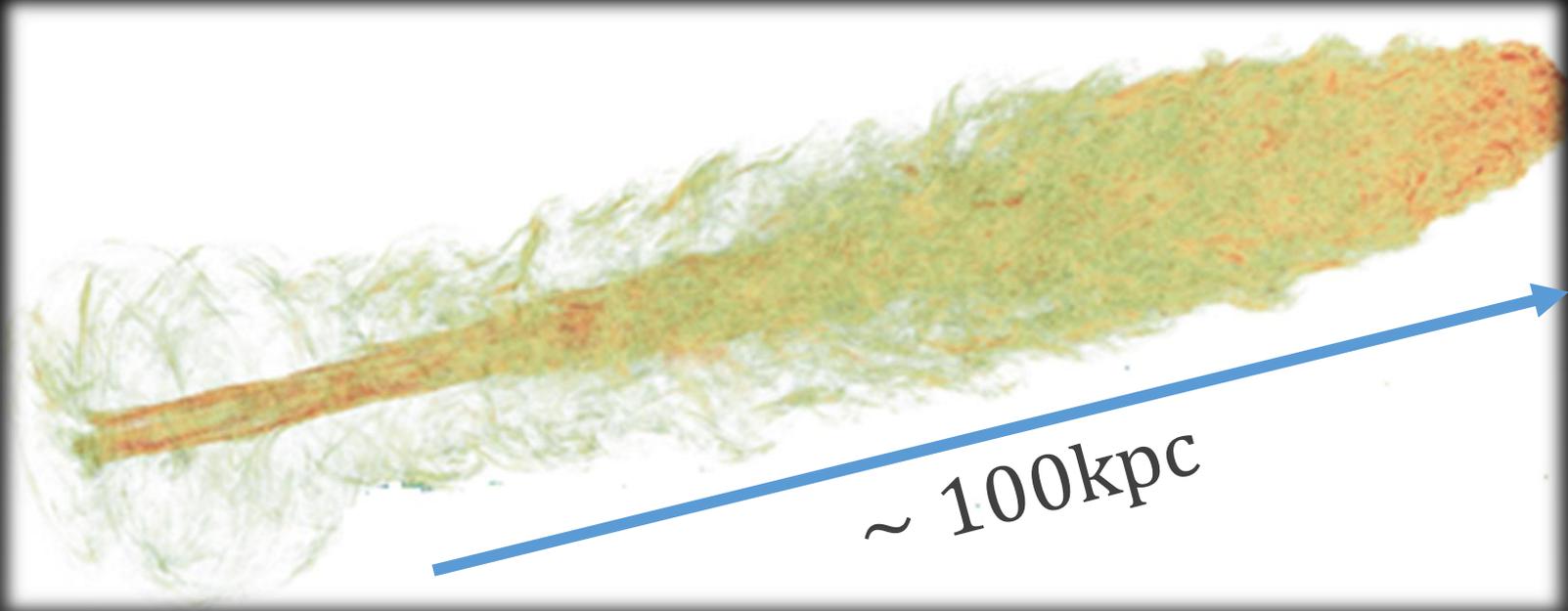


Preliminary

Three dimensional modeling of FR-II jets using two-temperature MHD simulation



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1. Introduction - FR II Radio Galaxies -

Most of FR II sources have powerful jets

- Edge brightened sources (hotspot, Radio-lobe)
Synchrotron radiation from relativistic electrons in the magnetic fields of lobes.
- Diffuse thermal X-ray emitted like a rugby ball.



Blue:X-ray, Red: Radio

1. Introduction - Two-temperature plasma - X-ray cavity in the galaxy cluster

-> Jet plasma is hotter and lighter than ICM

$$[n_{\text{ICM}} \sim 10^{-2} - 10^{-3} \text{cm}^{-3}, T_{\text{ICM}} \sim 10^7 - 10^8 \text{K}]$$

- Relaxation time scale of electrons and ions through Coulomb coupling is longer than the propagation time scale.

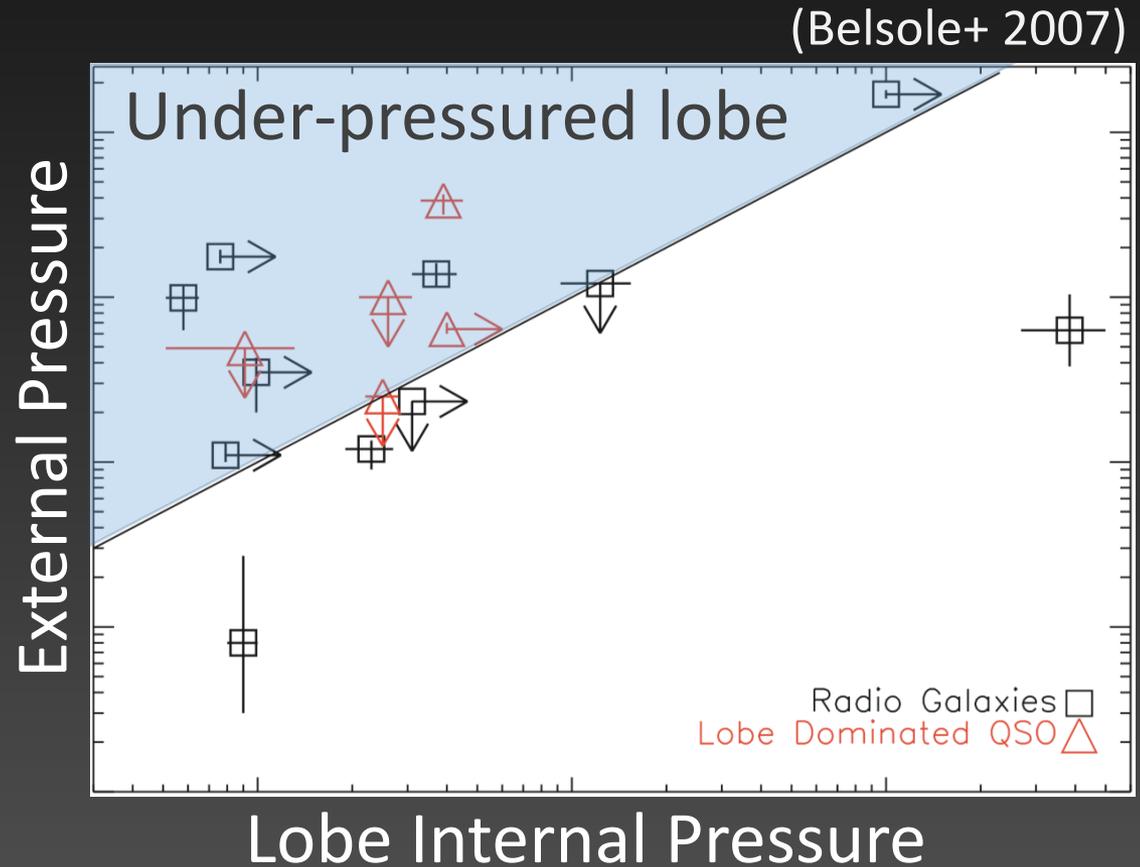
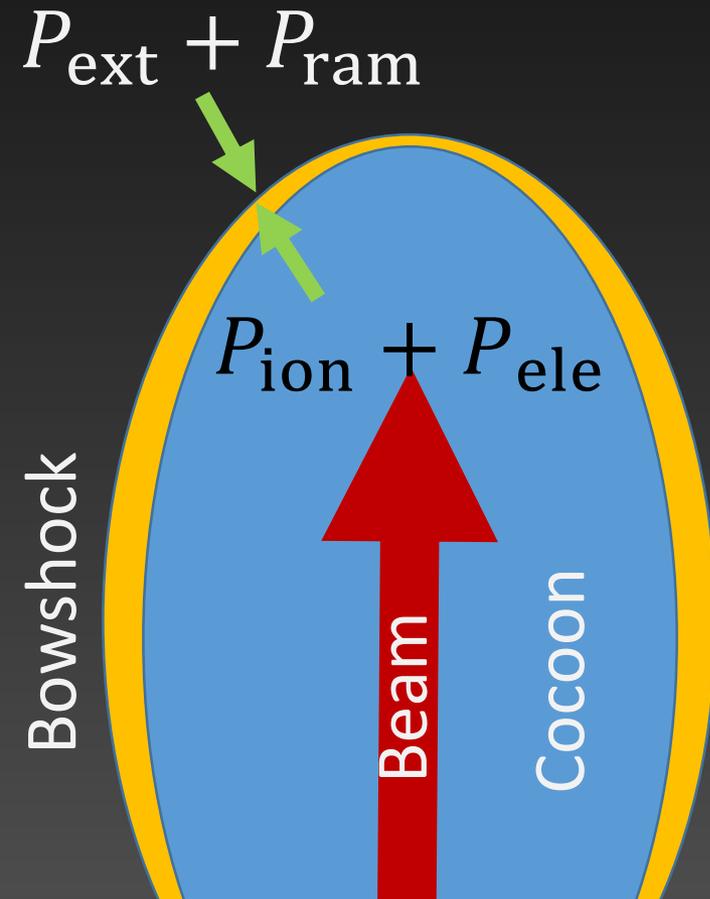
$$\tau_{\text{ei}} \geq \tau_{\text{dyn}} \quad \tau_{\text{ei}} \propto n^{-1} T_e^{3/2}$$

- Electrons and ions could be decoupled in jets.

 $p_{\text{electron}} \neq p_{\text{ion}}$

1. Introduction - Energetics of radio-lobe -

- Pressure balance of radio-lobe, $P_{\text{ext}} \leq P_{\text{ion}} + P_{\text{ele}}$
- The under-pressured lobe indicate the existence of the non-radiating ions in the FR-II jets.

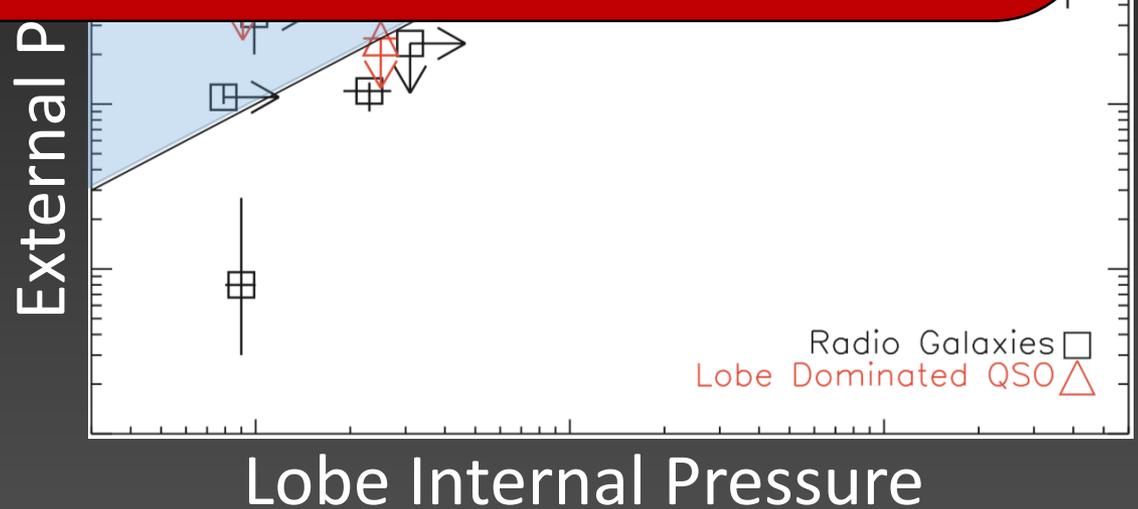


1. Introduction - Energetics of radio-lobe -

- Pressure balance of radio-lobe, $P_{\text{ext}} \leq P_{\text{ion}} + P_{\text{ele}}$

- We can only estimate electron energy through radiation.

- The ion energy are undetectable with current observational techniques!!



1. Introduction - Motivations -

We investigate

- the spatial energy distribution of electron and ion through two-temperature MHD jets simulation.
- the pressure balance between jet and surrounding medium, especially possibility of electron under-pressured lobe.

Method

- Basic equations
- The fraction of electron heating
- Model

2. Method - Basic equations -

Assumption

- Electron-Ion plasma, $n = n_i = n_e$
- One fluid approximation
- Electron and Ion have distinct temperature

MHD Eqs. + Electron and Ion. energy Eqs.

$$\left[\begin{array}{l} nT_i \frac{ds_i}{dt} = -q^{ie} + (1 - f_e)q^{\text{heat}} \\ nT_e \frac{ds_e}{dt} = +q^{ie} + f_e q^{\text{heat}} - q^{\text{rad}} \end{array} \right.$$

[Colulomb] [Dissipation] [Radiation]

2. Method - The fraction of electron heating -

Model 1: Turbulent dissipation into the electron and ions.

(Kawazura+ 2019)

$$f_e = f_e(T_e/T_i, \beta_i)$$

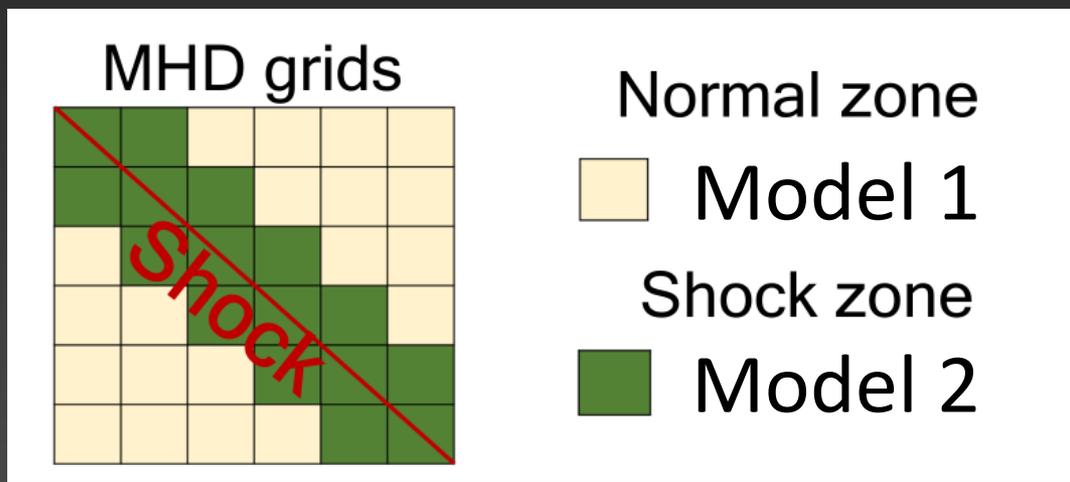
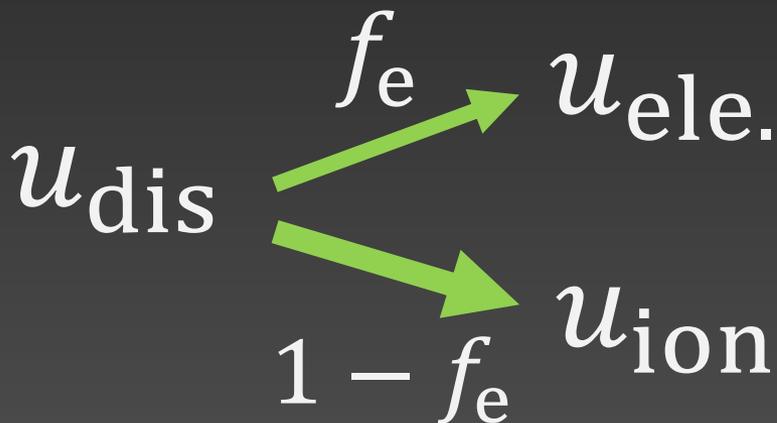
Magnetic energy dominant region : $f_e \rightarrow 1$ (electron heating),

Thermal energy dominant region : $f_e \rightarrow 0$ (ion heating)

Model 2: Dissipation in collisionless shock.

$f_e = 0.05$ Shocks heat ions (e.g., Ghavamian+2013, Matsukiyo+2010).

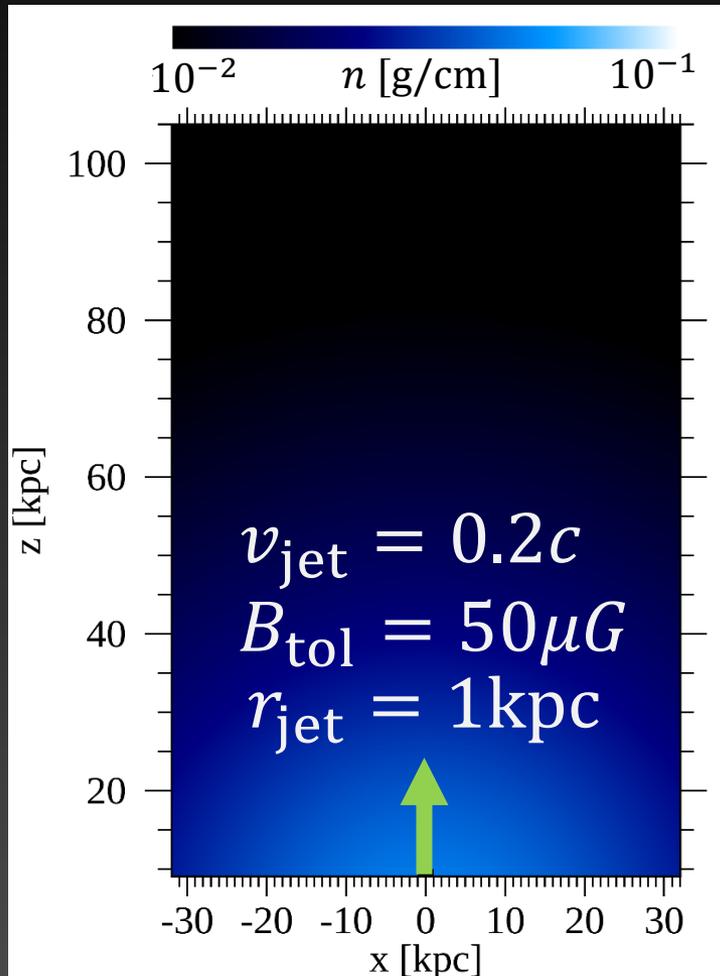
$$0 < f_e < 1$$



2. Method - Model -

MHD CODE : CANS+ (Matsumoto+2019, Ohmura+ 2019)

Uniform Grids: (800,800,1200)



Jets parameter

$$L_{\text{kin}} = 2.0 \times 10^{46} \text{erg/s}$$

$$T_e = T_i = 3.0 \times 10^9 \text{K}$$

Initial ICM profile

β -profile (King+ 1972)

$$\rho(r) = \rho_0 \{1 + (r/r_c)^2\}^{-3\beta/2}$$

$$T_e = T_i = 2.3 \times 10^7 \text{K (2KeV)}$$

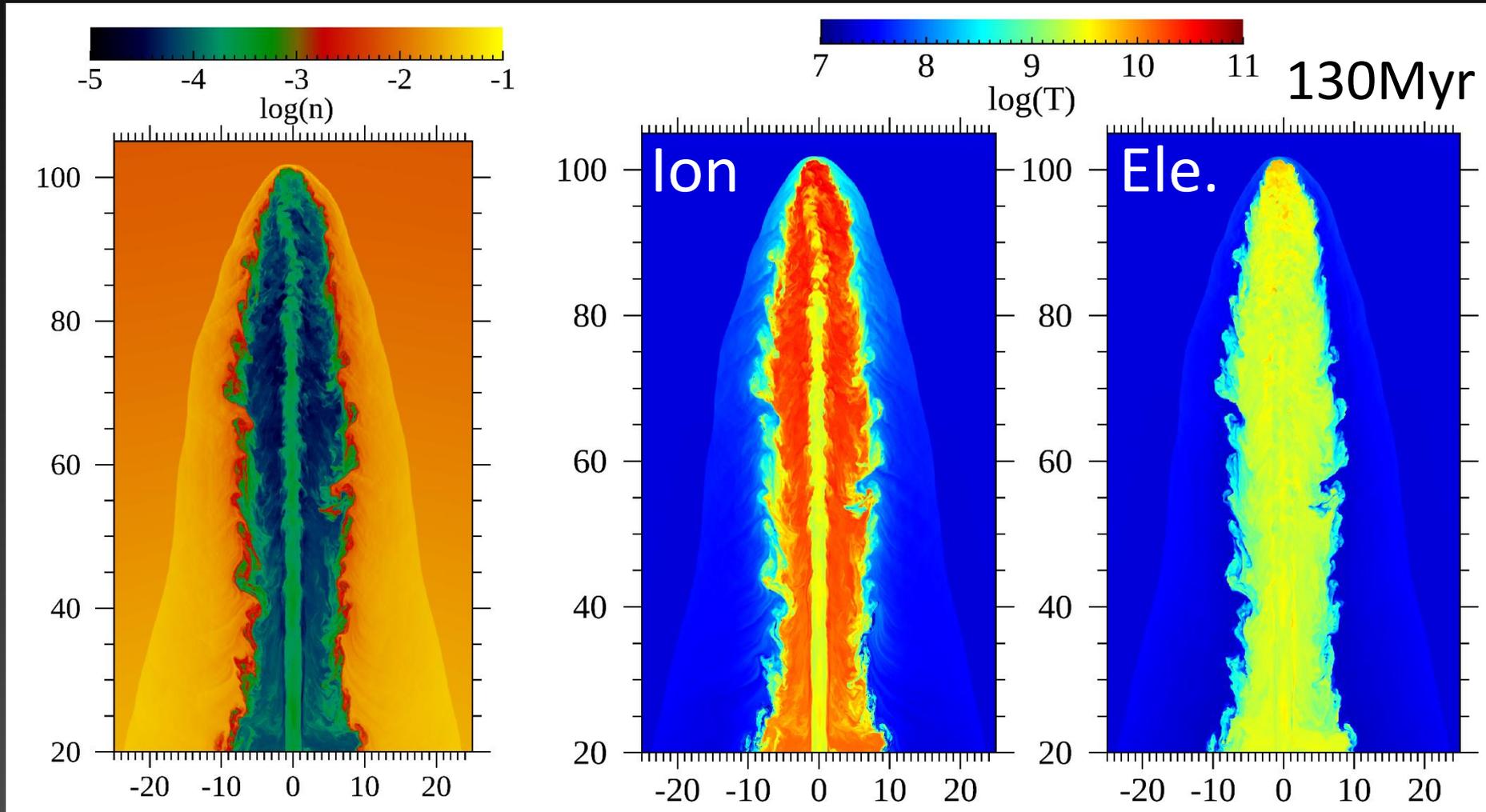
$$\rho_0 = 8.35 \times 10^{-26} \text{g/cc}, r_c = 30\text{kpc}, \beta = 1/2$$

Results

- Overview
- Energetics in jets
- Under-pressured lobe

3. Results - Overview -

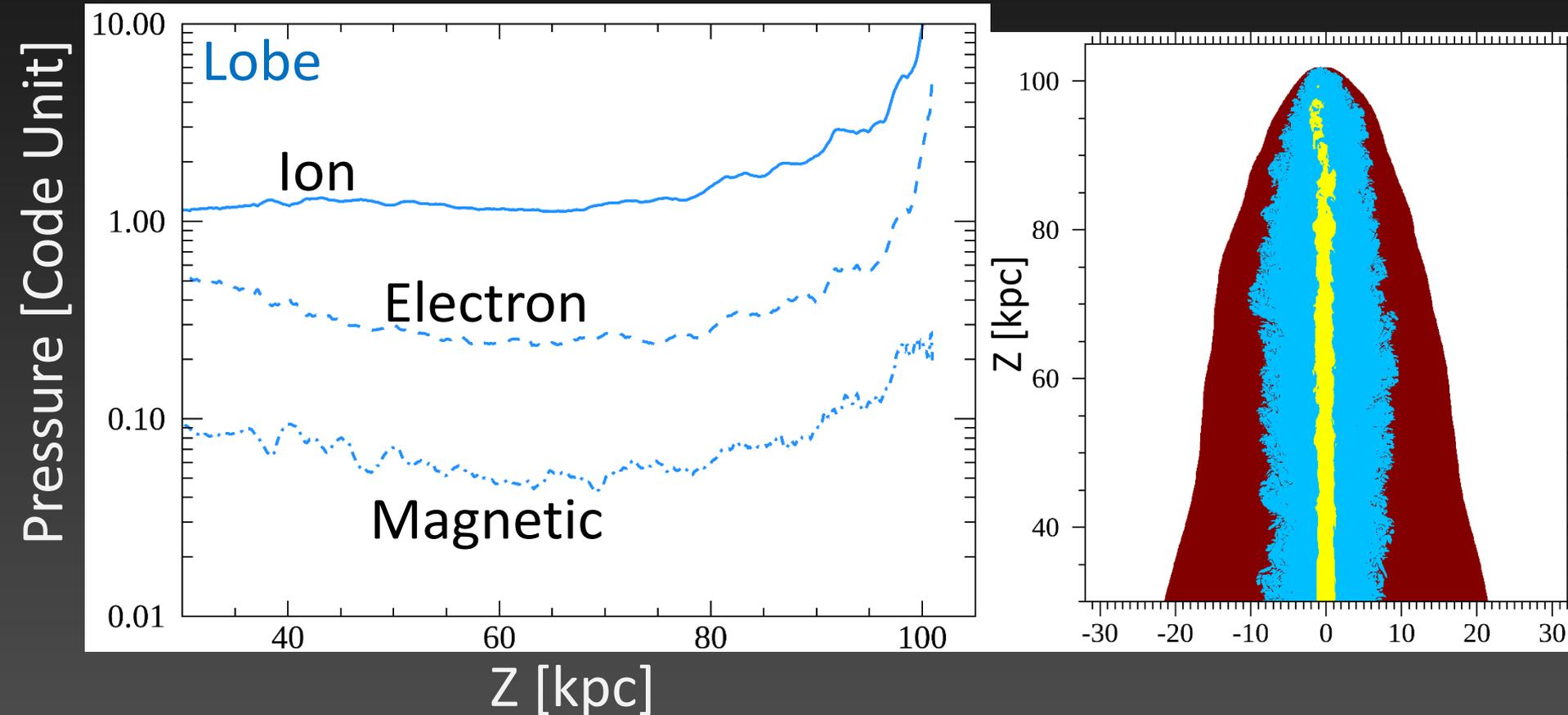
- Beam can reach termination, and form the terminal shocked.
- The most of heating energy goes to ions.



3. Results - Energetics in jets -

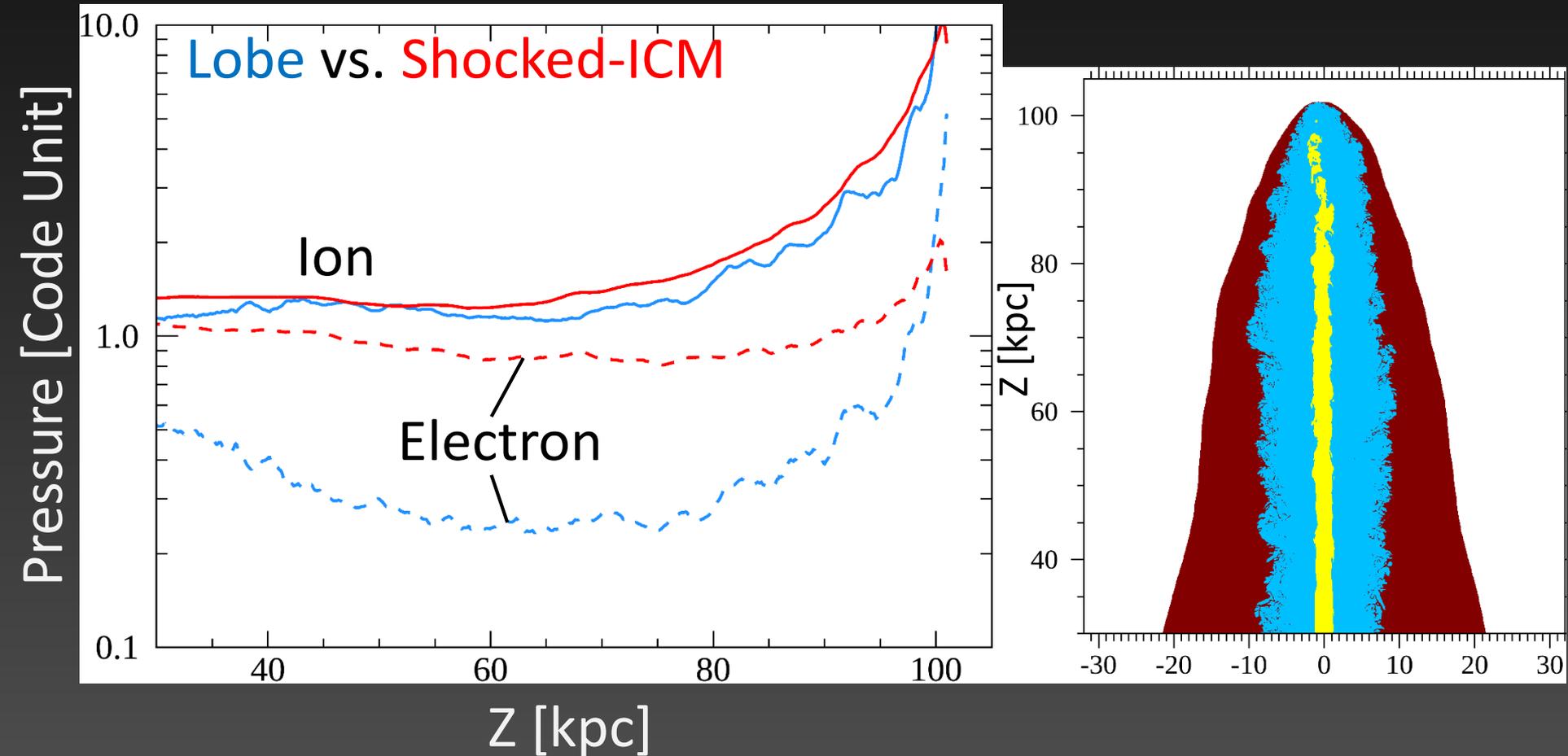
Volume-weighted pressure in radio-lobe.

$$P_{i,e,B}(z) = \frac{\int p_{i,e,B} 2\pi r dr dz}{\int 2\pi r dr dz} \quad \rho_B = B^2 / 8\pi$$



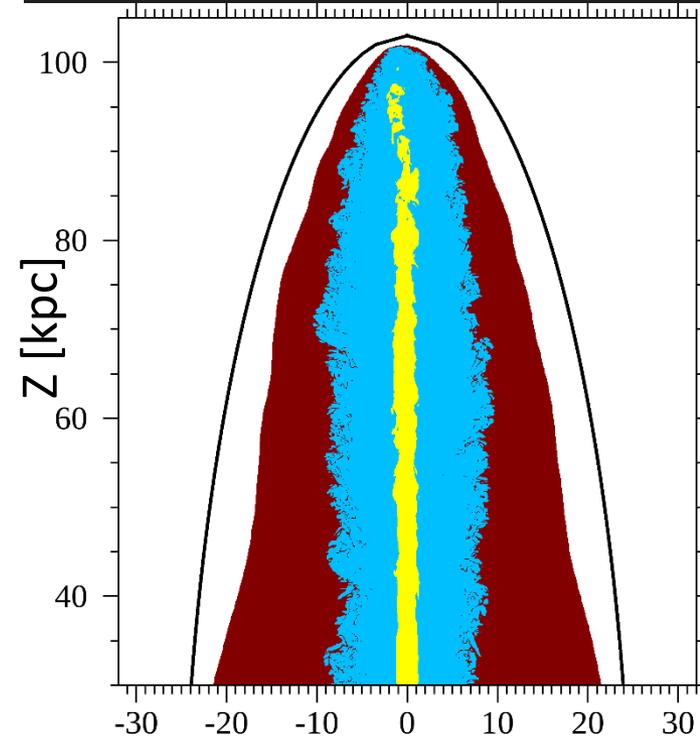
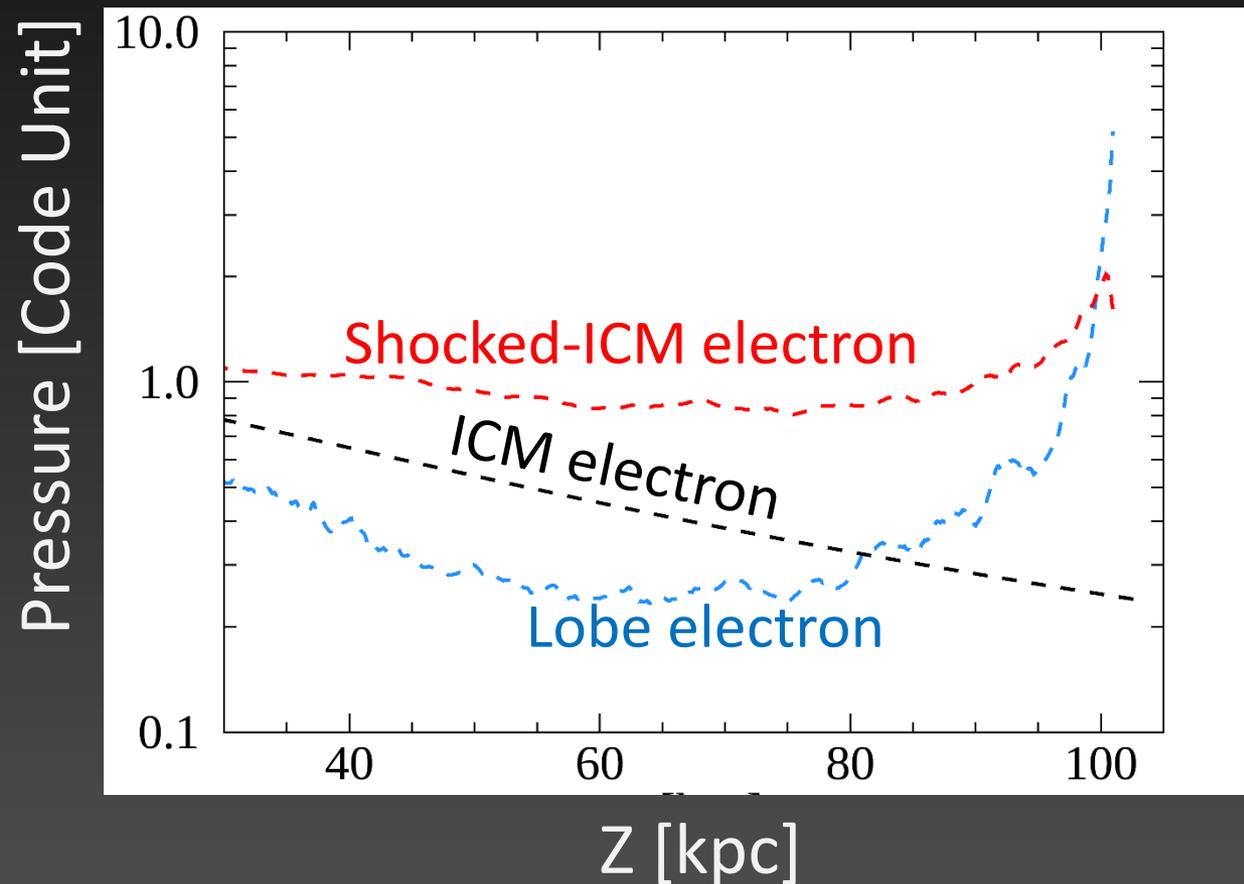
3. Results - Energetics in jets -

- In shocked-ICM, ions and electrons decouple by bowshock. But, Coulomb interaction is effective because density is high.
- The lobe electron pressure is lower than shocked-ICM.



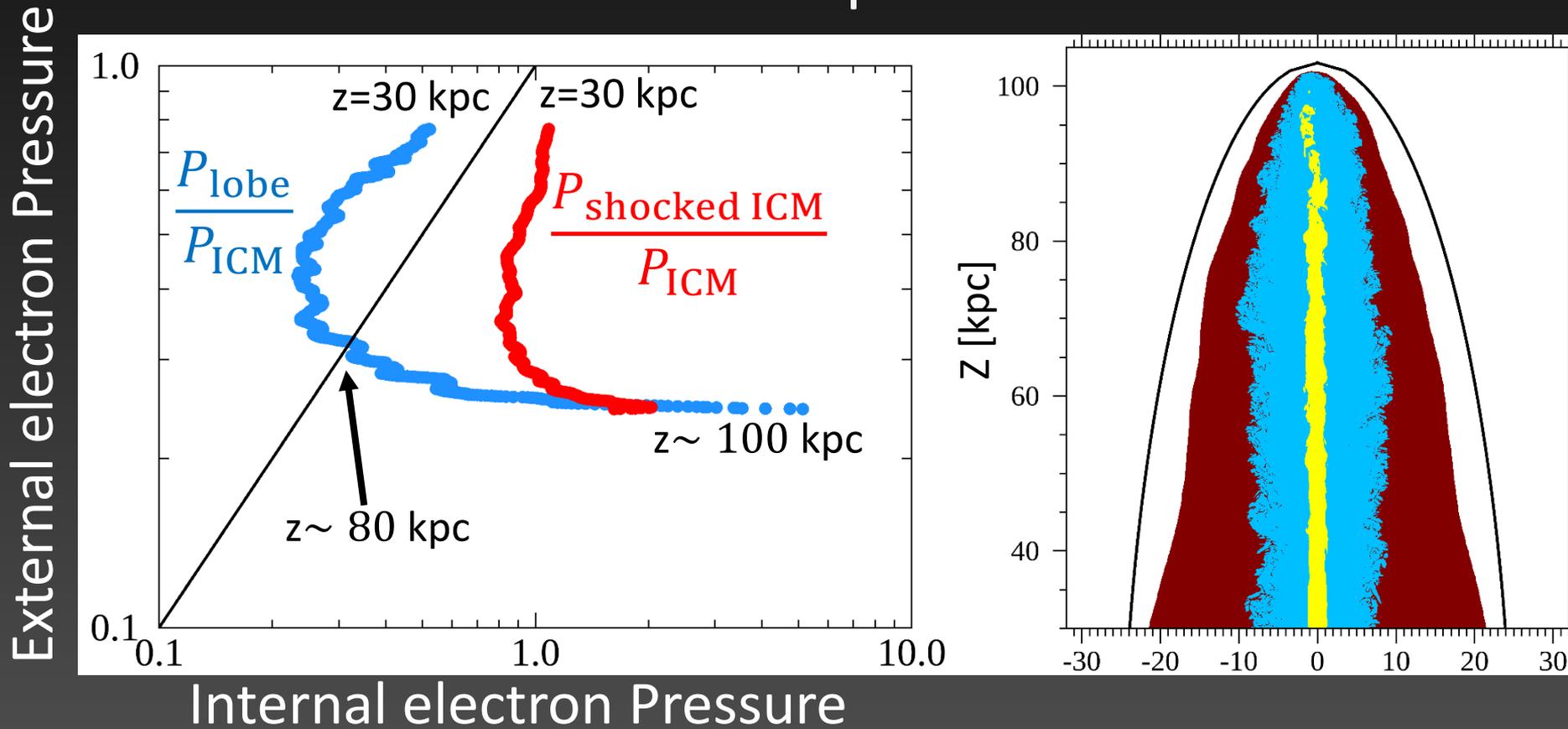
3. Results - Energetics in jets -

- Shocked-ICM electron pressure is higher than surrounding electron pressure everywhere.
- Lobe electrons have lower energy than ICM electrons ($z < 80$).



3. Results - Under-pressured lobe -

- Shocked-ICM electron pressure is higher than surrounding electron pressure everywhere.
 - On the other hand, lobe electron pressure is low.
- Under-pressured lobe

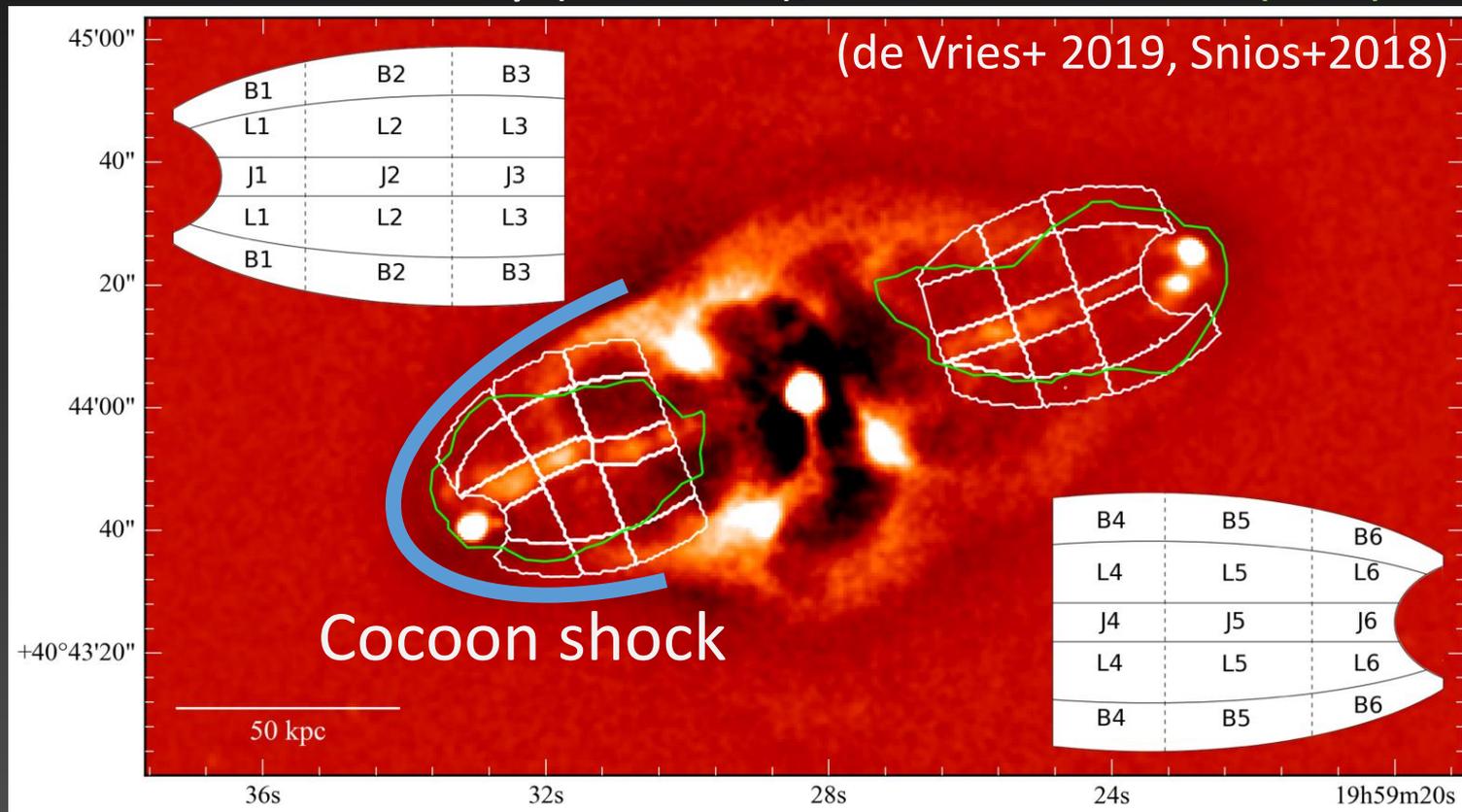


Discussion

4. Discussion - Under-pressured lobe -

- The bow (cocoon) shocks of Cyg.A are clearly seen in high-resolution Chandra observations.
- Shock compression is clearly observed at these shocks.

0.5-7.0 KeV X-ray (Chandra), **contour 4.5 GHz (VLA)**



4. Discussion - Under-pressured lobe -

- However, the pressure of the eastern lobe, which is estimated by X-ray Inverse Compton, is lower than the rim pressure, *i.e.*, **Under-pressured lobe**.

- These observational analysis indicates that

$$P_{\text{lobe}} < P_{\text{ICM}} < P_{\text{shocked-ICM}}$$

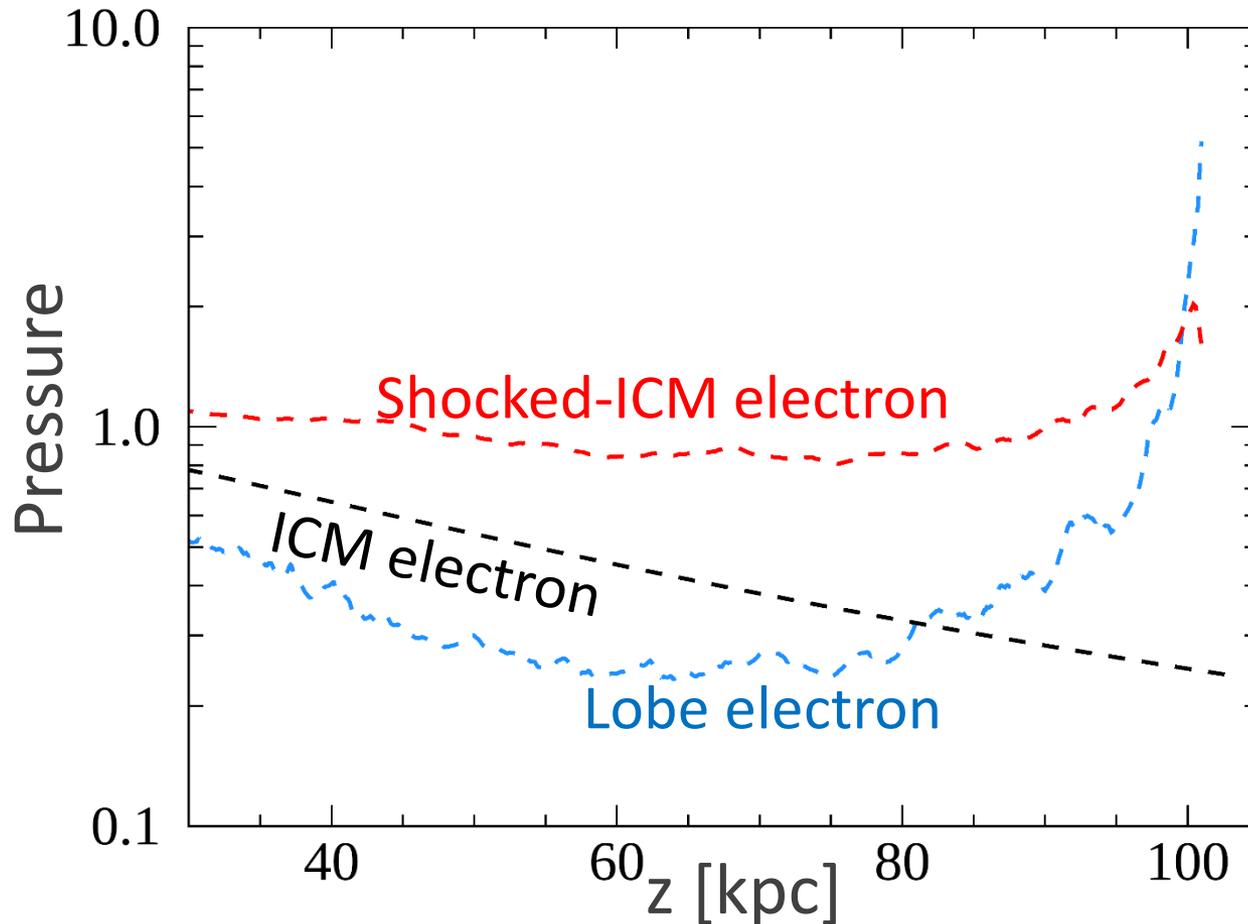
	$P_{\text{rim}}^{\text{a}}$ (10^{-10} erg cm^{-3})	$\gamma_{\text{min}}^{\text{b}}$	P_{IC}^{c} (10^{-10} erg cm^{-3})
East	10.4 ± 0.4	1 10	$5.8^{+2.0}_{-1.4}$ $2.2^{+0.4}_{-0.4}$
West	8.4 ± 0.2	1 10	140^{+1690}_{-116} 18^{+40}_{-10}

(de Vries+ 2019)

4. Discussion - Under-pressured lobe -

Our two-temperature jets model could explain this pressure condition.

$$P_{\text{lobe}} < P_{\text{ICM}} < P_{\text{shocked-ICM}}$$



5. Summary

- We have carried out high-resolution two-temperature MHD jets propagate simulations.
- Beam, which propagating up to 100 kpc, can reach termination of jets, and terminal shock (hotspot) is formed.
- The magnetic energy is much less than the ion thermal energy, and hence most of heating energy goes to ions.
- The lobe electron pressure is lower than that of shocked-ICM and surrounding ICM, and under-pressured lobe is formed .